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**Abstract**

**Full Text**

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**MATHEMATICS**

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## HARMONIC ANALYSIS ON THE ONE-SHEETED HYPERBOLOID

*(Presented by Academician I. M. Vinogradov on 6 III 1966)*

1. Let a bilinear form be given in the  $n$ -dimensional real space  $R^n$ :  $[x, y] = -x_1y_1 + x_2y_2 + \dots + x_ny_n$ ; let  $G$  be the connected group of linear transformations with determinant 1 that preserve this form.  $G$  acts transitively on the one-sheeted hyperboloid  $X : [x, x] = 1$ . The shifts  $f(x) \rightarrow f(xg)$  form a unitary representation of  $G$  in  $L^2(X)$  with respect to the invariant measure  $dx = |x_1|^{-1}dx_2 \dots dx_n$ .  $X$  is an example of a homogeneous space with a noncompact stationary subgroup. Harmonic analysis on such spaces has been studied only in special cases. An analogue of the Plancherel formula on the one-sheeted hyperboloid for  $n = 4$  ( $G$  is the Lorentz group) was obtained by the horisphere method in <sup>(1)</sup> (see also <sup>(2)</sup>).

In the present work  $L^2(X)$  is decomposed into irreducibles with the aid of spherical functions.

Let  $W$  be the stationary subgroup of the point  $x^0 = (0, 1, 0, \dots, 0)$ ; let  $T_g$  be an irreducible unitary representation of  $G$  in a Hilbert space  $H$ , and let  $\theta$  be a linear functional defined on an everywhere dense set  $D$  and invariant with respect to  $W$ . Then the expression  $(T_g\theta, \theta)$ , understood in the regularized sense (see § 5), defines a function on  $X$ :  $\Phi(x) = (T_g\theta, \theta)$ ,  $x = x^0g$ , which is called spherical. The main problem consists in expanding the  $\delta$ -function on  $X$  in spherical functions.

2. **Representations of  $G$  associated with the cone.** Representations of the group  $G$  were studied in <sup>(3-5)</sup> (see also <sup>(6)</sup>, Ch. X, § 2).

Let  $\tau$  be a complex number;  $D_\tau$  the space of infinitely differentiable functions  $\varphi(\xi)$  on the upper part of the cone  $X_0^+ : [x, x] = 0, x_1 > 0$ , homogeneous of degree  $\sigma = (2-n+\tau)/2$ :  $\varphi(t\xi) = t^\sigma\varphi(\xi)$ ,  $t > 0$ . In  $D_\tau$  there acts a representation  $T^\tau$  of the group  $G$ :  $\varphi(\xi) \rightarrow \varphi(\xi g)$ . A function  $\varphi \in D_\tau$  is determined by its values on the set  $Y : \xi_1 - \xi_2 = 1$ ; a point  $\eta \in Y$  has the form

$$\eta = (\eta_1, \eta_2, \eta) = \left( \frac{1}{2}(\|\eta\|^2 + 1), \frac{1}{2}(\|\eta\|^2 - 1), \eta_3, \dots, \eta_n \right),$$

where  $\|\eta\|^2 = \eta_3^2 + \dots + \eta_n^2$ .  $D_\tau$  is realized as a certain aggregate of functions from  $C^\infty(Y)$ .

The operator

$$A_\tau : \varphi(\eta) \rightarrow \frac{1}{\Gamma(-\tau/2)} \int \|\eta - \eta'\|^{2-n-\tau} \varphi(\eta') d\eta' \quad (d\eta = d\eta_3 \dots d\eta_n)$$

maps  $D_\tau$  into  $D_{-\tau}$  and commutes with the action of  $G$ .

**Continuous series.**  $\tau = i\rho$ ,  $-\infty < \rho < \infty$ .  $T^\tau$  is irreducible and preserves the scalar product

$$(\varphi, \varphi) = \int |\varphi(\eta)|^2 d\eta,$$

$T^\tau$  and  $T^{-\tau}$  are equivalent.

**Discrete series.** For  $\tau = 2 - n - 2l$ ,  $l = 0, 1, 2, \dots$ , the operator  $A_\tau$  has in  $D_\tau$  its own null subspace  $F_\tau$ , which is irreducible with respect to the restriction of  $T^\tau$  (except for  $n = 3$ ) and in which there exists an invariant scalar product

$$B(\varphi, \varphi) = c_{nl} \int \|\eta - \eta'\|^{2l} \nu_n \|\eta - \eta'\| \varphi(\eta) \overline{\varphi(\eta')} d\eta d\eta'.$$

The constant  $c_{nl}$  here is chosen equal to  $(-1)^{l+1} 2\pi^{2-n} \Gamma(l+n-2) / \Gamma(l+1) \Gamma(-\tau/2)$ . The representations in  $F'_\tau$  form a discrete series.

$D_{ip}$  and  $F_{2-n-2l}$  are completed to Hilbert spaces  $H_p$  and  $H_l$  with respect to the scalar products already present.

**3. Functions invariant with respect to the stationary subgroup.** The representation  $T^\tau$  is extended from  $D_\tau$  to the space  $D'_\tau$  of generalized functions  $\theta(\eta)$  on  $X_0^+$ , homogeneous of degree  $\sigma$ ,  $\theta : \varphi(\eta) \rightarrow \int \theta(\eta) A_\tau \varphi(\eta) d\eta$ . For  $\tau = 2 - n - 2l$  denote by  $F'_\tau$  the null subspace of the operator  $A_\tau$ . In  $D'_\tau$  there exist two linearly independent functions invariant with respect to  $W$ , namely  $(\xi_2)_\pm^\sigma$ , if  $\sigma \neq -m$ ,  $m$  natural, and  $\xi_2^{-m}$ ,  $\delta^{(m-1)}(\xi_2)$ , if  $\sigma = -m$ . Only one of the functions  $\xi_2^{2-n-l}$ ,  $\delta^{(n-3+l)}(\xi_2)$  belongs to  $F'_\tau$ : the first for odd  $n$ , the second for even  $n$ . Put  $\theta_l(\eta) = \beta_l(\eta_2)$ , where  $\beta(t) = (-1)^{(n-3)/2} t^{2-n-l}$  for odd  $n$ ,  $\beta(t) = (-1)^t l \Gamma^{-1}(l+n-2) \delta^{(n-3+l)}(t)$  for even  $n$ .

**4. Integral representation of spherical functions.** Let

$$V = \{v \in R^n : [v, v] = 0, v_2 = 1\}, \quad V^\pm = V \cap \{v_1 \gtrless 0\}, \quad \tilde{\eta} = \eta_2 v.$$

**Continuous series.** To the generalized functions  $\theta_p^\pm(\eta) = (\eta_2)_\pm^\sigma$  there correspond (see §1) the spherical functions

$$\Phi_\rho^\pm(x) = \int_Y [x^0, \tilde{\eta}]_\pm^\sigma [x, \tilde{\eta}]_\pm^{\tilde{\sigma}} d\eta = \int_{V^\pm} [x, v]_\pm^{\tilde{\sigma}} dv \quad \left( dv = \frac{dv_3 \dots dv_n}{|v_1|} \right). \quad (1)$$

**Discrete series.** From  $\theta_l(\eta)$  we obtain (see §1) the spherical functions

$$\Psi_l(x) = p_{nl} \int_Y [x_0, \tilde{\eta}]_+^l \beta_l([x, \tilde{\eta}]) d\eta = p_{nl} \int_{V^+} \beta_l([x, v]) dv, \quad (2)$$

where  $p_{nl} = 2^{-\tau} \pi^{(4-n)/2} / \Gamma(l+1)$ . (For  $n = 3$ ,  $\Psi_l$  is the sum of two spherical functions.) The integrals (1), (2) are understood in the sense of the regularized value, see §5.

**5. Regularization of integrals.** Consider the integral

$$R_\lambda(x) = \int_{V^+} [x, v]_+^\lambda dv,$$

where  $\lambda$  is complex,  $x \in X$ . If  $|x_2| > 1$ , then it converges absolutely for values of  $\lambda$  lying in a certain domain, and is analytically continued to the points  $\sigma$  of interest to us from (1). If  $|x_2| < 1$ , then  $R_\lambda(x)$  diverges for all  $\lambda$ . However, then the integral

$$R_{\lambda, \mu}(x) = \int_{V^+} [x, v]_+^\lambda |v_1|^\mu dv$$

converges absolutely for  $\text{Re}(\lambda + \mu) < 3 - n$ ,  $\text{Re} \lambda > -1$ , and is analytically continued to the point  $\lambda = \sigma$ ,  $\mu = 0$ . We proceed analogously with the integrals (2) and (3).

**6. Expression of the spherical functions in terms of Legendre functions.**

**Continuous series.** Let  $\nu = (n-3)/2$ ,  $E(\rho) = \frac{2}{\pi} (2\pi)^\nu |\Gamma(\sigma+1)|^2$ . Then

$$\frac{\Phi_\rho^+(x)}{E(\rho)} = \begin{cases} -e^{-i\nu\pi} \sin \sigma\pi (x_2^2 - 1)^{-\nu/2} Q_{\nu+\sigma}^\nu(x_2), & \text{for } x_2 > 1, x_1 > 0, \\ -e^{-i\nu\pi} \sin \bar{\sigma}\pi (x_2^2 - 1)^{-\nu/2} Q_{-\nu-\bar{\sigma}-1}^\nu(x_2), & \text{for } x_2 > 1, x_1 < 0, \\ \frac{\pi}{2} (1-x_2^2)^{-\nu/2} P_{\nu+\sigma}^\nu(-x_2), & \text{for } |x_2| < 1, \\ 0, & \text{for } x_2 < -1, x_1 > 0, \\ \pi \cos \pi\nu (x_2^2 - 1)^{-\nu/2} P_{\nu+\sigma}^\nu(-x_2), & \text{for } x_2 < -1, x_1 < 0. \end{cases}$$

The expressions for  $\Phi_\rho^-(x)$  are obtained by replacing  $x_1$  by  $-x_1$ .

**Discrete series.** If we denote

$$f(z) = (z^2 - 1)^{-\nu/2} Q_{l+\nu}^\nu(z),$$

then

$$\Psi_l(x) = 2^{\nu-l} \gamma_l \Gamma^{-1}(l+n-2) [f(x_2 + i0) - (-1)^n f(x_2 - i0)].$$

All spherical functions are locally integrable.

**7. Expansion of the  $\delta$ -function in spherical functions.** Let  $\delta(x)$  denote the  $\delta$ -function on  $X$ , concentrated at  $x^0$ . For the expansion of  $\delta(x)$  in spherical functions we use the method of M. Riesz, as in <sup>(1,2)</sup>. Take the generalized function of one variable  $a$

$$s(a, \lambda) = \frac{1}{2^{2\lambda} \Gamma(\lambda + \frac{1}{2})} |a-1|^{2\lambda} a_+^{-\lambda-\nu}, \quad \nu = (n-3)/2,$$

and put

$$S(x, \lambda) = \int_Y s([x, v], \lambda) dv, \quad x \in X. \quad (3)$$

The integral is understood in the sense of the regularization of Sect. 5. It is computed explicitly:

$$S(x, \lambda) = \{\cos \pi \lambda \cdot (1 - x_2)_-^\lambda + (1 - x_2)_+^\lambda\} S_1(x) S_2(\lambda), \quad (4)$$

where  $S_1(x) = (x_2 + 1)^{-\nu}$  for odd  $n$ ;  $S_1(x) = (x_2 + 1)_+^{-\nu}$  for even  $n$ ;

$$S_2(\lambda) = \pi^{-1/2-\nu} 2^{-\lambda+\nu+1} \Gamma(-\nu - \lambda).$$

Let  $K(\lambda, \rho)$  be such that

$$s(a, \lambda) = \int_{-\infty}^{\infty} a_+^{(2-n+i\rho)/2} K(\lambda, \rho) d\rho.$$

Then, denoting  $\Phi_\rho = \Phi_\rho^+ + \Phi_\rho^-$ , we obtain from (1), (2)

$$S(x, \lambda) = \int_{-\infty}^{\infty} K(\lambda, \rho) \overline{\Phi_\rho(x)} d\rho. \quad (5)$$

Put  $\lambda = -\nu - 1$ . Using the results of (7), Ch. III, § 2, item 2, and also the relation  $\Phi_\rho = \Phi_{-\rho}$ , we obtain from (4) and (5)

$$\varepsilon_n \delta(x) = \frac{1}{2} \pi^{-\nu-2} \Gamma(\nu + 1) L(x) + \int_{-\infty}^{\infty} \omega(\rho) \overline{\Phi_\rho(x)} d\rho,$$

where

$$L(x) = (x_2 - 1)^{-1} (1 - x_2^2)^{-\nu} \quad \text{for odd } n; \quad L(x) = (x_2 - 1)^{-1} (1 - x_2^2)_+^{-\nu} \quad \text{for even } n;$$

$\varepsilon_n = \frac{1}{2}$  or 1, respectively, and

$$2^{2n-3} \pi^{n-1} \omega(\rho) = \begin{cases} \rho \tanh \frac{\pi \rho}{2} \cdot \prod_{k=1}^{\nu} [(2k-1)^2 + \rho^2], & n \text{ odd,} \\ \prod_{k=0}^{[\nu]} (4k^2 + \rho^2), & n \text{ even.} \end{cases}$$

We expand  $L(x)$  in  $\Psi_l(x)$ , for which purpose we use the expressions of  $\Psi_l$  in terms of Legendre functions. Finally:

$$\delta(x) = \sum_{l=0}^{\infty} a_l \overline{\Psi_l(x)} + \int_{-\infty}^{\infty} \omega(\rho) \overline{\Phi_\rho(x)} d\rho,$$

where

$$a_l = \frac{1}{2} \pi^{-n/2-1} (2l + n - 2) \Gamma(l + n - 2) 4^{2-n-l}.$$

**8. An analogue of the Plancherel formula.** Let  $C_c^\infty(X)$  be the space of compactly supported infinitely differentiable functions on  $X$ . We shall call the following functions on  $Y$  the Fourier components of a function  $f \in C_c^\infty(X)$ :

$$F_\rho^\pm(\eta) = \int f(x)[x, \eta]_\pm^\sigma dx, \quad \sigma = \frac{2-n+i\rho}{2},$$

$$F_l(\eta) = \int f(x)\beta_l([x, \eta]) dx.$$

For each  $\rho$  or  $l$ , these formulas define a mapping of  $C_c^\infty(X)$  into  $D_\tau$ ,  $\tau = i\rho$  or  $\tau = 2-n-2l$ , respectively. Under this, the shifts by  $g$  pass into  $T_g^\tau$ .

From the expansion of  $\delta(x)$  and the integral representations (1), (2), it follows that  $f \in C_c^\infty(X)$  is reconstructed from its Fourier components as follows:

$$f(x) = \sum_{l=0}^{\infty} a_l B(F_l, \beta_l, ([x, \tilde{\eta}])) + \int_{-\infty}^{\infty} \omega(\rho) \{ (F_\rho^+, [x, \tilde{\eta}]_+^\sigma) + (F_\rho^-, [x, \tilde{\eta}]_-^\sigma) \} d\rho.$$

The correspondence

$$f \rightarrow \{F_\rho^+, F_\rho^-, F_l\}$$

can be extended to

$$L^2(X) \rightarrow \{H_\rho, H_\rho, H_l\}.$$

Thus, the representation in  $L^2(X)$  of the group  $G$  decomposes into representations of the continuous series with multiplicity 2 and representations of the discrete series with multiplicity 1. There is an analogue of the Plancherel formula

$$\int |f(x)|^2 dx = \sum_{l=0}^{\infty} a_l B(F_l, F_l) + \int_{-\infty}^{\infty} \omega(\rho) \{ (F_\rho^+, F_\rho^+) + (F_\rho^-, F_\rho^-) \} d\rho,$$

where  $a_l, \omega_\rho$  are determined by the formulas preceding (7).

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*Note: Figure translations are in progress. See original paper for figures.*

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