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Abstract

Full Text

THEORY OF ELASTICITY

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TRANSFER OF LOAD FROM A STRINGER OF FINITE LENGTH TO AN INFINITE AND A SEMI-INFINITE PLATE

(Presented by Academician Yu. N. Rabotnov on 18 VII 1963)

An infinite and a semi-infinite ($\text{Im } z > 0$, $z = x + iy$) elastic plate are considered, to which, on the segment $0 \leq x \leq a$ of the x -axis, a stringer of length a is riveted, loaded at the end $x = 0$ by a force P directed opposite to the x -axis. The distributed load $q(x)$ transmitted from the stringer to the plate is sought.

Using the known solution of the problem of a concentrated force applied to an infinite plane and to the boundary of a half-plane ⁽¹⁾, we find the deformation of the plate ε_x on the x -axis due to an arbitrary load $q(x)$ directed along the x -axis.

Then, writing the equilibrium equation for a part of the stringer and using the compatibility condition for the deformations of the stringer and the plate, we obtain for both problems the following singular integral equation:

$$\int_0^1 \frac{q(\tau) d\tau}{\tau - \tau_0} = \frac{\lambda\pi^2}{4} \int_{\tau_0}^1 q(\tau) d\tau, \quad (1)$$

whose solution satisfies the equilibrium condition for the entire stringer

$$\frac{\lambda\pi}{2} \int_0^1 q(\tau) d\tau = 1. \quad (2)$$

Here $\tau = x/a$, $\tau_0 = x_0/a$, $0 \leq x$, $x_0 \leq a$.

Moreover, in the case of an infinite plate

$$\lambda = \frac{16}{\pi(3 - \nu)(1 + \nu)} \frac{E\delta}{E_{cFc}} a, \quad q(\tau) = \frac{(3 - \nu)(1 + \nu)E_{cFc}}{8E\delta} \frac{q(x)}{P}, \quad (3)$$

whereas in the case of a semi-infinite plate

$$\lambda = \frac{2}{\pi} \frac{E\delta}{E_c F_c}, \quad q(\tau) = \frac{E_c F_c q(x)}{E\delta P}. \quad (4)$$

In formulas (3) and (4), E , ν , δ are, respectively, the modulus of elasticity, Poisson's ratio, and thickness of the plate; E_c is the modulus of elasticity, and F_c the cross-sectional area of the stringer.

By the methods developed by L. I. Sedov ⁽²⁾, the original equation (1) is reduced to the equivalent Fredholm equation

$$L[q] = A/\sin \theta + B \operatorname{ctg} \theta, \quad (5)$$

where the change of variable has been introduced

$$\tau = \sin^2 \frac{\sigma}{2}, \quad \tau_0 = \sin^2 \frac{\theta}{2}, \quad q(\sigma) = q\left(\sin \frac{\sigma}{2}\right). \quad (6)$$

A is an arbitrary constant determined from condition (2),

$$B = \frac{\lambda}{8} \int_0^\pi q(\sigma) \sigma \sin \sigma \, d\sigma, \quad (7)$$

$$L[q] = q(\theta) - \frac{\lambda}{8} \int_0^\pi \sin \sigma \ln \left| \frac{\sin \frac{\sigma - \theta}{2}}{\sin \frac{\sigma + \theta}{2}} \right| q(\sigma) \, d\sigma. \quad (8)$$

The solution of equation (5) has, at the ends of the interval $\theta = 0; \pi$, a singularity of the type $\operatorname{ctg} \theta$. With

$$q(\theta) = q_1(\theta) + \frac{A \cos 2\theta}{\sin \theta} + B \operatorname{ctg} \theta \quad (9)$$

we transform equation (5) to the form

$$L(q_1) = f(\theta), \quad (10)$$

where

$$f(\theta) = \frac{A}{4}(8 + \lambda) \sin \theta + (A \sin 2\theta + 2B \sin \theta) \frac{\lambda}{8} \ln \operatorname{tg} \frac{\theta}{2}. \quad (11)$$

If

$$q_1(\theta) = f(\theta) + \sum_{n=1}^{\infty} C_n \sin n\theta, \quad (12)$$

then the solution of equation (10) reduces to the solution of the infinite linear system of equations

$$C'_k - \lambda \sum_{n=1}^{\infty} C'_n \alpha_{kn} = b_k \quad (k = 1, 2, 3, \dots). \quad (13)$$

Here

$$\alpha_{kn} = n \{[(k-n)^2 - 1][(k+n)^2 - 1]\}^{-1} \quad (14)$$

for $k+n$ even, and $\alpha_{kn} = 0$ for $k+n$ odd,

$$C_n = -\frac{A\lambda^2 C'_n}{4} \quad (n = 1, 3, 5, \dots); \quad C_n = -\frac{B\lambda^2 C'_n}{4} \quad (n = 2, 4, 6, \dots), \quad (15)$$

$$b_1 = \frac{8+\lambda}{3\lambda} - \frac{1}{12}; \quad b_3 = -\frac{8+\lambda}{45\lambda} - \frac{104}{108 \cdot 24};$$

$$b_k = -\frac{8+\lambda}{\lambda k^2(k^2-4)} + \frac{2(1 + \frac{1}{3} + \dots + \frac{1}{k-4})}{(k^2-1)(k^2-9)} - \frac{3k^2(k+1) - 2(k-1)}{2(k^2-4)(k^2-1)k^2(k+3)} \quad (k = 5, 7, \dots); \quad (16)$$

$$b_2 = -\frac{1}{12}; \quad b_k = \frac{2(1 + \frac{1}{3} + \dots + \frac{1}{k+3})}{k^2(k^2-4)} - \frac{3k+2}{2k^2(k^2-1)(k+2)} \quad (k = 4, 6, \dots). \quad (17)$$

By virtue of (14), system (13) splits into two independent systems for the even and odd coefficients C'_k and is completely regular for any finite values of the parameters λ .

For large λ , $|C_k| \leq 1.756/\lambda$ for odd k , and $|C'_k| \leq 1/\lambda$ for even k .

If system (13) is solved, then the solution of the problem is found from formula (9), taking (12) and (15) into account. For the constants A and B we have the formulas

$$A = \frac{32}{\pi^2 \lambda \chi}, \quad B = AB_1; \quad (18)$$

$$B_1 = \frac{\pi^2 \lambda}{128} \frac{\chi}{1 + \frac{\lambda}{4} - \frac{\lambda^2 \omega}{32} - \frac{\lambda^3}{8} \sum_{2,4,\dots}^{\infty} \frac{C'_n n}{(n^2-1)^2}}; \quad (19)$$

$$\chi = 8 + \lambda - \lambda^2 C'_1 - \frac{\lambda}{3}, \quad \omega = 0.603599 \dots \quad (20)$$

For the stresses in the stringer $P(\theta)$, referred to the stresses $P(0)$ at the point $\theta = 0$, we have the formula

$$\begin{aligned} \frac{P(\theta)}{P(0)} = 1 - \frac{\theta}{\pi} - \frac{1}{\pi \chi} \left\{ B_1(8 + \lambda) \sin \theta + \frac{1}{2}(8 - \chi) \sin 2\theta + \frac{\lambda}{2} \left(\sin \theta - \frac{1}{3} \sin 3\theta - B_1 \sin 2\theta \right) \ln \operatorname{tg} \frac{\theta}{2} \right. \\ \left. + B_1 \lambda \omega_1(\theta) - \lambda^2 \sum_{2,3,4,\dots}^{\infty} C'_n \delta_n \left[\frac{\sin(n-1)\theta}{n-1} - \frac{\sin(n+1)\theta}{n+1} \right] \right\}; \quad (21) \end{aligned}$$

$$\delta_n = 1 \text{ for } n = 3, 5, \dots; \quad \delta_n = B_1 \text{ for } n = 2, 4, 6, \dots;$$

$$\omega_1(\theta) = \int_0^\theta \ln \operatorname{tg} \frac{\theta}{2} d\theta = -2 \sum_{1,3,5}^{\infty} \frac{\sin n\theta}{n^2}. \quad (22)$$

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REFERENCES

- ¹ N. I. Muskhelishvili, *Some Basic Problems of the Mathematical Theory of Elasticity*, 4th ed., Publishing House of the Academy of Sciences of the USSR, 1954. ² L. I. Sedov, *Plane Problems of Hydrodynamics and Aerodynamics*, 1950.

Note: Figure translations are in progress. See original paper for figures.

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