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Abstract

Full Text

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MATHEMATICAL PHYSICS

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ON CONDITIONS FOR THE ANALYTICITY OF THE SCATTERING MATRIX

(Presented by Academician V. A. Fock on 21 III 1964)

In this paper it is proved that the Jost function $f(k)$ is an entire function of the complex variable k of growth order not exceeding one and such that $f(k) - 1$ is quadratically integrable in k if and only if the potential is finite. Thus it is clarified that the class of potentials considered by Regge ⁽¹⁾ is the maximal admissible one for the function $f(k)$ to be entire of growth order not exceeding one.

Let us consider the Schrödinger equation

$$y'' + [k^2 - p(x)]y = 0 \quad (1)$$

under the assumption on the potential

$$\int_0^{\infty} x|p(x)| dx < \infty. \quad (2)$$

Consider the solution of equation (1) defined by the following asymptotics for large x :

$$\lim_{x \rightarrow \infty} e^{-ikx} f(x, k) = 1. \quad (3)$$

The function $f(k) = f(0, k)$ is called the Jost function. The basic properties of solutions of equation (1) are collected in the book ⁽²⁾. The analytic properties of the function $f(k)$, under the assumption that the potential $p(x)$ is finite, were established in ^(1,3).

Our task is to invert some of the results obtained by Regge. In this direction we shall prove the theorem:

Theorem 1. *Let $f(k)$ be an entire function of growth order not exceeding one, and let $f(k) - 1$ be quadratically integrable. Suppose that $f(k)$ is the Jost function of some scattering problem of the type (1)–(3), M_n being normalization constants.**

Then the potential $p(x)$ is finite.

Proof. We shall base the proof on the Wiener–Paley theorem ⁽⁴⁾ and on the properties of transformation operators ⁽²⁾. It is shown in ⁽²⁾ that

$$f(x, k) = e^{ikx} + \int_x^\infty A(x, y) e^{iky} dy. \quad (4)$$

From (4) we obtain for $f(k)$ the formula

$$f(k) = 1 + \int_0^\infty A(0, y) e^{iky} dy. \quad (5)$$

* Specially chosen; see the note to formula (13).

The function $A(x, y)$ satisfies the equation of V. A. Marchenko ⁽⁵⁾:

$$A(x, y) = F(x + y) + \int_x^\infty A(x, t) F(t + y) dt, \quad x \leq y, \quad (6)$$

where

$$F(t) = \frac{1}{2\pi} \int_{-\infty}^\infty [-S(k) + 1] e^{ikt} dk + \sum_{n=1}^m M_n^2 e^{-k_n t}, \quad (7)$$

$$S(k) = \frac{f(-k)}{f(k)}; \quad (8)$$

the points $k = ik_n$, $k_n > 0$, are poles of the function $f(k)$.

For us the essential relation between the potential $p(x)$ and the function $A(x, y)$ is:

$$p(x) = -2 \frac{d}{dx} A(x, x). \quad (9)$$

All these formulas are contained in the book ⁽²⁾.

From (6), for $x = 0$, we obtain the equation

$$A(y) = F(y) + \int_0^\infty A(t) F(t + y) dt, \quad y \geq 0, \quad (10)$$

where

$$A(y) = A(0, y). \quad (11)$$

By virtue of the Wiener–Paley theorem, from the conditions imposed on $f(k)$ and from equation (5) it follows that $A(y)$ is a finite function, square-integrable.

Transform (10) to the form

$$A(y) = F(y) + \int_y^\infty A(v - y)F(v) dv, \quad y \geq 0. \quad (12)$$

We shall show that from equation (12) it follows that $F(y) \equiv 0$ for sufficiently large y .

Equation (12) is a Volterra equation with a difference kernel. It can be shown (see the Appendix) that there exists a solution of equation (12) having the form*:

$$F(y) = A(y) + \int_y^\infty \Gamma(v - y)A(v) dv. \quad (13)$$

Since the function $A(y)$ is finite, $A(y) = 0$ for $y > y_0$. From (13) it follows that for $y > y_0$, $F(y) = 0$.

To prove the finiteness of the potential we use formula (9) and the following equation, which follows from (6):

$$A(x, x) = F(2x) + \int_x^\infty A(x, t)E(t + x) dt. \quad (14)$$

From equation (14) we conclude that for $2x > y_0$, $A(x, x) = 0$. Hence, from (9), we obtain the finiteness of $p(x)$. Theorem 1 is proved.

Appendix. Here we shall consider the question of the solvability

* With a suitable choice of the constants M_n . If the discrete spectrum is absent, then all $M_n = 0$, and the reasoning can be simplified.

equations of type (12). Let

$$g(x) = f(x) + \int_x^\infty A(y - x)f(y) dy. \quad (15)$$

Let, for some $c > 0$, $e^{cx}g(x) \in L_2(0, \infty)$, $e^{cx}A(x) \in L(0, \infty)$.

Theorem 2. Under the assumptions made, there exists a unique solution of equation (15) such that $e^{c'x}fx \in L_2(0, \infty)$, $0 < c'$. This solution has the form

$$f(x) = g(x) + \int_x^\infty \Gamma(y-x)g(y) dy, \quad (16)$$

where Γ is some function.

Proof. Multiply both sides of equation (15) by e^{bx} , $b > 0$. We obtain the equation

$$g_b(x) = f_b(x) + \int_x^\infty A_b(y-x)f_b(y) dy, \quad (17)$$

where, for example, $g_b(x) \equiv g(x)e^{bx}$.

Introduce the functions:

$$G(x) = \begin{cases} g_b(x), & x \geq 0, \\ 0, & x < 0; \end{cases} \quad (18a)$$

$$E(x) = \begin{cases} f_b(x), & x \geq 0, \\ f_b(x), & x < 0, \end{cases} \quad (18)$$

$$B(x) = \begin{cases} A_b(x), & x \geq 0, \\ 0, & x < 0. \end{cases} \quad (18)$$

Equation (17) can be rewritten in the form

$$G(x) = F(x) + \int_{-\infty}^\infty B(y-x)F(y) dy. \quad (19)$$

Take the Fourier transform of (19). We obtain⁶

$$\tilde{G}(k) = \tilde{F}(k) + \tilde{B}(-k)\tilde{F}(k). \quad (20)$$

Hence

$$\tilde{F}(k) = \frac{\tilde{G}(k)}{1 + \tilde{B}(-k)} = \tilde{G}(k) - \frac{\tilde{B}(-k)\tilde{G}(k)}{1 + \tilde{B}(-k)}. \quad (21)$$

Choosing $b > 0$ sufficiently large, we can ensure that $1 + \tilde{B}(-k) \neq 0$ for $-\infty < k < \infty$. Then, by a well-known theorem of Wiener⁵,

$$\frac{1}{1 + \widetilde{B}(-k)} = \widetilde{\mathcal{L}}(-k),$$

where $\widetilde{\mathcal{L}}(-k)$ is the Fourier transform of some function from $L(0, -\infty)$. Thus,

$$\widetilde{F}(k) = \widetilde{G}(k) + \widetilde{\Gamma}(-k)\widetilde{G}(k), \quad (22)$$

where $\widetilde{\Gamma}(-k) = \widetilde{B}(-k)\widetilde{\mathcal{L}}(-k)$ is the Fourier transform of some function from $L(0, -\infty)$. Inverting (22) and using the convolution theorem, we arrive at the formula:

$$F(x) = G(x) + \int_{-\infty}^{\infty} \Gamma(y-x)G(y) dy. \quad (23)$$

But for $x < 0$, $\Gamma(x) \equiv 0$. This follows from a theorem of Titchmarsh ((⁶, 172), if one takes into account that the function

$$\widetilde{\Gamma}(-k) = \frac{\widetilde{B}(-k)}{1 + \widetilde{B}(-k)}$$

is analytic in the half-plane $\text{Im } k \leq 0$, together with the functions $\widetilde{B}(-k)$ and

$$\frac{1}{1 + \widetilde{B}(-k)}.$$

(Recall that $\widetilde{B}(-k)$ is an entire function of the complex variable k , since $f(k)$ is, by assumption, an entire function.) Therefore (23) has the form (16). The theorem is proved.

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⁶ E. Titchmarsh, *Introduction to the Theory of Fourier Integrals*, Moscow, 1948.

Note: Figure translations are in progress. See original paper for figures.

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