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## Abstract

## Full Text

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## PHYSICAL CHEMISTRY

B. B. DAMASKIN, N. B. GRIGOR' EV

# EFFECT OF POTENTIAL ON THE ATTRACTION INTERACTION BETWEEN ADSORBED ORGANIC MOLECULES

(Presented by Academician A. N. Frumkin, June 4, 1962)

In work <sup>(1)</sup>, one of us considered the influence of the constant  $a$ , characterizing the degree of interaction of adsorbed organic molecules, on the shape of adsorption peaks of the differential capacitance, it being assumed that  $a = \text{const}$ . In reality,  $a$  may vary with the electrode potential both as a result of a change in the area  $S$  per adsorbed molecule as the quantity  $\theta = \Gamma/\Gamma_m$ , i.e., the ratio of adsorption to its limiting value  $\Gamma_m$  <sup>(2)</sup>, increases, and as a result of a change with  $\theta$  in the additional work of adsorption associated with the dipole moment of the adsorbing particles. In the first case, as shown in <sup>(2)</sup>, the relation between the surface charge and  $\theta$  is given by the equation:

$$\varepsilon = \varepsilon_0 [1 - k\theta + (k-1)\theta^2] + C'(\varphi - \varphi_N) [k\theta - (k-1)\theta^2], \quad (1)$$

where  $k = \frac{S_{\theta=0}}{S_{\theta=1}}$ ;  $\varepsilon_0 = \int_0^\varphi C_0 d\varphi$ ;  $C_0$  and  $C'$  are the capacitances, respectively, at  $\theta = 0$  and  $\theta = 1$ ;  $\varphi_N$  is the point of zero charge (p.z.c.) at  $\theta = 1$ ; all potentials  $\varphi$  are measured from the p.z.c. at  $\theta = 0$ . For  $k = 1$ , from (1) we obtain the equation valid for  $a = \text{const}$ :

$$\varepsilon = \varepsilon_0(1 - \theta) + C'(\varphi - \varphi_N)\theta. \quad (1a)$$

If  $k \neq 1$ , then, as shown in <sup>(2)</sup>:

$$a = a_0 + (k-1)P = a_0 + (k-1) \frac{\int_0^\varphi \varepsilon_0 d\varphi + \varphi C'(\varphi_N - \varphi/2)}{A}, \quad (2)$$

Fig. 1. Curves of differential capacitance in 0.9 N NaF (dashed line), and also with additions of tert.-C<sub>5</sub>H<sub>11</sub>OH at concentrations: 1–0.3 M; 2–0.1 M; 3–0.03 M; 4–0.01 M. I—experimental data, 400 Hz; II—calculation by equation (7) at  $A = 1.00$ ,  $B_0 = 25.6$ ,  $C' = 4.4$ ,  $\varphi_N = 0.3$ , and constant  $a = 1.6$ .

Figure 1: Fig. 1. Curves of differential capacitance in 0.9 N NaF (dashed line), and also with additions of tert.-C<sub>5</sub>H<sub>11</sub>OH at concentrations: 1–0.3 M; 2–0.1 M; 3–0.03 M; 4–0.01 M. I—experimental data, 400 Hz; II—calculation by equation (7) at  $A = 1.00$ ,  $B_0 = 25.6$ ,  $C' = 4.4$ ,  $\varphi_N = 0.3$ , and constant  $a = 1.6$ .

where  $A = RT\Gamma_m$ . The dependence of  $a$  on  $\varphi$ , corresponding to equation (2), has the form of a parabola and is realized, for example, in the adsorption on mercury of the cations [(C<sub>4</sub>H<sub>9</sub>)<sub>4</sub>N]<sup>+</sup> (3).

On the other hand, as was noted in (1), in the adsorption on mercury of tert.-C<sub>5</sub>H<sub>11</sub>OH, a linear dependence of  $a$  on  $\varphi$  is observed:

$$a = a_0 + \beta\varphi, \quad (3)$$

where  $a_0$  and  $\beta$  are constants equal, respectively, to 1.64 and 0.25.

Writing the adsorption isotherm in the form

$$Bc = B_0 e^{-Pc} = \frac{\theta}{1-\theta} \exp(-2a_0\theta) \exp(-2\theta\beta\varphi), \quad (4)$$

where  $c$  is the concentration of the organic substance, and using the properties of the total differential in the Gibbs equation, after algebraic transformations we obtain

$$\varepsilon = \varepsilon_0(1-\theta) + C'(\varphi - \varphi_N^0) + \beta A \theta^2. \quad (5)$$

Thus, in both cases, for  $a \neq \text{const}$ , the dependence of the electrode charge on  $\theta$  deviates from the linear dependence corresponding to a constant value of  $a$ . Comparing (1a) and (5) and introducing an additional term into equation (5)

in the quantity  $\varphi_N$ , we obtain

$$\varphi_N = \varphi_N^0 - \frac{\beta A}{C} \theta, \quad (6)$$

where  $\varphi_N^0$  is the value of  $\varphi_N$  at  $\theta = 0$ . In other words, the linear variation of  $a$  with  $\varphi$  observed experimentally can be explained by a change in  $\varphi_N$  with surface coverage. But since the quantity  $\varphi_N$  is associated with the presence, in the

Figure 2

Figure 2: Figure 2

**Fig. 1.** Curves of differential capacitance in 0.9 *N* NaF (dashed line), and also with additions of tert.-C<sub>5</sub>H<sub>11</sub>OH at concentrations: 1–0.3 *M*; 2–0.1 *M*; 3–0.03 *M*; 4–0.01 *M*. *I*—experimental data, 400 Hz; *II*—calculation by equation (7) at  $A = 1.00$ ,  $B_0 = 25.6$ ,  $C' = 4.4$ ,  $\varphi_N = 0.3$ , and constant  $a = 1.6$ .

adsorbing molecules of a dipole moment, the result obtained means a change in the orientation of the adsorbed molecules with increasing coverage, which leads to a change in the component of the dipole moment normal to the surface and to a corresponding change in the additional adsorption energy (2).

We measured  $C, \varphi$ -curves in 0.9 *N* NaF solutions with various additions of tert.-C<sub>5</sub>H<sub>11</sub>OH (Fig. 1 *I*), and from the quadratic dependence obtained for the peak potentials ( $\varphi^{\max}$ ) on  $\lg c$ , the constants used in (1) were refined:  $A = 1.00$ ,  $B_0 = 25.6$ ,  $C' = 4.4$ ,  $\varphi_N = 0.3$ . Under these conditions the maximum variation of  $\varphi_N$ , and consequently also of the additional energy

adsorption in the case of tert.-C<sub>5</sub>H<sub>11</sub>OH is  $\sim 17\%$ . However, as follows from (1), the height of the adsorption peaks should in this case change considerably more strongly. Thus, calculation of the  $C, \varphi$  curves from the equation

$$C = C_0(1 - \theta) + C'\theta + \frac{\theta(1 - \theta)}{1 - 2a\theta(1 - \theta)} \frac{[\varepsilon\delta - C'(\varphi - \varphi_N)]^2}{A}, \quad (7)$$

which follows from Frumkin's theory (2) and is identical to equation (9) in (4), can serve as a sensitive test both of the theory itself and of the existence of a dependence of the quantity  $a$  on the electrode potential.

In Fig. 1 *II* the  $C, \varphi$  curves in 0.9 *N* NaF with various additions of tert.-C<sub>5</sub>H<sub>11</sub>OH are shown, calculated by us with the above constants for constant  $a = 1.6$ , and in Fig. 2—for  $a$  varying linearly with  $\varphi$  according to equation (3). The dependence of  $\theta$  on  $\varphi$  needed in the calculation was determined graphically with the aid of equation (4). Fig. 2 also gives experimental capacitance data obtained in the corresponding solutions at 400 Hz. As can be seen from Fig. 2, the experimental data agree well with the calculated  $C, \varphi$  curves when the linear change of  $a$  with the electrode potential is taken into account. At the same time, it follows from Fig. 1 that the anodic peaks on the  $C, \varphi$  curves calculated for  $a = \text{const}$  are considerably smaller than the experimental ones, whereas the cathodic peaks,

Fig. 2. Curves of differential capacitance in 0.9 *N* NaF (dashed line) and also with additions of tert.-C<sub>5</sub>H<sub>11</sub>OH at concentrations: 1—0.3 *M*; 2—0.1 *M*; 3—0.03 *M*; 4—0.01 *M*. Solid lines—calculation by equation (7) with the same  $A, B_0, C', \varphi_N$  as in Fig. 1;  $a$  was found from equation (3)

for  $a_0 = 1.64$  and  $\beta = 0.25$ . Points—experimental data in the corresponding solutions, 400 Hz.

on the contrary, are somewhat higher than the experimental ones. This result shows that the forces of interaction between tert.-C<sub>5</sub>H<sub>11</sub>OH molecules adsorbed on mercury do not remain constant when the electrode potential changes. Taking into account the linear dependence of  $a$  on  $\varphi$ , Frumkin's theory (2) makes it possible to convey quantitatively the very complex shape of the differential-capacitance curves in the presence of an organic substance.

In work (1) it was shown that the capacitance values at the maximum of the adsorption peak ( $C^{\max}$ ) must be linearly related to the logarithm of the concentration of the organic substance \*:

$$C^{\max} = K_1 + K_2 \lg c, \quad (8)$$

where

$$K_1 = \frac{C_0 + C'}{2} + \frac{C_0 - C'}{2 - a} (a + \ln B_0) + \frac{(C' \varphi_N)^2}{2A} \cdot \frac{a}{(2 - a)};$$

$$K_2 = 2.3 \frac{C_0 - C'}{2 - a}, \quad (9)$$

[Figure 3]

Fig. 3. Dependence of the height of the cathodic capacitance peak in solutions of 0.9 N NaF + tert.-C<sub>5</sub>H<sub>11</sub>OH on the logarithm of the concentration of tert.-C<sub>5</sub>H<sub>11</sub>OH

Figure 3 gives the experimental dependence of  $C^{\max}$  on  $\lg c$ , obtained by us for the cathodic desorption peaks of tert.-C<sub>5</sub>H<sub>11</sub>OH, which satisfies relation (8) very well. A comparison of the experimental values of  $K_1$  and  $K_2$  with those calculated for the constants indicated above and for the mean value  $C_0 = 16.8 \mu\text{F}/\text{cm}^2$  is given in Table 1. As can be seen from the table, the experimental  $K_1$  and  $K_2$  agree considerably better with the calculation at  $a = 1.5$  than at  $a = 1.6$ , corresponding to the potential of maximum adsorption. This result once again indicates a decrease in the attraction constant  $a$  for tert.-C<sub>5</sub>H<sub>11</sub>OH with an increase in the negative value of the potential. It should be noted that the change, established by us, in the forces of interaction between adsorbed tert.-C<sub>5</sub>H<sub>11</sub>OH molecules with the electrode potential is in contradiction with the conclusion of Lorenz and Müller (5).

**Table 1**

	Experiment	$a = 1.6$	$a = 1.5$
$K_1$	135	163	130
$K_2$	54	71	57

Good agreement of the experimental data with the theory <sup>(1,2)</sup> makes it possible to determine the value of  $a$  from experimental differential-capacitance curves by the following methods:

- 1) From the height of the adsorption capacitance peaks, using the equation

$$C^{\max} = \frac{C_0 + C'}{2} + \frac{[\varepsilon_0 - C'(\varphi - \varphi_N)]^2}{A} \frac{1}{4 - 2a}, \quad (10)$$

which follows from equation (7) at  $\theta = 0.5$  (see <sup>(1)</sup>).

- 2) From the width of the capacitance peaks, respectively at 1/2 or 3/4 of their height, and the value  $\Phi \simeq d \ln c / d\varphi^{\max}$ , using equation (12) from work <sup>(1)</sup>.
- 3) From the slope of the straight line  $C^{\max} - \lg c$ , using equations (8) and (9).
- 4) From the shape of the adsorption isotherm by comparing experimental isotherms with those calculated theoretically at different  $a$  (see equation (3) in <sup>(1)</sup>).

All these methods, when  $a$  is variable, prove to be approximate, but the errors thereby obtained usually do not exceed 2-3%.

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Moscow State University  
named after M. V. Lomonosov

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\* An analogous conclusion was drawn in work <sup>(6)</sup> from an analysis of the Langmuir equation.

*Note: Figure translations are in progress. See original paper for figures.*

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