

# ON THE CHOICE OF $\zeta$ -FUNCTIONS

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**Abstract**

**Full Text**

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**PHYSICS**

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**ON THE CHOICE OF  $S$ -FUNCTIONS  
IN THE METHOD OF CONTACT TRANS-  
FORMATIONS**

*(Presented by Academician Ya. K. Syrkin, May 17, 1966)*

1. The method of contact transformations (c.t.) was first applied by Van Vleck to the calculation of the multiplet structure of the levels of diatomic molecules <sup>(1)</sup>, and subsequently by Jordahl in calculations of the paramagnetic susceptibility of a number of salts <sup>(2)</sup>. Later this method was used mainly for the analysis of vibrational-rotational spectra of polyatomic molecules <sup>(3)</sup>. In the works of Herman and Shaffer <sup>(4)</sup> and Amat, Nielsen, and Goldsmith <sup>(5, 6)</sup>,  $S$ -functions of the first and second c.t. were found for diagonalizing the matrix of the vibrational-rotational energy of a polyatomic molecule. In the literature,  $S$ -functions are unknown for the subsequent transformations of the energy operator. Nor are there indications of general methods for finding them.

In the present work a general formula is derived for the c.t. of the operator of vibrational-rotational energy, and a method is proposed for finding  $S$ -functions for diagonalizing operators containing terms of the type  $(p^m q^n + q^n p^m)$  with arbitrary  $m$  and  $n$ .

2. Let a Hermitian operator  $H$  be given in the form

$$H = \sum_{n=0}^{\infty} \lambda^n H_n,$$

where  $\lambda$  is a small parameter, and the orthonormal eigenbasis  $H_0(\Psi_0)$  and commutators of the type

$$[H_m, H_n], [H_m, [H_n, H_k]], \dots, m, n, k = 0, \dots, \infty \quad (1)$$

are known.

Owing to the Hermiticity of the operator  $H$ , there exists some unitary transformation  $T$  that diagonalizes the matrix of the operator  $H$  given in the representation  $\Psi_0$ , i.e.

$$\Psi = \Psi_0 T, \quad \Psi^+ = T^+ \Psi_0^+, \quad E = \Psi_0 T H T^+ \Psi_0^+. \quad (2)$$

Thus, the problem of determining the elements of the diagonal matrix  $E$  of the operator  $H$  reduces to transforming it by means of  $T$  into the operator  $THT^+$ , diagonal in the representation  $\Psi_0$ .

In the c.t. method  $T$  is specified in the form of the product

$$T = T_\infty \cdots T_k T_{k-1} \cdots T_2 T_1, \quad (3)$$

and the individual  $T_k$  in the form

$$T_k = \exp(i\lambda^k S_k), \quad T_k^+ = \exp(-i\lambda^k S_k), \quad (4)$$

which automatically ensures the orthonormality of the basis  $\Psi$ . Expanding  $T_k$  in a series in powers of  $\lambda$ , and restricting ourselves to the  $k$ -th term of (3), we obtain

$$T_k = \sum_{n=0}^{\infty} (n!)^{-1} (i\lambda^k S_k)^n, \quad T_k^+ = \sum_{n=0}^{\infty} (n!)^{-1} (-i\lambda^k S_k)^n; \quad (5)$$

$$T^{(k)} = \prod_k T_k = \prod_k \sum_{n=0}^{\infty} (n!)^{-1} (i\lambda^k S_k)^n; \quad (6)$$

$$E^{(k)} = \Psi_0 \left[ \prod_k \sum_{n=0}^{\infty} (n!)^{-1} (i\lambda^k S_k)^n H \prod_k \sum_{n=0}^{\infty} (n!)^{-1} (-i\lambda^k S_k)^n \right] \Psi_0^+. \quad (7)$$

where the expression in square brackets in (7) is the  $k$ -fold transformed operator ( $H^{(k)}$ ). The functions  $S_k$  must be chosen so that the operator  $H^{(k)}$  commutes with  $H_0$ . Then the operators  $H^{(k)}$  and  $H_0$  have a common orthonormal basis  $\Psi_0$ , and  $E^{(k)}$  is a diagonal matrix. Assigning to  $k$  the values  $1, 2, \dots, n, \dots$ , one can successively obtain the corresponding  $E^{(n)}$ , diagonal to within terms of the  $(n+1)$ -st order of smallness.

For  $k=1$ , for  $H^{(1)}$  we obtain the expression

$$\begin{aligned}
 H^{(1)} &= \sum_{n=0}^{\infty} \lambda^n H_n^{(1)} = \sum_{n=0}^{\infty} \lambda^n H_n + \sum_{l=1}^{\infty} (l!)^{-1} (i\lambda)^l \left[ S_1, \underbrace{[\dots [ S_1, \sum_{m=0}^{\infty} \lambda^m H_m] \dots]}_{l \text{ times}} \right] \\
 &= \sum_{n=0}^{\infty} \lambda^n H_n + \sum_{l=1}^{\infty} \sum_{m=0}^{\infty} (l!)^{-1} i^l \lambda^{l+m} K_{lm}^{(1)},
 \end{aligned} \tag{8}$$

where the commutators  $K_{lm}^{(1)} = [S_1, \underbrace{[\dots [ S_1, H_m] \dots]}_{l \text{ times}}]$  are expanded according to the usual rules.

For arbitrary  $k$ ,

$$H^{(k)} = \sum_{n=0}^{\infty} \lambda^n H_n^{(k-1)} + \sum_{l=1}^{\infty} \sum_{m=0}^{\infty} (l!)^{-1} i^l \lambda^{l+k+m} K_{lm}^{(k)}, \tag{9}$$

where  $K_{lm}^{(k)} = [S_k, [\dots [S_k, H_m^{(k-1)}] \dots]]$ . Expressions for the individual  $H_n^{(k)}$  can be obtained by equating terms with identical powers  $\lambda^n$  in the right- and left-hand sides of (9). All  $S$ -functions, up to  $S_k$ , are chosen so that the operator  $H^{(k)}$  is diagonal in the basis  $\Psi_0$ . The terms  $H_n^{(k)}$  with  $n > k$  are nondiagonal in this basis  $\Psi_0$ , but the nondiagonal elements  $H_n^{(k)}$  (with  $n > k$ ) give a correction to the eigenvalues of  $H_0$  only in the  $(n+1)$ -st approximation. Therefore, to compute corrections to the eigenvalues of  $H_0$  in the  $n$ -th approximation it is sufficient to make  $n-1$  transformations, whereas to compute corrections to the eigenvectors in the same  $(n$ -th) approximation it is necessary to make  $n$  transformations of  $H$ .

3. Suppose all  $S$ -functions up to  $S_{k-1}$  have been found and all operators  $H_n^{(k-1)}$  with  $n \leq k$  have been constructed. It follows from equation (9) that the function  $S_k$  must satisfy the condition

$$H_k^{(k)} = H_k^{(k-1)} + iK_{10}^{(k)}, \tag{10}$$

where  $H_k^{(k)}$  in the basis  $\Psi_0$  is a diagonal matrix.

Finding a function  $S_k$  satisfying condition (10) is possible if the commutation laws are known for the operators of which  $H_k^{-1}$  is composed, with  $H_0$ . For the vibrational-rotational energy operator of a diatomic molecule [7], the problem is simplified by the fact that  $H_k^{(k-1)}$  is expressed through combinations of operators of the dimensionless coordinate and momentum of the type  $(p^m q^n + q^n p^m)$ . Using the relation  $[p, q] = -i$ , one can obtain commutation relations for any combination  $(p^m q^n + q^n p^m)$  with  $H_0$ .\*

\* The commutation relations for  $(p^m q^n + q^n p^m)$  used below in the calculations are not given for lack of space.

**Table 1**

$(H_k^{(k-1)})_i$	$(S_k)_i \cdot \frac{\omega}{2}$	$(H_k^{(k)})_i$
$p$	$\frac{1}{2}q$	0
$q$	$-\frac{1}{2}p$	0
$p^2$	$\frac{1}{8}(pq + qp)$	$\frac{1}{2}(p^2 + q^2)$
$pq + qp$	$-\frac{1}{2}p^2$	0
$q^2$	$-\frac{1}{8}(pq + qp)$	$\frac{1}{2}(p^2 + q^2)$
$p^3$	$\frac{1}{4}(p^2q + qp^2) + \frac{1}{3}q^3$	0
$p^2q + qp^2$	$-\frac{1}{3}p^3$	0
$pq^2 + q^2p$	$\frac{1}{3}q^3$	0
$q^3$	$-\frac{1}{4}(pq^2 + q^2p) - \frac{1}{3}p^3$	0
$p^4$	$\frac{5}{32}(p^3q + qp^3) + \frac{3}{32}(pq^3 + q^3p)$	$\frac{3}{8}(p^2 + q^2)^2 + \frac{3}{8}$
$p^3q + qp^3$	$-\frac{1}{4}p^4$	0
$p^2q^2 + q^2p^2$	$-\frac{1}{16}(p^3q + qp^3) + \frac{1}{16}(pq^3 + q^3p)$	$\frac{1}{4}(p^2 + q^2)^2 - \frac{3}{4}$
$pq^3 + q^3p$	$\frac{1}{4}q^4$	0
$q^4$	$-\frac{3}{32}(p^3q + qp^3) - \frac{5}{32}(pq^3 + q^3p)$	$\frac{3}{8}(p^2 + q^2)^2 + \frac{3}{8}$
$p^5$	$\frac{1}{4}(p^4q + qp^4) + \frac{1}{3}(p^2q^3 + q^3p^2) + \frac{4}{15}q^5 + q$	0
$p^4q + qp^4$	$-\frac{1}{5}p^5$	0
$p^3q^2 + q^2p^3$	$\frac{1}{6}(p^2q^3 + q^3p^2) + \frac{2}{15}q^4 - q$	0
$p^2q^3 + q^3p^2$	$-\frac{1}{6}(p^3q^2 + q^2p^3) - \frac{2}{15}p^5 + p$	0

$(H_k^{(k-1)})_i$	$(S_k)_i \cdot \frac{\omega}{2}$	$(H_k^{(k)})_i$
$pq^4 + q^4p$	$\frac{1}{5}q^5$	0
$q^5$	$-\frac{1}{4}(pq^4 + q^4p) -$	0
$p^6$	$\frac{1}{3}(p^3q^2 + q^2p^3) - \frac{4}{15}p^5 - p$ $\frac{11}{64}(p^5q + qp^5) + \frac{5}{64}(pq^5 +$ $q^5p) + \frac{5}{24}(p^3q^3 +$ $q^3p^3) + \frac{15}{16}(pq + qp)$	$\frac{5}{16}(p^2+q^2)^3 + \frac{25}{16}(p^2+q^2)$
$p^5q + qp^5$	$-\frac{1}{6}p^6$	0
$p^4q^2 + q^2p^4$	$-\frac{1}{32}(p^5q + qp^5) +$ $\frac{1}{32}(pq^5 + q^5p) +$ $\frac{1}{12}(p^3q^3 + q^3p^3) -$ $\frac{3}{8}(pq + qp)$	$\frac{1}{8}(p^2+q^2)^3 - \frac{19}{8}(p^2+q^2)$
$p^3q^3 + q^3p^3$	$-\frac{1}{8}(p^2q^4 + q^4p^2) +$ $\frac{1}{12}q^6 + \frac{3}{2}p^2$	0
$p^2q^4 + q^4p^2$	$-\frac{1}{32}(p^5q + qp^5) +$ $\frac{1}{32}(pq^5 + q^5p) -$ $\frac{1}{12}(p^3q^3 + q^3p^3) +$ $\frac{3}{8}(pq + qp)$	$\frac{1}{8}(p^2+q^2)^3 - \frac{19}{8}(p^2+q^2)$
$pq^5 + q^5p$	$\frac{1}{6}q^6$	0
$q^6$	$-\frac{5}{64}(p^5q + qp^5) -$ $\frac{11}{64}(pq^5 + q^5p) -$ $\frac{5}{24}(p^3q^3 + q^3p^3) -$ $\frac{15}{16}(pq + qp)$	$\frac{5}{16}(p^2+q^2)^3 + \frac{25}{16}(p^2+q^2)$

To choose the function  $S_k$  that diagonalizes  $H_k^{(k-1)}$ , we shall use the following device.

Let us specify  $S^k$  in the form of a linear combination of the operators  $(p^l q^j + q^j p^l)$

$$S_k = \sum_{lj} y_{lj} (p^l q^j + q^j p^l), \quad l + j \leq m + n, \quad (11)$$

and  $H_k^{(k)}$  in the form of the series

$$H_k^{(k)} = \sum_i x_i (p^2 + q^2)^i, \quad i \leq (m + n)/2, \quad (12)$$

where the operator  $(p^2 + q^2)$ , and hence all its powers, are diagonal in the vibrational quantum number  $v$ . Substituting (11) and (12) into (10) and equating the coefficients of identical combinations  $(p^m q^n + q^n p^m)$  on the right- and left-hand sides of (10), one can find all  $y_{lj}$  and  $x_i$ , and consequently the function  $S_k$  and, at the same time, the diagonal correction  $H_k^{(k)}$ .

If  $H_k^{(k-1)}$  contains many terms, they can be diagonalized separately: the complete function  $S_k$  will be equal to the sum of the  $(S_k)_i$  that diagonalize the separate terms  $(H_k^{(k-1)})_i$  in  $H_k^{(k-1)}$ , and the diagonal correction  $H_k^{(k)}$  will be the sum of the diagonal corrections  $(H_k^{(k)})_i$ .

As an example, let us find the function  $(S_k)_i$  for diagonalizing  $(H_k^{(k-1)})_i = c_i q^6$ , where  $c_i$  is a constant coefficient. We represent  $(S_k)_i$  and  $(H_k^{(k)})_i$  in the following form

$$\begin{aligned} (S_k)_i &= y_{51} (p^5 q + q p^5) + y_{15} (p q^5 + q^5 p) + y_{33} (p^3 q^3 + q^3 p^3) + y_{11} (p q + q p), \\ (H_k^{(k)})_i &= x_1 (p^2 + q^2) + x_3 (p^2 + q^2)^3. \end{aligned} \quad (13)$$

The remaining coefficients  $y_{lj}$  and  $x_i$  are equal to zero.\*

Substituting (13) and  $(H_k^{(k-1)})_i = c_i q^6$  into (10) and taking into account that

$$[(p^5 q + q p^5), (p^2 + q^2)] = 4ip^6 - 10i(p^4 q^2 + q^2 p^4) - 60ip^2,$$

$$[(p q^5 + q^5 p), (p^2 + q^2)] = -4iq^6 + 10i(p^2 q^4 + q^4 p^2) + 60iq^2,$$

$$[(p^3 q^3 + q^3 p^3), (p^2 + q^2)] = 6i(p^4 q^2 + q^2 p^4) - 6i(p^2 q^4 + q^4 p^2) + 18i(p^2 - q^2),$$

$$x_3 (p^2 + q^2)^3 = x_3 (p^6 + q^6) + \frac{3}{2} x_3 (p^4 q^2 + q^2 p^4 + p^2 q^4 + q^4 p^2) +$$

$$+4x_3(p^2 + q^2),$$

we obtain the following values for the coefficients  $y_{lj}$  and  $x_i$ :

$$y_{51} = -\frac{5}{64} \frac{2c_i}{\omega}; \quad y_{15} = -\frac{11}{64} \frac{2c_i}{\omega}; \quad y_{33} = -\frac{5}{24} \frac{2c_i}{\omega};$$

$$y_{11} = -\frac{15}{16} \frac{2c_i}{\omega}; \quad x_3 = \frac{5}{16} c_i; \quad x_1 = \frac{25}{16} c_i.$$

In this way the  $S$ -functions given in the table have been found for diagonalizing the operators  $(p^m q^n + q^n p^m)$  for values  $m + n \leq 6$ . The first column of the table contains the various  $(p^m q^n + q^n p^m)$  entering as separate terms in  $H_k^{(k-1)}$ , the second contains the function  $(S_k)_i$ , and the third contains the diagonal correction  $(H_k^{(k)})_i$ . We note that the diagonal correction from  $(p^m q^n + q^n p^m)$  is equal to zero if at least one of the exponents ( $m$  or  $n$ ) is odd. For even  $m$  and  $n$  in (12), only the terms with  $i = (m + n)/2$ ,  $(m + n)/2 - 2$ , etc., are nonzero.

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\* The coefficients  $y_{lj}$  that are different from zero are chosen with the aid of the table of commutators  $[(p^m q^n + q^n p^m), (p^2 + q^2)]$ .

*Note: Figure translations are in progress. See original paper for figures.*

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