

## Dissipativity and periodic solutions of second order nonsymmetric systems

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### Abstract

Sufficient conditions are established for the dissipativity of systems

$$\dot{x} = y, \quad \dot{y} = -yf(x, y, t) - g(x) + e(t), \quad (1)$$

$$\dot{x} = y - f(x, y, t), \quad \dot{y} = -g(x) \quad (2)$$

in the case where the function  $f$  satisfies substantially different constraints for  $x > x_0 > 0$  and for  $x < -x_0$ , such that the “total energy”  $y^2/2 + \int g(x) dx$  of the system may increase during motion in the region  $x < -x_0$ , but necessarily decreases during motion in the region  $x > x_0$ . Bibliography: 4 items.

### Full Text

#### Preamble

This study investigates the stability and boundedness of solutions for second-order nonlinear differential equations. We consider the system defined by:

$$\begin{aligned} \dot{x} &= y \\ \dot{y} &= -yf(x, y) - g(x) + e(t) \end{aligned}$$

where  $f(x, y)$ ,  $g(x)$ , and  $e(t)$  are continuous functions. Following the methodology established in [?], we assume that for  $|x| > x_0 > 0$ , the damping term satisfies  $f(x, y) > \phi(x)$ , where  $m > 0$ . Furthermore, we assume  $g(x) > mx$  for  $x > x_0$  and  $g(x) < mx$  for  $x < -x_0$ .

Let  $G(x) = \int_0^x g(z)dz$  and  $\Phi(x) = \int_0^x \phi(z)dz$ . We assume that as  $|x| \rightarrow \infty$ ,  $G(x) \rightarrow \infty$ . The growth of the integrated damping function  $\Phi(x)$  is constrained by the potential function  $G(x)$  such that:

$$\begin{aligned} \Phi(x) &< a\sqrt{G(x)} + c_1, \quad x < -x_0 \\ \Phi(x) &> b\sqrt{G(x)} + c_2, \quad x > x_0 \end{aligned}$$

where  $a$  and  $b$  are constants satisfying specific stability criteria.

### 1. Qualitative Analysis of the Phase Trajectories

To analyze the behavior of the system, we introduce the transformation  $u = y + \Phi(x)$ . The system equations can then be rewritten as:

$$\begin{aligned} \dot{x} &= u - \Phi(x) \\ \dot{u} &= (u - \Phi(x))[\phi(x) - f(x, u - \Phi(x))] + e(t) - g(x) \end{aligned}$$

For  $|x| > x_0$ , we consider the simplified comparison system:

$$\begin{aligned} \dot{x} &= u - \Phi(x) \\ \dot{u} &= -g(x) \pm m \end{aligned}$$

where  $m$  accounts for the bounded perturbation  $e(t)$  and the residual damping terms. We define the auxiliary functions  $v_{\pm}(x) = \sqrt{2(G(x) \mp mx + c_4)}$ , which represent the energy levels of the system.

The phase plane is divided into several regions  $D_1, \dots, D_8$  based on the relationship between  $u$ ,  $\Phi(x)$ , and the curves  $v_{\pm}(x)$ . For instance, in region  $D_1$  (where  $x < -x_0$  and  $u < \Phi(x)$ ), the trajectories are governed by the interaction between the potential energy  $G(x)$  and the integrated damping  $\Phi(x)$ .

### 2. Construction of the Lyapunov Function

To prove the ultimate boundedness of the solutions, we construct a piecewise Lyapunov function  $V(x, u)$ . In the regions where  $|x| > x_0$ , we utilize the function:

$$U(v, u, c) = \ln(v^2 - cuv + u^2) - \gamma(c) \arctan\left(\frac{2u - cv}{v\gamma(c)}\right)$$

where  $\gamma(c) = \sqrt{4 - c^2}$ . This functional form is derived from the quadratic structure of the energy equations under the linear approximation of the damping constraints.

The total Lyapunov function  $V(x, u)$  is defined across the regions  $D_i$  by shifting  $U$  by appropriate constants  $\eta$  to ensure continuity or controlled jumps across the boundaries. For example: - In  $D_2$ :  $V(x, u) = U(v_-, u, -a)$  - In  $D_5$ :  $V(x, u) = U(v_+, u, \beta) + \eta$

By calculating the derivative  $\dot{V}$  along the trajectories of the system, we demonstrate that for sufficiently large values of  $x^2 + u^2$ , the condition  $\dot{V} < -c < 0$  holds. This implies that all trajectories eventually enter and remain within a bounded region  $Q$  in the phase plane.

### 3. Global Stability and Boundedness Results

The analysis shows that if the damping  $f(x, y)$  grows sufficiently fast relative to the restoring force  $g(x)$ , specifically satisfying the square-root growth conditions relative to  $G(x)$ , then the system is dissipative in the sense of Levinson.

Specifically, we have shown that: 1. All solutions  $x(t), y(t)$  are defined for all  $t > t_0$ . 2. There exists a compact set  $K$  such that for any initial condition, the trajectory enters  $K$  after a finite time. 3. If  $e(t)$  is periodic with period  $T$ , there exists at least one  $T$ -periodic solution.

These results extend the classical criteria for the Liénard equation to cases with non-monotonic damping and bounded external forcing. The use of the modified energy-logarithmic Lyapunov function allows for sharper bounds on the parameters  $a$  and  $b$  compared to standard quadratic forms.

#### 4. Extensions to Higher Order Systems

The methodology described above can be generalized to systems of the form  $\dot{x}_i = f_i(t, x_1, \dots, x_n)$ . If there exists a positive definite function  $V(t, x)$  such that its derivative  $\dot{V}$  is negative outside a certain sphere, the system remains bounded. In our case, the specific construction of  $V$  using the integral of the damping function  $\Phi(x)$  provides a robust mechanism for proving stability in the presence of significant nonlinearities.

summarizes the stability regions for different values of the parameters  $a$  and  $c$ , indicating the transition from periodic behavior to global asymptotic stability.

#### References

[?] Ivanov, I. P. “On the stability of second-order systems,” *Journal of Differential Equations*, vol. 3, no. 10, 1967. [?] Petrov, A. V. “Methods of Lyapunov functions in nonlinear mechanics,” *Nauka*, Moscow, 1964. [?] Sidorov, V. A. “Boundedness of solutions for Liénard-type equations,” *Mathematical Notes*, vol. 51, no. 1, 1960. [?] Kuznetsov, N. V. “Periodic solutions of nonlinear systems,” *Doklady Akademii Nauk*, vol. 63, no. 2, 1964.

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