

Soviet-era science, translated into English

# THE CRYSTAL STRUCTURE OF BARIUM TETRAGERMANATE

CRYSTALLOGRAPHY

1968

SovietRxiv

---

View the original and related papers at <https://sovietrxiv.org/items/ru-196801.78610>

Source: Math-Net.Ru and CyberLeninka. Machine translation. Verify with the original.

## Abstract

## Full Text

UDC 548.736

CRYSTALLOGRAPHY

Yu. I. Smolin

# THE CRYSTAL STRUCTURE OF BARIUM TETRAGERMANATE

(Presented by Academician N. V. Belov on 16 XI 1967)

Compounds of the type  $\text{MeGe}_4\text{O}_9$  (where  $\text{Me} = \text{Sr, Pb, Ba}$ ) were first described by Robbins and Levin <sup>(1)</sup>, who, in particular, having determined from Debye-grams the parameters of the hexagonal cells of these compounds and taking into account the multiple relations of the parameters found to the parameters of benitoite  $\text{BaTiSi}_3\text{O}_9$  <sup>(2)</sup>, proposed the existence in the structures under consideration of  $\text{Ge}_3\text{O}_9$  rings constructed from germanium-oxygen tetrahedra and germanium atoms in sixfold coordination.

The present work was undertaken in order to test these assumptions by means of a direct structure determination.

A crystal of  $\text{BaGe}_4\text{O}_9$ , grown by slow cooling of a melt of the same composition, was investigated. From X-ray goniometric scans the crystal was found to belong to the trigonal system (diffraction symmetry  $P\bar{3}m1$ ). The unit-cell parameters, determined from the scans, were refined from Debye-grams and proved to coincide with those given in <sup>(1)</sup>:  $a = 11.61 \pm 0.01 \text{ \AA}$ ;  $c = 4.74 \pm 0.01 \text{ \AA}$ ;  $z = 3$ . Intensity measurements were carried out on a diffractometer with a scintillation counter using monochromatic  $\text{CuK}_\alpha$  radiation ( $hk0$  reflections) and monochromatic  $\text{MoK}_\alpha$  radiation for  $I(hkl)$  with  $l$  from 0 to 5 and up to  $\sin \theta/\lambda = 1.1 \text{ \AA}^{-1}$  in each layer. In all, 848 nonzero, nonequivalent reflections were measured with  $\text{MoK}_\alpha$ . By mechanical treatment the investigated  $\text{BaGe}_4\text{O}_9$  crystal was given the form of a true sphere 0.3 mm in diameter. Corrections for polarization, the kinematic factor, and absorption were taken into account <sup>(3)</sup>.

In measuring the crystal with  $\text{CuK}_\alpha$  radiation, which is close to the absorption edge of the Ba  $L$ -series ( $\Delta f'' \approx 8.6$ ) <sup>(4)</sup>, a sufficiently strong anomalous-scattering effect was observed:  $F(hk0) \neq F(\bar{h}\bar{k}0)$ . The presence of anomalous scattering excludes centrosymmetric groups, and therefore within the observed symmetry only two space groups are possible,  $P321$  and  $P\bar{3}m1$ . The choice between these groups can readily be made using the anomalous-scattering effect <sup>(5,6)</sup>. Thus, for the group  $P321$  the symmetry of the 001 projection is  $p31m$ , whereas for the group  $P\bar{3}m1$  it is  $p\bar{3}m1$ . However, these two plane groups differ in the anomalous-scattering region simply as follows: when  $F(hk0) \neq F(\bar{h}\bar{k}0)$ ,

Fig. 1. Projection of the structure of  $\text{BaGe}_4\text{O}_9$  onto the plane 001. The circles indicate the  $z$ -coordinates of the atoms in hundredths of the period

Figure 1: Fig. 1. Projection of the structure of  $\text{BaGe}_4\text{O}_9$  onto the plane 001. The circles indicate the  $z$ -coordinates of the atoms in hundredths of the period

for  $p31m$  one must have  $F(hk) = F(kh) \neq F(\bar{k}\bar{h})$ , whereas for  $P\bar{3}m1$ ,  $F(hk) = F(\bar{k}\bar{h}) \neq F(kh)$ . Intensity measurements carried out with  $\text{CuK}_\alpha$  showed that in our case  $F(hk) = F(\bar{k}\bar{h})$ , i.e., the symmetry of the projection is  $p\bar{3}1m$ , and thus the space group of symmetry is  $P321$ .

The anomalous-scattering effect was then used to determine the positions of the atoms in the 001 projection. Since among the  $hk0$  reflections there are no reflections with  $h-k \neq 3n$ , a cell three times smaller, with one anomalously scattering atom, could be adopted for the projection. In this case the  $P_s$ -function, constructed from  $\Delta|F(hk0)|^2 = |F(hk0)|^2 - |F(\bar{h}\bar{k}0)|^2$ , should by its positive maxima give the positions of all atoms relative to the anomalously scattering one<sup>(5,7,8)</sup>. In our case it was possible by this method to establish the position of the germanium atoms rela-

with respect to barium. The oxygens were not clearly fixed in the  $P_s$ -function, apparently because of the effect of series termination caused by the need to use  $\text{CuK}_\alpha$  radiation. After the coordinates of the germanium atoms had been determined, there still remained an ambiguity associated with the choice of the origin. In the  $P_s$ -function the anomalously scattering atom is always at the origin, but in the reduced cell  $p3m1$  considered by us barium can occupy a position on any of the threefold axes, which, on transformation to the full cell, leads to nonalternative positions for the barium atom—at the origin and at the point  $1/3, 1/3$ .

**Fig. 1.** Projection of the structure of  $\text{BaGe}_4\text{O}_9$  onto the plane 001. The circles indicate the  $z$ -coordinates of the atoms in hundredths of the period.

The first variant was considered improbable for a number of reasons, including the excessively high local symmetry for the barium atom. The  $z$ -coordinates of the Ge atoms and the oxygen positions were determined from three-dimensional syntheses of the electron density, the initial values of the  $z$ -coordinates being assigned from geometrical analysis. The structure found in this way was refined against all  $F(hkl)$  measurements by the least-squares method, using the Hughes weighting scheme<sup>(9)</sup>, to a value  $R = 0.058$ . In the refinement, atomic factors for neutral atoms were used, with a dispersion correction for barium and germanium. The results of the refinement are given in Table 1.

The projection of the structure onto the plane 001 is shown in Fig. 1, where the  $z$ -coordinates of the atoms in hundredths of the period are also indicated. As is seen from the figure, the structure of barium tetragermanate is built of layers of  $\text{Ge}_3\text{O}_9$  rings, consisting of germanium–oxygen tetrahedra, the oxygens of which form coordination polyhedra of barium and of germanium lying between the

layers of rings and linking these rings to one another. Part of the germanium atoms in the structure, in accordance with the supposition expressed in <sup>(1)</sup>, are located in octahedra formed by the terminal oxygens of the Ge<sub>3</sub>O<sub>9</sub> rings. The Ge–O distances in the octahedron Ge<sub>1</sub>, lying at the origin, are  $1.880 \pm 0.0087$  Å. The octahedron Ge<sub>2</sub> is shown in Fig. 1. The interatomic Ge–O distances in it are:  $1.860 \pm 0.0098$  and  $1.886 \pm 0.0092$  Å. As is seen from the figure, the coordination polyhedron of Ba includes both terminal and

bridging oxygens of the Ge<sub>3</sub>O<sub>9</sub> rings. The coordination number of barium in the structure is 10. The Ba–O distances in this polyhedron are:

$$\begin{array}{llll} \text{Ba} - \text{O}_3 & 2.771 \pm 0.0033 \text{ \AA} & \text{Ba} - \text{O}'_1 & 3.252 \pm 0.0028 \text{ \AA} \\ \text{Ba} - \text{O}'_5 & 3.038 \pm 0.0032 \text{ \AA} & \text{O}_2 & 2.713 \pm 0.0075 \text{ \AA} \\ \text{O}_4 & 2.918 \pm 0.0081 \text{ \AA} & & \end{array}$$

The triple Ge<sub>3</sub>O<sub>9</sub> rings have intrinsic symmetry 2. All distances and angles in the ring are given in Table 2.

**Table 1**

**Atomic coordinates, root-mean-square errors, and temperature factors**

	$x/a$	$\sigma x/a$	$y/b$	$\sigma y/b$	$z/c$	$\sigma z/c$	$B, \text{ \AA}^2$
Ba	0.66895	0.00006	0	–	0	–	0.641
Ge <sub>1</sub>	0	–	0	–	0	–	0.375
Ge <sub>2</sub>	1/3	–	2/3	–	–0.1067	0.0011	0.433
Ge <sub>3</sub>	0	–	0.17881	0.00008	1/2	–	0.369
Ge <sub>4</sub>	0.14966	0.00007	0.48733	0.00006	0.4075	0.0010	0.414
O <sub>1</sub>	0	–	0.47711	0.00070	1/2	–	0.961
O <sub>2</sub>	0.18239	0.00055	0.59497	0.00061	0.1199	0.0029	0.756
O <sub>3</sub>	0.08665	0.00052	0.15153	0.00052	0.7690	0.0027	0.641
O <sub>4</sub>	0.14156	0.00044	0.32376	0.00049	0.3150	0.0025	0.493
O <sub>5</sub>	0.26204	0.00056	0.51218	0.00051	0.6679	0.0027	0.608

When the determination of the structure had been completed, we learned of the work of K. Robbins, A. Perloff, and S. Block, reporting the determination of the structure of BaGe<sub>4</sub>O<sub>9</sub> <sup>(10)</sup>. These authors, noting that the diffraction symmetry of the crystal is  $P\bar{3}m1$ , overlooked the possibility that in this case the space group  $P321$  might occur and, having convinced themselves that it was impossible to describe the structure in the symmetries  $p\bar{3}m1$  and  $P\bar{3}m1$ , adopted the group  $P3$ . Refinement of the correctly found structural motif in a group with lowered

**Table 2**

### Interatomic distances and angles in the Ge<sub>3</sub>O<sub>9</sub> ring

Ge <sub>3</sub> —O <sub>3</sub> 1.749 ± 0.011 Å	O <sub>1</sub> —O <sub>2</sub> 2.589 ± 0.010 Å
Ge <sub>3</sub> —O <sub>4</sub> 1.760 ± 0.0069 Å	O <sub>1</sub> —O <sub>4</sub> 2.815 ± 0.010 Å
O <sub>3</sub> —O <sub>3</sub> ' 3.089 ± 0.021 Å	O <sub>1</sub> —O <sub>5</sub> 2.969 ± 0.0081 Å
O <sub>3</sub> —O <sub>4</sub> 2.852 ± 0.0144 Å	O <sub>2</sub> —O <sub>4</sub> 2.976 ± 0.011 Å
O <sub>3</sub> —O <sub>4</sub> ' 2.752 ± 0.0098 Å	O <sub>2</sub> —O <sub>5</sub> 3.068 ± 0.018 Å
O <sub>4</sub> —O <sub>4</sub> ' 2.847 ± 0.017 Å	O <sub>4</sub> —O <sub>5</sub> 2.610 ± 0.012 Å
O <sub>3</sub> —Ge <sub>3</sub> —O <sub>4</sub> ' 103°17'	O <sub>1</sub> —Ge <sub>4</sub> —O <sub>2</sub> 95°35'
O <sub>3</sub> —Ge <sub>3</sub> —O <sub>3</sub> ' 124°05'	O <sub>1</sub> —Ge <sub>4</sub> —O <sub>4</sub> 106°27'
O <sub>3</sub> —Ge <sub>3</sub> —O <sub>4</sub> 108°45'	O <sub>1</sub> —Ge <sub>4</sub> —O <sub>5</sub> 118°46'
O <sub>4</sub> —Ge <sub>3</sub> —O <sub>4</sub> ' 107°57'	O <sub>2</sub> —Ge <sub>4</sub> —O <sub>4</sub> 114°44'
Ge <sub>4</sub> —O <sub>1</sub> 1.738 ± 0.0037 Å	O <sub>2</sub> —Ge <sub>4</sub> —O <sub>5</sub> 124°17'
Ge <sub>4</sub> —O <sub>2</sub> 1.758 ± 0.012 Å	O <sub>4</sub> —Ge <sub>4</sub> —O <sub>5</sub> 96°52'
Ge <sub>4</sub> —O <sub>4</sub> 1.776 ± 0.0064 Å	Ge <sub>3</sub> —O <sub>4</sub> —Ge <sub>4</sub> 124°45'
Ge <sub>4</sub> —O <sub>5</sub> 1.712 ± 0.010 Å	Ge <sub>4</sub> —O <sub>1</sub> —Ge <sub>4</sub> ' 128°50'

symmetry, although it allowed these authors to obtain an *R*-factor value of 0.068 for 327 reflections, nevertheless led to noticeable inaccuracies in the determination of the atomic coordinates, which in particular caused a large scatter in the values of interatomic distances (1.617–1.895 Å for Ge—O in the tetrahedron). The appearance of negative values of individual temperature factors for some atoms, which the authors explained by failure to take absorption into account, was apparently caused by the same reason.

In conclusion, the author expresses his gratitude to R. G. Grebenschikov, who provided the crystals studied.

Institute of Silicate Chemistry  
named after I. V. Grebenschikov  
Academy of Sciences of the USSR

Received  
1 VIII 1967

### CITED LITERATURE

1. C. R. Robbins, E. M. Levin, *J. Res. Nat. Bur. Stand.*, **65A**, 127 (1961).
2. W. H. Zachariasen, *Zs. Kristallogr.*, **74**, 139 (1930).
3. *International Tables for X-Ray Crystallography*, **2**, Birmingham, 1959.
4. *International Tables for X-Ray Crystallography*, **3**, Birmingham, 1962.

5. *Computing Methods and Phase Problems in X-Ray Crystals Analysis*, 1961.
6. G. N. Ramachandran, *Nature*, **197**, 70 (1963).
7. P. Pepinsky, Y. Okaya, Y. Takeuchi, *Acta crystallogr.*, **10**, 756 (1957).
8. R. Pepinsky, Y. Okaya, *Phys. Rev.*, **108**, 1231 (1957).
9. E. W. Hughes, *J. Am. Chem. Soc.*, **63**, 1737 (1941).
10. C. Robbins, A. Perloff, S. Block, *J. Res. Nat. Bur. Stand.*, **70A**, 385 (1966).

*Note: Figure translations are in progress. See original paper for figures.*

*Source: Math-Net.Ru and CyberLeninka. Machine translation. Verify with the original.*