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Abstract

Full Text

Mathematics

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ON A BOUNDARY-VALUE PROBLEM

(Presented by Academician I. G. Petrovskii on 29 V 1962)

Consider the differential equation

$$L(u) \equiv (-1)^{[s/2]+1} \frac{\partial^s u}{\partial t^s} + (-1)^{m+1} \sum_{|i|=|j|=m} D^i A^{ij}(x, t) D^j u + B(u) = f(x, t), \quad (1)$$

$$B(u) = \sum_{|i|/2m+s_1/s < 1} B^{(i, s_1)}(x, t) \frac{\partial^{s_1}}{\partial t^{s_1}} D^i u, \quad (2)$$

where $s \geq 1$, $m \geq 1$ are certain integers; $x = (x_1, \dots, x_n) \in \Omega_0$; Ω_0 is a domain bounded by a sufficiently smooth closed surface γ in x -space; $t \in [0, T]$; $(x, t) \in Q \equiv \Omega_0 \times (0, T)$; $\Gamma = \gamma \times [0, T]$. The functions $A^{ij}(x, t) = A^{ji}(x, t)$, $B^{(i, s_1)}(x, t)$ are sufficiently smooth in \bar{Q} ; for simplicity we shall regard them as real-valued; $f(x, t) \in \mathcal{L}_2(Q)$;

$$i = (i_1, \dots, i_n), \quad j = (j_1, \dots, j_n); \quad |i| = i_1 + \dots + i_n; \quad D^i = \partial^{|i|} / \partial x_1^{i_1} \dots \partial x_n^{i_n},$$

$$\sum_{|i|=|j|=m} \xi_1^{i_1} \dots \xi_n^{i_n} A^{ij}(x, t) \xi_1^{j_1} \dots \xi_n^{j_n} \geq \theta^2 > 0 \quad (3)$$

for $|\xi| = 1$, $(x, t) \in \bar{Q}$.

We shall be interested in the question of existence and uniqueness of a solution of equation (1) in the domain Q under the boundary conditions

$$\left. \frac{\partial^r u}{\partial n_{(x,t)}^r} \right|_{\Gamma} = \varphi_r(x, t), \quad 0 \leq r \leq m-1; \quad (4)$$

$$\left. \frac{\partial^r u}{\partial t^r} \right|_{t=0} = \psi_r(x); \quad 0 \leq r \leq k, \quad \text{if } s = 2k+1; \quad (5)$$

$$0 \leq r \leq k-1, \quad \text{if } s = 2k;$$

$$\left. \frac{\partial^r u}{\partial t^r} \right|_{t=T} = \chi_r(x), \quad 0 \leq r \leq k-1, \quad (6)$$

where $n_{(x,t)}$ is the normal to Γ at the point (x, t) ;

$$\varphi_r(x, t) \in W_{t,x,2}^{(2m-r-1/2)/s, 2m-r-1/2}(\Gamma), \quad \psi_r(x) \in W_{x,2}^{2m(1-r/s-1/2s)}(\Omega_0),$$

$$\chi_r(x) \in W_{x,2}^{2m(1-r/s-1/2s)}(\Omega_T), \quad \Omega_T = \bar{Q} \cap (t = T).$$

We shall assume, in addition, that the boundary functions satisfy the natural compatibility conditions.

Let us note that in the case of odd $s = 2k + 1$, instead of problem (1), (4)–(6) one may consider the problem ($\tilde{1}$), ($\tilde{4}$), ($\tilde{5}$), ($\tilde{6}$), where ($\tilde{1}$) differs from (1) by the sign of $\partial^s u / \partial t^s = \partial^{2k+1} u / \partial t^{2k+1}$, and the conditions ($\tilde{5}$) and ($\tilde{6}$) are obtained from (5) and (6) if the latter are interchanged.

Equation (1) for $s = 1$ is parabolic, and the problem considered for it has been well studied. For $s = 3$ the problem (1), (4)–(6) in the case,

when $B(u) \equiv 0$, was studied by A. A. Dezin; in ^(1,2) he constructed a generalized solution of this problem*.

Theorem 1 on an a priori estimate. Let $u(x, t) \in W_{t,x,2}^{(p,q)}(Q)$, $p \geq s$, $q \geq 2m$. Then there exists a constant C , depending only on the coefficients of (1) and on the domain Q , such that

$$\begin{aligned} \|u\|_{W^{(p,q)}(Q)}^2 &\leq C^2 \left(\|u\|_{\mathcal{L}_2(Q)}^2 + \|Lu\|_{W^{(p-s,q-2m)}(Q)}^2 + \right. \\ &\left. + \sum_{r=0}^{m-1} \|\varphi_r\|_{W^{(p_r,q_r)}(\Gamma)}^2 + \sum_{r=0}^k \|\psi_r\|_{W^{q'_r}(\Omega_0)}^2 + \sum_{r=0}^{k-1} \|\chi_r\|_{W^{q'_r}(\Omega_T)}^2 \right), \quad (7) \end{aligned}$$

where $p_r = p(1 - r/q - 1/2q)$, $q_r = q(1 - r/q - 1/2q)$, $q'_r = q(1 - r/p - 1/2p)$.

If $p = s$, $q = 2m$, then estimate (7) can be obtained in the following way: introduce the new variable $w(x, t) = u(x, t) - z(x, t)$, where $z(x, t) \in W^{(s,2m)}(Q)$ and satisfies the boundary conditions (4)–(6) (such a function $z(x, t)$ can be constructed according to ⁽³⁾); $w(x, t)$ satisfies the homogeneous conditions (4)–(6) (we shall call them conditions (4_0)–(6_0)) and equation (1) with right-hand side $F(x, t) = f(x, t) - L(z)$; moreover,

$$\begin{aligned} \|F\|_{\mathcal{L}_2(Q)}^2 &\leq C_1^2 \left(\|f\|_{\mathcal{L}_2(Q)}^2 + \|z\|_{W^{s,2m}(Q)}^2 \right) \leq \\ &\leq C_2^2 \left(\sum \|\varphi_r\|_{W^{(p_r,q_r)}(\Gamma)}^2 + \sum \|\psi_r\|_{W^{q_r}(\Omega_0)}^2 + \|f\|_{\mathcal{L}_2(Q)}^2 + \sum \|\chi_r\|_{W^{q_r}(\Omega_T)}^2 \right). \end{aligned}$$

The estimate of $\|w\|_{W^{(s,2m)}(Q)}^2$ is obtained by multiplying (1) by $\partial^s w / \partial t^s$ and integrating the resulting equality over Q .

In the remaining cases $p \geq s$, $q \geq 2m$, to obtain (7) one should use a partition of unity in \bar{Q} , the “parametrix” of the fundamental solution for interior points of Q , or the “parametrix” of the Green function in a neighborhood of boundary points. With the aid of the “parametrix” of the Green function one can obtain an a priori estimate also under more general boundary conditions.

Theorem 2 on an a priori estimate. Let $u(x, t) \in W^{p,q}(Q)$. Then there exists a constant C , independent of $u(x, t)$, such that

$$\begin{aligned} \|u\|_{W^{(p,q)}(Q)}^2 &\leq C^2 \left(\|u\|_{\mathcal{L}_2(Q)}^2 + \|L(u)\|_{W^{p-s,q-2m}(Q)}^2 + \right. \\ &\left. + \sum \|A_r(x, t, u)\|_{W^{\tilde{p}_r, \tilde{q}_r}(\Gamma)}^2 + \sum \|B_r(x, u)\|_{W^{\tilde{q}_r}(\Omega_0)}^2 + \sum \|C_r(x, u)\|_{W^{\tilde{q}_r}(\Omega_T)}^2 \right), \end{aligned} \quad (8)$$

where

$$\begin{aligned} A_r(x, t, u) &= \sum_{j=0}^{l_r} \alpha_{rj}(x, t) \frac{\partial^j u}{\partial n_{(x,t)}^j}, \quad 0 \leq r \leq m-1; \\ B_r(x, u) &= \sum_{j=0}^{m_r} \beta_{rj}(x) \frac{\partial^j u}{\partial t^j}; \quad 0 \leq r \leq k \quad \text{for } s = 2k+1; \\ &0 \leq r \leq k-1 \quad \text{for } s = 2k; \\ C_r(x, u) &= \sum_{j=0}^{n_r} \gamma_{rj}(x) \frac{\partial^j u}{\partial t^j}, \quad 0 \leq r \leq k-1; \end{aligned}$$

* **Note added in proof.** Recently it has become known to the author that in ⁽¹⁰⁾ the existence of a generalized solution of problem (1), (4)–(6) was proved under the conditions: a) A_{ij} are constants; b) $B(u) \equiv 0$; c) $s < m$; d) either $s = 4s_1$, $m = 2m_1 + 1$, or $s = 4s_1 + 2$, $m = 2m_1$.

$\alpha_{ri}, \beta_{ri}, \gamma_{ri}$ are sufficiently smooth functions given on $\Gamma, \Omega_0, \Omega_T$, respectively; the operators A_r, B_r, C_r cover, in the sense of (4), the operator L ; $\tilde{p}_r = p(1 - l_r/q - 1/2q)$; $\tilde{q}_r = q(1 - l_r/q - 1/2q)$; $\tilde{\tilde{q}}_r = q(1 - m_r/p - 1/2p)$; $\tilde{\tilde{\tilde{q}}}_r = q(1 - n_r/p - 1/2p)$, $p \geq \max(s, m_r, n_r)$, $q \geq \max(2m, l_r)$. In the parabolic case, $s = 1$, Theorem 2 was formulated in [5].

Consider the special case of equation (1)

$$L_0(u) \equiv (-1)^{[s/2]+1} \frac{\partial^s u}{\partial t^s} - (-\Delta)^m u = f(x, t). \quad (9)$$

Theorem 3. The problem (9), (4)–(6) is uniquely solvable in $W^{(s, 2m)}(Q)$ for arbitrary $f(x, t) \in \mathcal{L}_2(Q)$; $\varphi_r(x, t) \in W^{(2m-r-1/2s)2m, (2m-r-1/2)}(\Gamma)$, $\psi_r \in W^{2m(1-r/s-1/2s)}(\Omega_0)$; $\chi_r \in W^{2m(1-r/s-1/2s)}(\Omega_T)$.

To prove Theorem 3, again by means of the substitution $u(x, t) = w(x, t) - z(x, t)$, we reduce the problem (9), (4)–(6) to the problem (9), (4₀)–(6₀) with right-hand side in (9) $F(x, t) = f(x, t) - L_0(z)$. The latter problem is solved as follows: first a generalized solution $w(x, t) \in \overset{\circ}{W}^{[s/2], m}(Q)$ is constructed, satisfying the integral identity

$$\left[\frac{\partial^{[s/2]} w}{\partial t^{[s/2]}}, \frac{\partial^{[s/2]+1} v}{\partial t^{[s/2]+1}} \right] - \sum_{|i|=m} [D^i w, D^i v] = [f, v]. \quad (10)$$

The brackets $[\cdot, \cdot]$ denote integration over Q , and the function $v(x, t)$ is any function from $\overset{\circ}{W}^{[s/2]+1, m}(Q)$ satisfying the boundary conditions (4₀) and

$$v \Big|_{t=0}^{t=T} = \dots = \frac{\partial^{k-1} v}{\partial t^{k-1}} \Big|_{t=0}^{t=T} = \frac{\partial^k v}{\partial t^k} \Big|_{t=0}^{t=T} = 0.$$

Next, uniqueness of the generalized solution in $\overset{\circ}{W}^{[s/2], m}(Q)$ is proved, from which one can obtain the estimate

$$\|w\|_{\overset{\circ}{W}^{[s/2], m}(Q)}^2 \leq C_3^2 \|f\|_{L_2(Q)}^2, \quad (11)$$

where C_3 does not depend on $f(x, t)$. After this it is proved that the generalized solution is in fact a smooth solution from $W^{s, 2m}(Q)$ of equation (9). The proof of the smoothness of the solution inside the domain Q and near those portions of

its boundary which do not include some neighborhood of the set $(\Gamma \cap (t = 0)) \cup (\Gamma \cap (t = T))$ is carried out with the aid of a partition of unity and Hörmander's theorems. (We note that equation (1), as can be proved, is hypoelliptic.)

The proof of smoothness near $t = 0$ ($t = T$) (and hence also near the manifold $(\Gamma \cap (t = 0)) \cup (\Gamma \cap (t = T))$) is carried out in the following way: in (10), as the function $v(x, t)$ one substitutes

$$\rho(t)G_{N,\delta}(x, \xi, t, \tau) = \rho(t) \sum_{r=1}^N V_r(x)V_r(\xi) \int_0^\infty M_r(t, t_1)\xi_\delta(t-t_1) dt_1, \quad \delta > 0, N > l; \quad (12)$$

where $V_r(x)$ is the r -th eigenfunction of the operator $(-\Delta)^m$ in the domain Ω_0 , corresponding to the eigenvalue λ_r , under the boundary conditions $V_r|_\gamma = \dots$

$$\dots = D^{m-1}V_r|_\gamma = 0;$$

$M_r(t, t_1)$ is the Green's function of the operator

$$(-1)^{[s/2]} \frac{\partial^s}{\partial t^s} - \lambda_r,$$

satisfying the conditions

$$M_r(t, t_1)|_{t=0} = \dots = \frac{d^{[s/2]-1}}{dt^{[s/2]-1}} M_r(t, t_1)|_{t=0} = 0, \quad M_r|_{t=\infty} = 0,$$

$$\xi_\delta(t_1 - \tau) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{i\theta(t_1-\tau)} e^{-\delta\theta^2} d\theta, \quad \delta > 0;$$

$\rho(t)$ is a function from $C^\infty([0, T])$; $|\rho(t)| \leq 1$; $\rho(t) = 1$ for $0 \leq t \leq T/3$; $\rho(t) = 0$ for $T/2 \leq t \leq T$.

From equality (10) we obtain

$$w_{N\delta}(\xi, \tau) = [tp, G_{N\delta}] + \sum_{p \leq k} B_p [(u\rho^{2k+1-p})_{tp}, G_{N\delta}] + \sum_{p \geq k+1} B_p \left[(u\rho^{2k+1-p})_{t^k} \frac{\partial^{p-k} G_{N\delta}}{\partial t^{p-k}} \right], \quad (13)$$

where $w_{N\delta}(\xi, \tau) = [w(x, t), L_0^*(x, t)G_{N\delta}(x, \xi, t, \tau)]$ are smooth functions, and B_p , $p = 0, \dots, 2k$, are certain numbers.

It can be proved that for any $w(x, t) \in \mathcal{L}_2(Q)$ the sequence $w_{N\delta}(\xi, \tau)$ as $N \rightarrow \infty$ and $\delta \rightarrow 0$ converges in $\mathcal{L}_2(Q)$ to $w(\xi, \tau)$.

Using (13) and (11), one can prove

$$\left\| \frac{\partial^{[s/2]+1} w_{N\delta}}{\partial \tau^{[s/2]+1}} \right\|_{\mathcal{L}_2(Q)}^2 \leq C_4^2 \|f\|_{\mathcal{L}_2(Q)}^2,$$

whence, in turn, follows the existence of the $([s/2] + 1)$ -st derivative with respect to τ in $\mathcal{L}_2(Q \cap (0 \leq t \leq T/3))$ of $w(\xi, \tau)$. Analogously, by means of a suitably chosen $\tilde{G}_{N\delta}(x, \xi, t, \tau)$, the existence of the $([s/2] + 1)$ -st derivative in $\mathcal{L}_2(Q \cap (2T/3 \leq t \leq T))$ is proved, and hence, by virtue of (6), also in $\mathcal{L}_2(Q)$. Then the existence in $\mathcal{L}_2(Q)$ of the $([s/2] + 2)$ -nd derivative of $w(\xi, \tau)$ with respect to τ , etc., up to the derivative with respect to τ of order s of $w(\xi, \tau)$, is established. In this case the estimate

$$\|w_{t^s}\|_{\mathcal{L}_2(Q)}^2 \leq C_5^2 \|f\|_{\mathcal{L}_2(Q)}^2$$

holds.

From this estimate, equation (9) for $w(x, t)$, and a known theorem from (8), the proof of Theorem 3 follows.

By the method of continuation with respect to a parameter (for the parabolic equation (1) ($s = 1$) this method is presented in (9)), with the aid of Theorem 1 on an a priori estimate, the following theorems are proved.

Theorem 4. For problem (1), (4)–(6) and its adjoint problem (1*), (4*)–(6*), all Fredholm theorems are valid.

In particular, for the existence of a unique solution of (1), (4)–(6) it is necessary and sufficient that problem (1₀), (4₀)–(6₀) have only the trivial solution.

Theorem 5. If for every smooth finite function $v(x, t)$ in Q the inequality

$$[(-1)^m \sum D^i A^{ij}(x, t) D^{jv} - B(v), v] \geq C_6^2 [v, v]$$

holds with a constant C_6 independent of v , then problem (1), (4)–(6) is uniquely solvable in $W^{(s, 2m)}(Q)$ for all right-hand sides and boundary functions under consideration.

Theorem 6. Problem (1), (4)–(6) is uniquely solvable in $W^{s, 2m}(Q)$, if the height T of the cylinder Q is sufficiently small.

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Note: Figure translations are in progress. See original paper for figures.

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