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# MECHANICS

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**Abstract**

**Full Text**

## MECHANICS

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### A NEW SOLUTION OF THE PROBLEM OF THE MOTION IN A NEWTONIAN FORCE FIELD OF A BODY HAVING CAVITIES FILLED WITH FLUID

*(Presented by Academician P. Ya. Kochina on 9 III 1964)*

Assuming that the body with cavities filled with fluid has a fixed point, we write the equations of its motion in the notation of work [1]\*

$$A_1 \frac{d\omega_1}{dt} = (A_2 - A_3)(\omega_2\omega_3 - \varepsilon\gamma_2\gamma_3) + \lambda_2\omega_3 - \lambda_3\omega_2 + e_2\gamma_3 - e_3\gamma_2, \quad (1)$$

$$\frac{d\gamma_1}{dt} = \omega_3\gamma_2 - \omega_2\gamma_3 \quad (1 \ 2 \ 3).$$

The known integrals are

$$\gamma_1^2 + \gamma_2^2 + \gamma_3^2 = \Gamma^2, \quad (2)$$

$$(A_1\omega_1 + \lambda_1)\gamma_1 + (A_2\omega_2 + \lambda_2)\gamma_2 + (A_3\omega_3 + \lambda_3)\gamma_3 = m, \quad (3)$$

$$A_1\omega_1^2 + A_2\omega_2^2 + A_3\omega_3^2 + \varepsilon(A_1\gamma_1^2 + A_2\gamma_2^2 + A_3\gamma_3^2) - 2(e_1\gamma_1 + e_2\gamma_2 + e_3\gamma_3) = 2h. \quad (4)$$

Under the conditions

$$A_1 = A_2 + A_3^{**}, \quad e_1 = 0, \quad \lambda_1 = 0,$$

$$(A_2^2\lambda_2^2 + A_3^2\lambda_3^2)\varepsilon = A_2^2e_2^2 + A_3^2e_3^2$$

the equations (1) have a solution in which

$$\gamma_1 = \frac{(A_2 + A_3)(A_2^2\lambda_2^2 + A_3^2\lambda_3^2)}{(A_2 - A_3)(A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3)}(\omega_1 - s), \quad (5)$$

$$\begin{aligned} \gamma_2 = (A_2^2e_2^2 + A_3^2e_3^2)^{-1} & \left\{ A_2A_3(e_3\lambda_2 + e_2\lambda_3) \left( \omega_3 - \frac{\lambda_3}{A_2} \right) + \right. \\ & \left. + (A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3) \left( \omega_2 - \frac{\lambda_2}{A_3} \right) \right\}, \quad (6) \end{aligned}$$

$$\begin{aligned} \gamma_3 = (A_3^2e_3^2 + A_2^2e_2^2)^{-1} & \left\{ A_3A_2(e_2\lambda_3 + e_3\lambda_2) \left( \omega_2 - \frac{\lambda_2}{A_3} \right) + \right. \\ & \left. + (A_3^2e_3\lambda_3 - A_2^2e_2\lambda_2) \left( \omega_3 - \frac{\lambda_3}{A_2} \right) \right\}. \end{aligned}$$

\* The quantities  $\mu_i, R_i$  of work [1] are denoted here by  $e_i, \gamma_i$ , respectively, and the notation of the constant integral (2) has been changed.

\*\* In article [2] an example is given showing that this condition can be realized in a body with cavities filled with fluid.

Substituting (5), (6)' into the integrals (2), (3), we determine  $\omega_2, \omega_3$  as functions of  $\omega_1$ :

$$\begin{aligned} & \left( \omega_2 - \frac{\lambda_2}{A_3} \right)^2 + \left( \omega_3 - \frac{\lambda_3}{A_2} \right)^2 = \\ & = (A_2^2e_2^2 + A_3^2e_3^2) \left\{ \frac{\Gamma^2}{A_2^2\lambda_2^2 + A_3^2\lambda_3^2} - \frac{(A_2 + A_3)^2(A_2^2\lambda_2^2 + A_3^2\lambda_3^2)}{(A_2 - A_3)^2(A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3)}(\omega_1 - s)^2 \right\}, \\ & (A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3) \left[ A_2 \left( \omega_2 - \frac{\lambda_2}{A_3} \right)^2 - A_3 \left( \omega_3 - \frac{\lambda_3}{A_2} \right)^2 \right] + \quad (7) \\ & + A_2A_3(A_2 + A_3)(e_3\lambda_2 + e_2\lambda_3) \left( \omega_2 - \frac{\lambda_2}{A_3} \right) \left( \omega_3 - \frac{\lambda_3}{A_2} \right) + \\ & + (A_2 + A_3)(A_2^2\lambda_2^2 + A_3^2\lambda_3^2) \left[ \left( \omega_2 - \frac{\lambda_2}{A_3} \right) \frac{e_2}{A_3} + \left( \omega_3 - \frac{\lambda_3}{A_2} \right) \frac{e_3}{A_2} \right] = \\ & = (A_2^2e_2^2 + A_3^2e_3^2) \left\{ m - \frac{(A_2 + A_3)^2(A_2^2\lambda_2^2 + A_3^2\lambda_3^2)}{(A_2 - A_3)(A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3)}\omega_1(\omega_1 - s) \right\}. \end{aligned}$$

Thus, from (6),  $\gamma_2, \gamma_3$  are determined as functions of  $\omega_1$ , after which the first equation (1) determines, by quadrature, the dependence of  $\omega_1$  on  $t$ .

The integral (4) in the indicated solution is dependent; it can be composed from the integrals (5), (6), (7). The constants  $h, m, \Gamma, s$  are connected by the relation

$$\begin{aligned} & (A_2 + A_3)(A_2^2\lambda_2^2 + A_3^2\lambda_3^2)^2 \left[ \frac{A_2\lambda_2^2}{A_3^2} + \frac{A_3\lambda_3^2}{A_2^2} + (A_2 + A_3)s^2 - 2h \right] + \\ & + 2(A_2 - A_3)(A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3)(A_2^2\lambda_2^2 + A_3^2\lambda_3^2)m + \\ & + [4(A_2^2e_2\lambda_2 - A_3^2e_3\lambda_3)^2 + A_2A_3(A_2 + A_3)(e_3\lambda_2 + e_2\lambda_3)^2] A_2A_3\Gamma^2 = 0. \end{aligned}$$

The indicated solution contains 10 independent parameters:

$$A_2, A_3, \frac{e_2}{e_3}, \lambda_2, \lambda_3, \Gamma, s, m, \omega_1^0, \psi_0.$$

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## REFERENCES CITED

1. P. V. Kharlamov, *Journal of Applied Mechanics and Technical Physics*, No. 4, 17 (1963).
2. E. I. Kharlamova, *Doklady AN SSSR*, **125**, No. 5, 996 (1959).

*Note: Figure translations are in progress. See original paper for figures.*

*Source: Math-Net.Ru and CyberLeninka. Machine translation. Verify with the original.*