



Soviet-era science, translated into English

MATHEMATICS

V. A. TRENOGIN

1964

SovietRxiv

View the original and related papers at <https://sovietrxiv.org/items/ru-196401.47851>

Source: Math-Net.Ru and CyberLeninka. Machine translation. Verify with the original.

Abstract

Full Text

MATHEMATICS

V. A. TRENOGIN

EXISTENCE AND ASYMPTOTICS OF SOLUTIONS OF THE “SOLITARY WAVE” TYPE FOR DIFFERENTIAL EQUATIONS IN A BANACH SPACE

(Presented by Academician S. L. Sobolev on 20 I 1964)

The problem considered in this note describes a phenomenon well known in mechanics under the name “solitary wave” and first studied by M. A. Lavrent’ev⁽¹⁾ (see also^(2,3)). Recently A. M. Ter-Krikorov and the author showed⁽⁴⁾ that this phenomenon occurs for a certain class of quasilinear elliptic equations in a rectilinear unbounded strip. Below, the validity of analogous results is established for a second-order differential equation in a Banach space. Our investigation makes essential use of the theory of semigroups.

Lemma 1. Let an operator B act in a Banach space E , where:

- 1) B is a closed linear unbounded operator with domain dense in E ;
- 2) $-B$ is the infinitesimal generator of a strongly continuous semigroup $\exp(-\xi B)$, $\xi \geq 0$, such that

$$\|\exp\{-\xi B\}\| \leq \exp\{-\xi\beta\}, \quad \beta = \text{const} > 0.$$

Define $B^{1/2}$ to be a closed linear unbounded operator with domain dense in E , as the inverse of the operator (see⁽⁵⁾)

$$B^{-1/2} = \frac{1}{\sqrt{\pi}} \int_0^{+\infty} \exp\{-\eta B\} \eta^{-1/2} d\eta.$$

Then the operator $-B^{1/2}$ is the infinitesimal generator of the strongly continuous semigroup $\exp\{-\xi B^{1/2}\}$, $\xi \geq 0$, and moreover

$$\|\exp\{-\xi B^{1/2}\}\| \leq \exp\{-\xi\sqrt{\beta}\}.$$

Proof. A related result is available in⁽⁶⁾. Our proof is based on the following integral representation, valid for $\lambda \in (-\sqrt{\beta}, +\infty)$:

$$(B^{1/2} + \lambda I)^{-1} = \frac{1}{\sqrt{\pi}} \int_0^{+\infty} \exp\{-\eta B\} \psi(\lambda\sqrt{\eta}) \eta^{-1/2} d\eta,$$

where

$$\psi(x) = 1 - 2xe^{x^2} \int_x^{+\infty} e^{-t^2} dt > 0.$$

From this representation one derives an estimate for the norm of the operator $(B^{1/2} + \lambda I)^{-1}$, which makes it possible to apply the Hille-Yosida-Phillips theorem (7) and obtain the assertion of the lemma.

We now introduce, for $\delta > 0$, E_δ —the Banach space of abstract even functions $f(\xi)$, continuous on $(-\infty, +\infty)$, with norm

$$|f(\xi)|_\delta = \sup_{(-\infty, +\infty)} \|e^{\delta|\xi|} f(\xi)\|.$$

Lemma 2. Suppose that the conditions of Lemma 1 are fulfilled for the operator B , and suppose $h(\xi) \in E_\delta$, $h'(\xi) \in E_\delta$, where $0 < \delta < \sqrt{\beta}$. Then the boundary-value problem

$$-\frac{d^2 z}{d\xi^2} + Bz = h(\xi), \quad \lim_{\xi \rightarrow \pm\infty} z = 0$$

has a solution $z = z(\xi)$, $z^{(k)}(\xi) \in E_\delta$, $k = 0, 1, 2$. This solution is given by the explicit formula

$$z(\xi) = \frac{1}{2} \int_{-\infty}^{+\infty} \exp\{-|\xi - \eta| B^{1/2}\} B^{-1/2} h(\eta) d\eta.$$

The validity of Lemma 2 is verified by direct computations. We now pass to the problem of a solitary wave; namely, we consider in the Banach space E the following nonlinear boundary-value problem:

$$-\frac{d^2 y}{d\eta^2} + Ay = F(\lambda, y), \quad -\infty < \eta < +\infty, \quad \lim_{\eta \rightarrow \pm\infty} y(\eta) = 0. \quad (1)$$

Here λ is a real parameter; A is a closed linear unbounded operator with dense domain of definition in E ; $F(\lambda, y)$ is a nonlinear operator acting in E , analytic in the Fréchet sense with respect to λ, y in some neighborhood of the point $y = 0$ for all λ , and $F(\lambda, 0) = 0$. Problem (1) always has the trivial solution. Our aim is to give conditions sufficient for a small nontrivial solution to branch off from the trivial solution for some $\lambda = \lambda_0$.

Suppose there exists λ_0 such that the operator $B = A - \partial F(\lambda_0, 0)/\partial y$ satisfies the following conditions:

- 1) Zero is a simple isolated eigenvalue of the operator B , with corresponding null element φ .
- 2) For solvability of the equation $By = h$ it is necessary and sufficient that $\psi(h) = 0$, where ψ is some linear functional in E ; moreover φ and ψ can be normalized so that $\psi(\varphi) = -1$.
- 3) Decompose E into the direct sum $E = E^1 \dot{+} E^{\infty-1}$, where E^1 is the null subspace of the operator B , and $E^{\infty-1}$ is its range. In $E^{\infty-1}$ the operator B is invertible. We require that, for B considered on $E^{\infty-1}$, the conditions of Lemma 1 be fulfilled. Put now $\lambda = \lambda_0 + \varepsilon$ and write $F(\lambda, y)$ in the form

$$F(\lambda_0 + \varepsilon, y) = F_0 y + \sum_{j=2}^{\infty} F_{0j} y^j + \sum_{i=1}^{\infty} \sum_{j=1}^{\infty} F_{ij} \varepsilon^i y^j.$$

The subsequent exposition uses the definition of a generalized Jordan chain, which generalizes the previously given definitions (see ^(4, 8, 9)).

Let the operator $\Phi(y)$ be analytic in the Fréchet sense in some neighborhood of zero, with $\Phi^{(i)}(0) = 0$, $i = 0, 1, \dots, m-1$, but $\Phi^{(m)}(0) \neq 0$. Let $\Phi_i(y_1, y_2, \dots, y_i)$ be the coefficients in the formal expansion

$$\Phi \left(\sum_{k=1}^{\infty} y_k \varepsilon^k \right) = \sum_{i=1}^{\infty} \Phi_i(y_1, y_2, \dots, y_i) \varepsilon^{m+i-1}.$$

Definition. We shall say that the operator B has, relative to the operator $\Phi(y)$, a Jordan chain of length ρ , if there exist ρ linearly independent elements

$\varphi_1, \varphi_2, \dots, \varphi_p$, satisfying the relations $B\varphi_1 = 0$, $B\varphi_k = \Phi_{k-1}(\varphi_1, \dots, \varphi_{k-1})$, $k = 2, \dots, p$, and moreover

$$\psi(\Phi_p(\varphi_1, \dots, \varphi_p)) \neq 0.$$

Theorem 1. Suppose that the operator B has, relative to the operator

$$\Phi(y) = \sum_{j=2}^{\infty} F_{0j} y^j,$$

a Jordan chain of length p , and suppose that $\alpha = \psi(F_{11}\varphi) > 0$. Put $a_p = \psi(\Phi_p(\varphi_1, \dots, \varphi_p))$.

Then, if $p(m-1)$ is odd, for all λ sufficiently close to λ_0 , and if $p(m-1)$ is even, for all λ sufficiently close to λ_0 and satisfying the inequality $a_p(\lambda_0 - \lambda) > 0$,

there exists a nontrivial solution of problem (1) in the space E_δ , where $0 < \delta < \min(\sqrt{\beta}, 2(m-1)\sqrt{\alpha})$. The following asymptotics hold ($\lambda > \lambda_0$, $N > p$):

$$y(\eta, \varepsilon) = \sum_{i=1}^N y_i(\eta\sqrt{\lambda - \lambda_0})(\lambda - \lambda_0)^{[1+(i-1)(m-1)]/p(m-1)} + \\ + O((\lambda - \lambda_0)^{[1+N(m-1)]/p(m-1)}).$$

Remark. Together with $y(\eta)$, $y(\eta + c)$ will be a solution of problem (1) for any c , since (1) is invariant with respect to shifts in η . The asymptotics is constructed in the same way as in (4); its principal term has the form

$$\left[\frac{\alpha(\lambda_0 - \lambda)}{a_p} \right]^{1/p(m-1)} \left[\operatorname{ch} \frac{1}{2} p(m-1) \sqrt{\alpha|\lambda_0 - \lambda|} \eta \right]^{-2/p(m-1)} \varphi.$$

For the proof of the existence of a solution, Lemma 2 is used.

The solitary wave found in Theorem 1 has exponential decay at infinity. We give conditions under which the decay is of power type.

Let L_{ij} be the coefficients of the branching equation composed for finding small solutions of the equation $Ay = F(\lambda_0 + \varepsilon, y)$ (see (10, 11)). Introduce E_* , the Banach space of abstract even functions $f(\xi)$, continuous on $(-\infty, +\infty)$, with norm

$$\|f(\xi)\|_* = \sup_{(-\infty, +\infty)} \|f(\xi)\omega(\xi)\|,$$

where

$$\omega(\xi) = (1 + \xi^2)^{3/2} \ln(e + |\xi|).$$

Theorem 2. Suppose $L_{11} = L_{12} = L_{20} = 0$, $L_{21} \neq 0$, $L_{30} < 0$. Then, for all λ sufficiently close to λ_0 , there exists a nontrivial solution of problem (1) in the space E_* . The asymptotics holds

$$y(\eta, \varepsilon) = \sum_{i=1}^N y_i(\eta(\lambda - \lambda_0))(\lambda - \lambda_0)^i + O((\lambda - \lambda_0)^{N+1}).$$

The scheme of the proof is the same. We note only that in Lemma 2 the space E_δ can be replaced by the space E_* , and lemmas analogous to Lemmas 1 and 2 from (4) can be proved. The principal term of the asymptotics here will be

$$\frac{6(\lambda - \lambda_0)}{L_{21}} \left[\eta^2 (\lambda - \lambda_0)^2 - \frac{9}{2} L_{30} \right]^{-1} \varphi.$$

Example. $E = L_q(\Omega)$, where Ω is a simply connected bounded domain in R^n with boundary Γ ,

$$Ay \equiv - \sum_{i,j=1}^n a_{ij}(\xi) \frac{\partial^2 y}{\partial \xi_i \partial \xi_j} + \sum_{i=1}^n a_i(\xi) \frac{\partial y}{\partial \xi_i} + a(\xi)y,$$

a_{ij}, a_i, a are sufficiently smooth, $\xi = (\xi_1, \dots, \xi_n)$, and

$$\sum_{i,j=1}^n a_{ij}(\xi) \gamma_i \gamma_j \geq k \sum_{i=1}^n \gamma_i^2, \quad k = \text{const} > 0.$$

The domain of definition of A consists of functions belonging to $W_q^{(2)}(\Omega)$ for which $y|_{\Gamma} = 0$. Let

$$F(\lambda, y) = \sum_{i=1}^n F_i(\lambda, \xi) y^i,$$

where $F_i(\lambda, \xi)$ are sufficiently smooth. In the infinite cylinder $\Omega \times (-\infty, +\infty)$ consider the elliptic equation

$$\frac{\partial^2 y}{\partial \eta^2} + \sum_{i,j=1}^n a_{ij}(\xi) \frac{\partial^2 y}{\partial \xi_i \partial \xi_j} - \sum_{i=1}^n a_i(\xi) \frac{\partial y}{\partial \xi_i} - a(\xi)y + F(\lambda, \xi, y) = 0$$

with boundary conditions

$$y|_{\Gamma \times (-\infty, +\infty)} = 0, \quad \lim_{\eta \rightarrow \pm\infty} y = 0.$$

For this problem Theorems 1 and 2 can be rephrased. The conditions of Lemma 1 are verified in ⁽¹²⁾. One can also indicate conditions under which the solution found will be classical.

Moscow Institute of Physics and Technology

Received
15 I 1964

CITED LITERATURE

1. M. A. Lavrent' ev, *36 Trudy Inst. Mat. AN USSR*, 1946.
2. K. Friedrichs, D. Hyerz, *Comm. Pure and Appl. Math.*, **7**, No. 3 (1954).
3. A. M. Ters-Krikorov, *Zhurn. vychisl. matem. i matem. fiz.*, **1**, No. 6 (1961).
4. A. M. Ters-Krikorov, V. A. Trenogin, *Matem. sborn.*, **62** (104), No. 3 (1963).
5. M. A. Krasnosel' skii, P. E. Sobolevskii, *DAN*, **129**, No. 3 (1959).
6. R. S. Phillips, *Pacific J. Math.*, No. 2 (1952).
7. N. Dunford, J. T. Schwartz, *Linear Operators*, IL, 1962.
8. M. I. Vishik, L. A. Lyusternik, *UMN*, **15**, issue 3 (1960).
9. V. A. Trenogin, *DAN*, **140**, No. 2 (1961).
10. M. M. Vainberg, V. A. Trenogin, *UMN*, **17**, issue 2 (1962).
11. M. M. Vainberg, V. A. Trenogin, *UMN*, **18**, issue 5 (1963).
12. P. E. Sobolevskii, *Tr. Mosk. matem. obshch.*, **10** (1961).

Note: Figure translations are in progress. See original paper for figures.

Source: Math-Net.Ru and CyberLeninka. Machine translation. Verify with the original.