

# TENSOR INVARIANTS OF SUBGROUPS OF THE LORENTZ GROUP

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**Abstract**

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**MATHEMATICS**

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## TENSOR INVARIANTS OF SUBGROUPS OF THE LORENTZ GROUP

*(Presented by Academician L. I. Sedov, December 2, 1968)*

For every subgroup  $\mathcal{L}_k$  of the Lorentz group  $\mathcal{L}$  that can be specified by the requirement of invariance of some finite set of tensors  $H_{(1)}, \dots, H_{(n)}$ , any other tensor invariant  $T$  of it, of arbitrary rank, can be represented in the form of a linear combination of tensors  $T_{(1)}, \dots, T_{(p)}$  constructed by means of the operations of tensor multiplication and contraction from the basic tensors  $H_{(1)}, \dots, H_{(n)}$ :

$$T = k_1 T_{(1)} + \dots + k_p T_{(p)},$$

where the tensors  $T_{(1)}, \dots, T_{(p)}$  form a basis of the linear space of tensors of rank  $r$  invariant with respect to the group  $\mathcal{L}_k$ , the dimension of which is denoted by  $p$ .\*

Thus, the construction of tensor invariants for subgroups of the Lorentz group  $\mathcal{L}$  reduces to finding the basic tensors  $\{H_m\}$  for all subgroups  $\mathcal{L}_k$  of the Lorentz group  $\mathcal{L}$  that can be specified with the aid of tensors. All these subgroups  $\mathcal{L}_k$  are Lie subgroups of the Lorentz group with a finite component group, whose classification is given in <sup>(6)</sup>. The method for choosing all subgroups  $\mathcal{L}_k$  and obtaining their basic tensors  $\{H_m\}$  is indicated in <sup>(5)</sup>.

Table 1 lists all subgroups  $\mathcal{L}_k$  of the Lorentz group  $\mathcal{L}$  that are specified with the aid of tensors, and for all these subgroups the corresponding sets of tensors are written out. In some cases the group can be specified by a smaller number of invariant tensors; however, the tensors indicated are sufficient to obtain the entire algebra of tensor invariants.

At the beginning of each series of formulas the infinitesimal operators of the maximal connected subgroup  $\mathcal{G}_l$  of the group  $\mathcal{L}_k$  under consideration are indicated, while the symbols of A. V. Shubnikov\*\* are used to denote its finite component group <sup>(6)</sup>.

The infinitesimal operators  $X_q$  of the proper Lorentz group of transformations of the variables  $x^i$  ( $i = 1, 2, 3, 4$ ), which leave invariant the quadratic form  $x^{12} + x^{22} + x^{32} - x^{42}$ , have the form

$$X_1 = x^2\partial_3 - x^3\partial_2, \quad X_2 = x^3\partial_1 - x^1\partial_3, \quad X_3 = x^1\partial_2 - x^2\partial_1,$$

$$X_4 = x^1\partial_4 + x^4\partial_1, \quad X_5 = x^2\partial_4 + x^4\partial_2, \quad X_6 = x^3\partial_4 + x^4\partial_3,$$

where  $\partial_i = \partial/\partial x^i$ .

The vectors  $e_1, e_2, e_3, e_4$  form an orthonormal basis. Products and powers of the vectors  $e_i$  are understood as tensor products,  $e_{ij\dots k} = e_i e_j \dots e_k$ . The symbols  $[\ ]$ ,  $(\ )$  attached to indices denote alternation and symmetrization.

\* Basic tensors  $\{H_{(m)}\}$  and tensor bases  $\{T_{(s)}\}$  of the first, second, third, and fourth ranks for subgroups of the group  $O(3)$  were established in  $(1-4)$ .

\*\*  $n$  denotes a rotation about the  $x^3$  axis through the angle  $2\pi/n$ ;  $1:2$  denotes a rotation about the  $x^2$  axis through the angle  $\pi$ ; reflection of  $x^3$  is denoted by the symbol  $:m$ , or by a bar over the reflected  $x^3$ ; reflection of  $x^2$  by the symbol  $\cdot m$ , reflection of  $x^4$  by a subscript bar,  $\cdot m'$  denotes reflection of  $x^1$ ,  $1:2'$  denotes rotation about the  $x^1$  axis through the angle  $\pi$ .  $3/2$ ,  $3/4$ , and  $3/5$  denote the symmetry groups of the tetrahedron, octahedron, and icosahedron.

**Table 1**

Subgroup	Parameter	Tensor invariants	Subgroup	Parameter	Tensor invariants
$\mathcal{G}_1(X_1, X_2, X_3, X_4, X_5, X_6)$	$\underline{\underline{2}}$	$\mathfrak{g}, E$		$1 \cdot \underline{m}$	$\mathfrak{g}, E, \mathbf{e}_4^2, \mathbf{e}_3$
	$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}$		$1 \cdot m$	$\mathfrak{g}, \mathbf{e}_4^2, \mathbf{e}_3$
$\mathcal{G}_2(X_1, X_2, X_3)$				$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}, \mathbf{e}_3^2, \mathbf{e}_4$
	1	$\mathfrak{g}, E, \mathbf{e}_4$		$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}, \mathbf{e}_{[34]}, \mathbf{e}_4^2$
	$\frac{\underline{\underline{2}}}{2}$	$\mathfrak{g}, \mathbf{e}_4$		$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}, E, \mathbf{e}_3^2, \mathbf{e}_4^2$
	$\frac{1}{2}$	$\mathfrak{g}, \mathbf{e}_{[123]}$		$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}, \mathbf{e}_{[124]}, \mathbf{e}_4^2$
	$\frac{\underline{\underline{2}}}{2}$	$\mathfrak{g}, E, \mathbf{e}_4^2$		$\frac{\underline{\underline{2}}}{2} \cdot m \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_3^2, \mathbf{e}_4^2$
	$\frac{\underline{\underline{2}}}{2} \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_4^2$	$\mathcal{G}_8(X_6)$		
$\mathcal{G}_3(X_3, X_4, X_5)$	$1 \cdot \underline{m}$	$\mathfrak{g}, E, \mathbf{e}_3$		$n : \underline{m}$	$\mathfrak{g}, E, D_{nh}, \mathbf{e}_{[12]}$
	$1 \cdot m$	$\mathfrak{g}, \mathbf{e}_3$		$\underline{m} \cdot n : \underline{m}$	$\mathfrak{g}, E, D_{nh}, \mathbf{e}_2^2 (n = 1)$
	$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}, E, \mathbf{e}_3^2$		$m \cdot n : \underline{m}$	$\mathfrak{g}, D_{nh}, \mathbf{e}_{[34]}$
	$\frac{\underline{\underline{2}}}{2}$	$\mathfrak{g}, \mathbf{e}_{[124]}$		$n : \underline{m} \times \underline{1}$	$\mathfrak{g}, D_{nh}, \mathbf{e}_{[12]}$
	$\frac{\underline{\underline{2}}}{2} \cdot m$	$\mathfrak{g}, \mathbf{e}_{[124]}$		$m \cdot n : \underline{m} \times \underline{1}$	$\mathfrak{g}, D_{nh}, \mathbf{e}_2^2 (n = 1)$

Subgroup	Parameter	Tensor invariants	Subgroup	Parameter	Tensor invariants	
$\mathcal{G}_4(X_3, X_4 - X_2, X_1 + X_5)$	$\frac{2}{2} \cdot m \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_3^2$		$2n : \underline{m}$	$\mathfrak{g}, \mathbf{e}_{[34]} D_{nh}, \mathbf{e}_{[12]}$	
				$m \cdot 2n : \underline{m}$	$\mathfrak{g}, \mathbf{e}_{[34]} D_{nh}$	
		$1$	$\mathfrak{g}, E, \mathbf{e}_3 + \mathbf{e}_4$	$\mathcal{G}_9(X_4 - X_2)$	$1$	$\mathfrak{g}, E, \mathbf{e}_2, \mathbf{e}_3 + \mathbf{e}_4$
		$1 : \underline{2}$	$\mathfrak{g}, \mathbf{e}_{[123]} + \mathbf{e}_{[124]}$		$2$	$\mathfrak{g}, E, \mathbf{e}_2^2, \mathbf{e}_3 + \mathbf{e}_4$
		$\underline{2}$	$\mathfrak{g}, E, (\mathbf{e}_3 + \mathbf{e}_4)^2$		$1 : \underline{2}$	$\mathfrak{g}, \mathbf{e}_{[123]} + \mathbf{e}_{[124]}, \mathbf{e}_2$
$\mathcal{G}_5(X_3, X_6)$	$1 \cdot m$	$\mathfrak{g}, \mathbf{e}_3 + \mathbf{e}_4$		$\underline{2}$	$\mathfrak{g}, E, \mathbf{e}_2^2, \mathbf{e}_2(\mathbf{e}_3 + \mathbf{e}_4)$	
	$\frac{1}{2} \cdot m$	$\mathfrak{g}, (\mathbf{e}_3 + \mathbf{e}_4)^2$		$\frac{2}{2} : \underline{2}$	$\mathfrak{g}, \mathbf{e}_{[123]} + \mathbf{e}_{[124]}, \mathbf{e}_2^2$	
				$1 : \underline{m}$	$\mathfrak{g}, E, (\mathbf{e}_3 + \mathbf{e}_4)^2 \mathbf{e}_2$	
		$\underline{2}$	$\mathfrak{g}, E, \mathbf{e}_{[12]}$		$2 : \underline{m}$	$\mathfrak{g}, E, \mathbf{e}_2^2, (\mathbf{e}_3 + \mathbf{e}_4)^2$
	$\frac{2}{2} \cdot m$	$\mathfrak{g}, E, \mathbf{e}_3^2 - \mathbf{e}_4^2$		$1 \cdot m$	$\mathfrak{g}, \mathbf{e}_{[134]}, \mathbf{e}_3 + \mathbf{e}_4$	
	$\frac{2}{2} \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_{[12]}$		$2 \cdot m$	$\mathfrak{g}, \mathbf{e}_2^2, \mathbf{e}_3 + \mathbf{e}_4$	
	$\underline{2} \cdot m$	$\mathfrak{g}, \mathbf{e}_{[34]}$		$m' \cdot 1 : \underline{m}$	$\mathfrak{g}, (\mathbf{e}_3 + \mathbf{e}_4)^2, \mathbf{e}_2$	
	$\underline{2} \cdot m \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_3^2 - \mathbf{e}_4^2$		$m \cdot 1 : \underline{m}$	$\mathfrak{g}, \mathbf{e}_{[134]}, (\mathbf{e}_3 + \mathbf{e}_4)^2$	
	$\mathcal{G}_6(X_4 - X_2, X_1 + X_5)$				$m \cdot 2 : \underline{m}$	$\mathfrak{g}, \mathbf{e}_2^2, (\mathbf{e}_3 + \mathbf{e}_4)^2$
		$n$	$\mathfrak{g}, E, \mathbf{e}_3 + \mathbf{e}_4, \omega_n$		$\underline{2} \cdot m$	$\mathfrak{g}, \mathbf{e}_2^2, \mathbf{e}_{[14]} + \mathbf{e}_{[13]}$
$\underline{2n}$		$\mathfrak{g}, E, (\mathbf{e}_3 + \mathbf{e}_4)^2, \omega_n (n \text{ odd}), \omega_{2n}, (\mathbf{e}_3 + \mathbf{e}_4) \omega_n (n \text{ even})$		$1 : \underline{2}'$	$\mathfrak{g}, \mathbf{e}_{[134]}, \mathbf{e}_2(\mathbf{e}_3 + \mathbf{e}_4)$	
$n : \underline{2}$		$\mathfrak{g}, \mathbf{e}_{[123]} + \mathbf{e}_{[124]}, \omega_n$		$1 \cdot m'$	$\mathfrak{g}, \mathbf{e}_2, \mathbf{e}_3 + \mathbf{e}_4$	
$n : \underline{m}$		$\mathfrak{g}, E, (\mathbf{e}_3 + \mathbf{e}_4)^2, \omega_{2n}, (\mathbf{e}_3 + \mathbf{e}_4) \omega_n (n \text{ odd}), \omega_n (n \text{ even})$				

Subgroup	Parameter	Tensor invariants	Subgroup	Parameter	Tensor invariants
	$n \cdot m$	$\mathfrak{g}, \mathbf{e}_3 + \mathbf{e}_4, \omega_n$		$\underline{2} \cdot m'$	$\mathfrak{g}, \mathbf{e}_2^2, \mathbf{e}_2(\mathbf{e}_3 + \mathbf{e}_4)$
	$m \cdot n : \underline{m}$	$\mathfrak{g}, (\mathbf{e}_3 + \mathbf{e}_4)^2, \omega_{2n}, (\mathbf{e}_3 + \mathbf{e}_4)\omega_n$	<b>Finite groups</b> ( $n$ odd), $\omega_n$ ( $n$ even)		
	$\underline{2n} \cdot m$	$\mathfrak{g}, (\mathbf{e}_3 + \mathbf{e}_4)^2, \omega_n$	( $n$ odd), $\omega_{2n}, (\mathbf{e}_3 + \mathbf{e}_4)\omega_n$ ( $n$ even)	$n$	$\mathfrak{g}, E, D_{nh}, \mathbf{e}_3, \mathbf{e}_4$
$\mathcal{G}_7(X_3)$				$\underline{2n}$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_4 D_{nh}, \mathbf{e}_3$
	1	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_3 \mathbf{e}_4$		$n \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_{[12]}, D_{nh}, \mathbf{e}_3$
	$\underline{1}$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_3$		$n : \underline{2}$	$\mathfrak{g}, E, \mathbf{e}_4, \mathbf{e}_3 D_{nh}$ ( $n$ odd), $D_{nh}$ ( $n$ even)
	$1 : \underline{2}$	$\mathfrak{g}, E, \mathbf{e}_3^2, \mathbf{e}_4$		$n : \underline{2}$	$\mathfrak{g}, \mathbf{e}_{[123]}, \mathbf{e}_3 \mathbf{e}_4, \mathbf{e}_3 D_{nh}$ ( $n$ even), $D_{nh}$ ( $n$ odd)
	$1 : 2$	$\mathfrak{g}, \mathbf{e}_{[123]}, \mathbf{e}_{[34]}$		$\underline{2n} : 2$	$\mathfrak{g}, \mathbf{e}_{[123]}, \mathbf{e}_3 \mathbf{e}_4 D_{nh}$ ( $n$ odd), $\mathbf{e}_4 D_{nh}$ ( $n$ even)
	$1 : \underline{2}$	$\mathfrak{g}, \mathbf{e}_3 \mathbf{e}_{[12]}$		$n : \underline{2} \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_{[123]}, \mathbf{e}_3 D_{nh}$ ( $n$ odd), $D_{nh}$ ( $n$ even)
	$\underline{2}$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_4$		$\underline{2n}$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_3 D_{nh}, \mathbf{e}_4$
	2	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_{[34]}, \mathbf{e}_4^2$		$\underline{2n}$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_3 D_{nh}, \mathbf{e}_3 \mathbf{e}_4$
	$\frac{2}{2} \times \underline{1}$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_4^2$		$\underline{2n} \times$	$\mathfrak{g}, \mathbf{e}_{[12]}, \mathbf{e}_3 D_{nh}$
	$1 \cdot m$	$\mathfrak{g}, \mathbf{e}_{[34]}, \mathbf{e}_4$			

(continued)

$n : m$	$g, \mathbf{e}_{[12]}, D_{nh}, \mathbf{e}_4$	$3/2$	$g, \mathbf{e}_{[123]}, T_d, \mathbf{e}_4$
$n : \underline{m}$	$g, \mathbf{e}_{[12]}, D_{nh}, \mathbf{e}_3 \mathbf{e}_4$	$3/2 \times \underline{1}$	$g, \mathbf{e}_{[123]}, T_d$
$\underline{2n} : m$	$g, \mathbf{e}_{[12]}, \mathbf{e}_4 D_{nh}$	$\bar{6}/2$	$g, T_h, \mathbf{e}_4$
$\underline{2n} : \underline{m}$	$g, \mathbf{e}_{[12]}, \mathbf{e}_3 \mathbf{e}_4 D_{nh}$	$\bar{6}/2$	$g, E, T_h$
$n : m \times \underline{1}$	$g, \mathbf{e}_{[12]}, D_{nh}, \mathbf{e}_3^2$	$\bar{6}/2 \times \underline{1}$	$g, T_h$
$n \cdot m$	$g, P_{nh}, \mathbf{e}_3, \mathbf{e}_4$	$3/4$	$g, \mathbf{e}_{[123]}, O_h, \mathbf{e}_4$
$n \cdot \underline{m}$	$g, D_{nh}, \mathbf{e}_4 \mathbf{e}_{[12]}, \mathbf{e}_3$	$3/\bar{4}$	$g, \mathbf{e}_{[123]}, \mathbf{e}_4 T_d$
$\underline{2n} \cdot m$	$g, \mathbf{e}_4 D_{nh}, \mathbf{e}_3$	$3/\bar{4} \times \underline{1}$	$g, \mathbf{e}_{[123]}, O_h$
$n \cdot m \times \underline{1}$	$g, D_{nh}, \mathbf{e}_3, \mathbf{e}^2$ ( $n = 3/\bar{4}$ )	1)	$g, T_d, \mathbf{e}_4$
$m \cdot n : m$	$g, D_{nh}, \mathbf{e}_4, \mathbf{e}_2^2$ ( $n = 3/\bar{4}$ )	1)	$g, E, T_d$
$\underline{m} \cdot n : m$	$g, D_{nh}, \mathbf{e}_4 \mathbf{e}_{[12]}$	$3/\bar{4} \times \underline{1}$	$g, T_d$
$\underline{m} \cdot n : \underline{m}$	$g, E, D_{nh}, \mathbf{e}_4^2, \mathbf{e}_2^2$ ( $n \neq 4$ )	1)	$g, O_h, \mathbf{e}_4$
$m \cdot n : \underline{m}$	$g, D_{nh}, \mathbf{e}_3 \mathbf{e}_4$	$\bar{6}/4$	$g, \mathbf{e}_4 T_h$
$m \cdot \underline{2n} : m$	$g, \mathbf{e}_4 D_{nh}, \mathbf{e}_2^2$ ( $n = \bar{6}/4$ )	1)	$g, E, O_h$

$m \cdot \underline{2n} : \underline{m}$	$g, \mathbf{e}_3 \mathbf{e}_4 D_{nh}$	$\overline{6}/4$	$g, \mathbf{e}_4 T_d$
$m \cdot n : m \times \underline{1}$	$g, D_{nh}, \mathbf{e}_3^2, \mathbf{e}_2^2$	$(n = \overline{6}/4 \times \underline{1})$	$g, O_h$
$\overline{2n} \cdot m$	$g, \mathbf{e}_3 D_{nh}, \mathbf{e}_4$	$3/5$	$g, \mathbf{e}_{[123]}, Y_h, \mathbf{e}_4$
$\overline{2n} \cdot m$	$g, \mathbf{e}_3 D_{nh}, \mathbf{e}_3 \mathbf{e}_4$	$3/5 \times \underline{1}$	$g, \mathbf{e}_{[123]}, Y_h$
$\overline{2n} \cdot \underline{m}$	$g, E, \mathbf{e}_3 D_{nh}, \mathbf{e}_2^2$	$(n = \overline{3}/\overline{10})$	$g, Y_h, \mathbf{e}_4$
		1)	
$\overline{2n} \cdot \underline{m}$	$g, \mathbf{e}_3 D_{nh}, \mathbf{e}_4 \mathbf{e}_{[12]}$	$3/\overline{10}$	$g, \mathbf{e}_4 Y_h$
$\overline{2n} \cdot m \times \underline{1}$	$g, \mathbf{e}_3 D_{nh}, \mathbf{e}_2^2$	$(n = 3/\overline{10} \times \underline{1})$	$g, Y_h$
		1)	

For some tensors special notations have been introduced:  $g = e_1^2 + e_2^2 + e_3^2 - e_4^2$  is the metric tensor specifying the Lorentz group  $\mathcal{L}$ ;  $E = e_{[1234]}$ ;  $O_h = e_1^4 + e_2^4 + e_3^4$ ;  $T_d = e_{(123)}^4$ ;  $T_h = e_1^2 e_2^2 - e_2^2 e_1^2 + e_2^2 e_3^2 - e_3^2 e_2^2 + e_3^2 e_1^2 - e_1^2 e_3^2$ ;  $Y_h = 5e_1^6 + e_1^4 e_2^2 + e_1^2 e_2^4 + 5e_2^6 - e_1^5 e_3 + e_1^3 e_2^2 e_3 + e_1 e_2^4 e_3 + 2e_1^2 e_3^4 + 2e_2^2 e_3^4 + 2e_3^6$ , where the square bracket above denotes summation over all isomers (differing only in the order of the factors);  $D_{nh} = \text{Re}(e_1 + ie_2)^n = e_1^n - e_1^{n-2} e_2^2 + e_1^{n-4} e_2^4 - \dots$ ;  $\omega_n = \text{Re}(e_{[31]} + e_{[41]} + ie_{[32]} + ie_{[42]})^n$ .

Using the known theorems and formulas of the theory of finite-dimensional representations of classical groups (see, for example, <sup>(7)</sup>), one can, by means of the methods of analytic continuation, induction, and the theory of characters of compact groups, derive simple formulas for the dimension  $p_i = p_i(r, \mathcal{G}_i)$  of the space of invariant tensors of rank  $r$  for all connected subgroups  $\mathcal{G}_i$  of the Lorentz group:  $p_1 = (C_r^{r/2}(r/2 + 1))^2$  for even  $r$ , and  $p_1 = 0$  for odd  $r$ ;  $p_2 = p_3 = p_4 = C_{2r}^r/(r + 1)$ ,  $p_5 = (2C_{r-1}^{r/2})^2$  for even  $r$ , and  $p_5 = 0$  for odd  $r$ ;  $p_6 = (C_r^{[r/2]})^2$ ,  $p_7 = p_8 = p_9 = C_{2r}^r$ .

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*Note: Figure translations are in progress. See original paper for figures.*

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