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# MATHEMATICS

1961

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**Abstract**

**Full Text**

## MATHEMATICS

**N. I. AKHIEZER**

### A CONTINUOUS ANALOG OF POLYNOMIALS ORTHOGONAL ON AN ARC OF A CIRCLE

*(Presented by Academician S. N. Bernstein, July 7, 1961)*

1. Orthogonal polynomials on a circle, in contrast to orthogonal trigonometric sums, contain only nonnegative powers

$1, z, z^2, \dots$  of the quantity  $z = e^{it}$ . Accordingly, their continuous analog—a family  $P(x, \lambda)$ , continuous in  $x$  ( $\geq 0$ ), of entire transcendental functions of finite degree in  $\lambda$ —is obtained by means of a “continuous” orthogonalization of the family  $e^{it\lambda}$ , where the parameter  $t$  ranges only over nonnegative real values. The characteristic properties of the family  $P(x, \lambda)$  are as follows: a)  $P(0, \lambda) = 1$ ; b) if  $x > 0$ , then  $P(x, \lambda)$  is an entire function of exact degree  $x$ , bounded in the half-plane  $\text{Im } \lambda \geq 0$ , and such that

$\lim_{|\lambda| \rightarrow \infty} P(x, \lambda)e^{-ix\lambda} = 1$  ( $\text{Im } \lambda \leq 0$ ); c) there exists a nondecreasing function  $\sigma(\lambda)$  ( $-\infty < \lambda < \infty$ ), for which

$$\int_{-\infty}^{\infty} \frac{d\sigma(\lambda)}{1 + \lambda^2} < \infty, \quad (1)$$

and such that for every continuous finite function  $f(x)$  ( $x \geq 0$ )

$$\int_0^{\infty} |f(x)|^2 dx = \int_{-\infty}^{\infty} \left| \int_0^{\infty} f(x)P(x, \lambda) dx \right|^2 d\sigma(\lambda). \quad (2)$$

Property c) expresses the fact that the operator  $U$  in  $\mathcal{L}^2(0, \infty)$ , defined by the formula

$$Uf = \text{l.i.m.}_{n \rightarrow \infty} \int_0^n f(x)P(x, \lambda) dx = F(\lambda), \quad (3)$$

transforms  $\mathcal{L}^2(0, \infty)$  into a certain linear manifold  $\Delta_\sigma \subset \mathcal{L}_\sigma^2$  and is isometric. The inversion of equality (3) has the form

$$f(x) = \text{l.i.m.}_{n \rightarrow \infty} \int_{-n}^n F(\lambda) \overline{P(x, \lambda)} d\sigma(\lambda). \quad (4)$$

The simplest example of a family  $P(x, \lambda)$  is given by the Fourier integral:

$$P(x, \lambda) = e^{ix\lambda} \quad (x \geq 0), \quad \sigma(\lambda) = \frac{1}{2\pi} \lambda \quad (-\infty < \lambda < \infty).$$

M. G. Kreĭn <sup>(1)</sup> showed that to every distribution function  $\sigma(\lambda)$  satisfying relation (1) and some further rather general conditions there corresponds uniquely a family  $P(x, \lambda)$  with precisely the properties just listed. The general method of constructing the function  $P(x, \lambda)$  from the distribution function  $\sigma(\lambda)$  is not very effective and, like the analogous method in the inverse Sturm-Liouville problem, does not make it possible to clarify the influence of empty intervals in the spectrum of the function  $\sigma(\lambda)$ , if any exist.

The present article is devoted to one simple but, it seems to me, important case of this kind. Before turning to it, let us note that, together with  $P(x, \lambda)$ , it is useful to consider the function  $P_*(x, \lambda) = e^{ix\lambda} \overline{P(x, \lambda)}$ .

One of M. G. Krein's results is that the functions  $P(x, \lambda)$ ,  $P_*(x, \lambda)$  are connected by the differential relations

$$\frac{dP(x, \lambda)}{dx} = i\lambda P(x, \lambda) + D(x)P_*(x, \lambda), \quad \frac{dP_*(x, \lambda)}{dx} = \overline{D(x)}P(x, \lambda).$$

- Denote by  $E$  the real axis of the  $\lambda$ -plane with the interval  $(-1, 1)$  removed from it, and take an absolutely continuous distribution function  $\sigma(\lambda)$ , whose derivative is equal to

$$\sigma'(\lambda) = \frac{1}{2\pi} \sqrt{\frac{\lambda+1}{\lambda-1}} \quad (\lambda \in E); \quad \sigma'(\lambda) = 0 \quad (-1 < \lambda < 1).$$

For the construction of the family  $P(x, \lambda)$  we shall apply the method of the papers <sup>(2,3)</sup>, connected with certain considerations on the Riemann surface  $\mathfrak{F}$ , which is obtained from two copies ( $\mathfrak{G}$ ,  $\mathfrak{G}^*$ ) of the  $\lambda$ -plane, cut along the intervals  $(-\infty, -1)$ ,  $(1, \infty)$ , by gluing, under which these intervals become transition lines.

The radicals  $\sqrt{\lambda^2 - 1}$ ,  $\sqrt{\frac{\lambda+1}{\lambda-1}}$  will be regarded as positive on the upper bank of the cut  $(1, \infty)$  of the sheet  $\mathfrak{G}$ . With the aid of the method mentioned we find that in the case under consideration

$$P(x, \lambda) = \frac{1}{2} \left( 1 + \sqrt{\frac{\lambda-1}{\lambda+1}} \right) e^{\frac{1}{2}ix(\lambda + \sqrt{\lambda^2 - 1})} + \frac{1}{2} \left( 1 - \sqrt{\frac{\lambda-1}{\lambda+1}} \right) e^{\frac{1}{2}ix(\lambda - \sqrt{\lambda^2 - 1})},$$

whence it is not difficult to derive that

$$\frac{dP(x, \lambda)}{dx} = i\lambda P(x, \lambda) + \frac{1}{2i} P_*(x, \lambda), \quad \frac{dP_*(x, \lambda)}{dx} = -\frac{1}{2i} P(x, \lambda).$$

Thus, in the case under consideration  $D(x)$  is a constant.

We restrict ourselves to verifying Parseval' s equality (2). Let  $f(x)$  ( $x \geq 0$ ) be an arbitrary continuous finite function, which we may assume to be equal to zero in some neighborhood of the point  $x = 0$ , say for  $0 \leq x \leq \delta$ .

Introduce the usual Fourier transforms

$$\varphi(z) = \int_0^\infty f(x) e^{ixz} dx, \quad \psi(z) = \int_0^\infty \overline{f(x)} e^{-ixz} dx.$$

We can write the equality

$$\begin{aligned} F(\lambda) &= \int_0^\infty f(x) P(x, \lambda) dx = \\ &= \frac{1}{2} \left( 1 + \sqrt{\frac{\lambda-1}{\lambda+1}} \right) \varphi \left( \frac{\lambda + \sqrt{\lambda^2-1}}{2} \right) + \frac{1}{2} \left( 1 - \sqrt{\frac{\lambda-1}{\lambda+1}} \right) \varphi \left( \frac{\lambda - \sqrt{\lambda^2-1}}{2} \right), \end{aligned}$$

and, on the other hand, on the basis of Plancherel' s theorem,

$$\int_0^\infty |f(x)|^2 dx = \frac{1}{2\pi} \int_{-\infty}^\infty |\varphi(t)|^2 dt. \quad (5)$$

Our task now consists in expressing the right-hand side of this formula through  $F(\lambda)$ . To this end note that for  $\lambda \in E$

$$\begin{aligned} |F(\lambda)|^2 &= \frac{1}{2(\lambda+1)} \left[ (\lambda + \sqrt{\lambda^2-1}) \left| \varphi \left( \frac{\lambda + \sqrt{\lambda^2-1}}{2} \right) \right|^2 + \right. \\ &+ (\lambda - \sqrt{\lambda^2-1}) \left| \varphi \left( \frac{\lambda - \sqrt{\lambda^2-1}}{2} \right) \right|^2 \left. + \frac{1}{2(\lambda+1)} \left[ \varphi \left( \frac{\lambda + \sqrt{\lambda^2-1}}{2} \right) \psi \left( \frac{\lambda - \sqrt{\lambda^2-1}}{2} \right) + \right. \right. \\ &\left. \left. + \varphi \left( \frac{\lambda - \sqrt{\lambda^2-1}}{2} \right) \psi \left( \frac{\lambda + \sqrt{\lambda^2-1}}{2} \right) \right] \right]. \end{aligned}$$

We must compute, along the upper bank of  $E$  on the sheet  $\mathfrak{G}$ , the integral

$$\int_E |F(\lambda)|^2 d\sigma(\lambda) = J_1 + J_2.$$

First of all let us show that  $J_2 = 0$ . Indeed,

$$\begin{aligned} J_2 &= \frac{1}{2\pi} \operatorname{Re} \int_E \varphi \left( \frac{\lambda + \sqrt{\lambda^2 - 1}}{2} \right) \psi \left( \frac{\lambda - \sqrt{\lambda^2 - 1}}{2} \right) \frac{d\lambda}{\sqrt{\lambda^2 - 1}} = \\ &= \frac{1}{2\pi} \operatorname{Re} \int_{-\infty}^{\infty} \varphi \left( \frac{\lambda + \sqrt{\lambda^2 - 1}}{2} \right) \psi \left( \frac{\lambda - \sqrt{\lambda^2 - 1}}{2} \right) \frac{d\lambda}{\sqrt{\lambda^2 - 1}}. \end{aligned}$$

Applying the residue theorem, we find that

$$J_2 = -\frac{1}{2\pi} \operatorname{Re} \lim_{N \rightarrow \infty} \int_{K_N} \varphi \left( \frac{\lambda + \sqrt{\lambda^2 - 1}}{2} \right) \psi \left( \frac{\lambda - \sqrt{\lambda^2 - 1}}{2} \right) \frac{d\lambda}{\sqrt{\lambda^2 - 1}},$$

where  $K_N$  is the semicircle  $\lambda = Ne^{i\theta}$  ( $0 \leq \theta \leq \pi$ ). On this semicircle

$$\varphi \left( \frac{\lambda + \sqrt{\lambda^2 - 1}}{2} \right) = e^{-\delta N \sin \theta} O(1), \quad \psi \left( \frac{\lambda - \sqrt{\lambda^2 - 1}}{2} \right) = O(1),$$

and therefore  $J_2 = 0$ . We now turn to the integral  $J_1 = J'_1 + J''_1$ , where

$$\begin{aligned} J'_1 &= \frac{1}{4\pi} \int_E \left| \varphi \left( \frac{\lambda + \sqrt{\lambda^2 - 1}}{2} \right) \right|^2 \left( 1 + \frac{\lambda}{\sqrt{\lambda^2 - 1}} \right) d\lambda, \\ J''_1 &= -\frac{1}{4\pi} \int_E \left| \varphi \left( \frac{\lambda - \sqrt{\lambda^2 - 1}}{2} \right) \right|^2 \left( 1 - \frac{\lambda}{\sqrt{\lambda^2 - 1}} \right) d\lambda. \end{aligned}$$

Obviously,

$$J'_1 = \frac{1}{2\pi} \left( \int_{-\infty}^{-1/2} + \int_{1/2}^{\infty} \right) |\varphi(t)|^2 dt, \quad J''_1 = -\frac{1}{2\pi} \left( \int_0^{-1/2} + \int_{1/2}^0 \right) |\varphi(t)|^2 dt,$$

and, in order to prove Parseval's equality (2), it remains to recall (5).

If there is an empty interval in the spectrum, the operator  $U$ , introduced in No. 1, is unitary, i.e.  $\Delta U = \mathcal{L}_\sigma^2$ . Therefore, from the considerations of the present section there follows

**Theorem 1.** If  $F(\lambda)$  is measurable and

$$\int_E |F(\lambda)|^2 \sqrt{\frac{\lambda+1}{\lambda-1}} d\lambda < \infty,$$

then

$$\begin{aligned} F(\lambda) &= \int_0^\infty f(x) \left\{ \frac{1}{2} \left( 1 + \sqrt{\frac{\lambda-1}{\lambda+1}} \right) e^{\frac{1}{2}ix(\lambda+\sqrt{\lambda^2-1})} + \right. \\ &\quad \left. + \frac{1}{2} \left( 1 - \sqrt{\frac{\lambda-1}{\lambda+1}} \right) e^{\frac{1}{2}ix(\lambda-\sqrt{\lambda^2-1})} \right\} dx, \\ \frac{1}{2\pi} \int_E |F(\lambda)|^2 \sqrt{\frac{\lambda+1}{\lambda-1}} d\lambda &= \int_0^\infty |f(x)|^2 dx. \end{aligned}$$

3. Let us generalize the proposition just proved, following the same path as in (4).

**Lemma.** Every polynomial  $\Omega(\lambda)$ , positive on the set  $E$ , can be represented, and in a unique way, in the form

$$\Omega(\lambda) = \omega(\lambda)\omega^*(\lambda),$$

where

$$\omega(\lambda) = \sum_{k=-n}^n a_k (\lambda + \sqrt{\lambda^2 - 1})^k, \quad \omega^*(\lambda) = \sum_{k=-n}^n \bar{a}_k (\lambda - \sqrt{\lambda^2 - 1})^k$$

and  $a_{-k} = \bar{a}_k$ , while all zeros of  $\omega(\lambda)$  lie on the sheet  $\mathfrak{S}^*$ , and all zeros of  $\omega^*(\lambda)$  on the sheet  $\mathfrak{S}$ ;  $2n$  is the degree of the polynomial  $\Omega(\lambda)$ .

Consider the space  $\mathcal{L}_\tau^2$ , where  $\tau(\lambda)$  is absolutely continuous and

$$\tau'(\lambda) = \frac{1}{\Omega(\lambda)} \sigma'(\lambda),$$

so that

$$\tau'(\lambda) = \frac{1}{2\pi} \frac{1}{\Omega(\lambda)} \sqrt{\frac{\lambda+1}{\lambda-1}} \quad (\lambda \in E); \quad \tau'(\lambda) = 0 \quad (-1 < \lambda < 1),$$

where  $\Omega(\lambda)$  is some polynomial of degree  $2n$ , positive on  $E$ .

Let us find its representation according to the lemma and construct the family of functions

$$Q(x, \lambda) = \frac{1}{2} \left( 1 + \sqrt{\frac{\lambda-1}{\lambda+1}} \right) \omega(\lambda) e^{\frac{1}{2}ix(\lambda + \sqrt{\lambda^2-1})} + \frac{1}{2} \left( 1 - \sqrt{\frac{\lambda-1}{\lambda+1}} \right) \omega^*(\lambda) e^{\frac{1}{2}ix(\lambda - \sqrt{\lambda^2-1})}.$$

The function  $Q(x, \lambda)$ , for every  $x \geq 0$ , is entire with respect to  $\lambda$ , and moreover of degree  $x$ . Let us consider polynomials  $Q_k(\lambda)$  of degree  $k = 0, 1, 2, \dots, n-1$ , orthonormal in the space  $\mathcal{L}_\tau^2$ .

**Theorem 2.** *A function  $F(\lambda)$  ( $\lambda \in E$ ) belongs to  $\mathcal{L}_\tau^2$  if and only if it admits the representation*

$$F(\lambda) = \sum_{k=0}^{n-1} c_k Q_k(\lambda) + \int_0^\infty f(x) Q(x, \lambda) dx, \quad (6)$$

$$f(x) \in \mathcal{L}^2(0, \infty).$$

In this case

$$\int_E |F(\lambda)|^2 d\tau(\lambda) = \sum_{k=0}^{n-1} |c_k|^2 + \int_0^\infty |f(x)|^2 dx.$$

To prove the necessity of the representation (6) for  $F(\lambda)$  to belong to the space  $\mathcal{L}_\tau^2$ , let us find, for the given function  $F(\lambda) \in \mathcal{L}_\tau^2$ , the coefficients

$$c_k = \int_E F(\lambda) Q_k(\lambda) d\tau(\lambda) \quad (k = 0, 1, \dots, n-1)$$

and introduce the function

$$G(\lambda) = \frac{1}{\Omega(\lambda)} \left[ F(\lambda) - \sum_{k=0}^{n-1} c_k Q_k(\lambda) \right], \quad (7)$$

which, evidently, belongs to the space  $\mathcal{L}_\sigma^2$  together with its product by any polynomial of degree  $\leq n$  in  $\lambda$ , and

$$\int_E G(\lambda) \lambda^k d\sigma(\lambda) = 0 \quad (k = 0, 1, \dots, n-1). \quad (8)$$

Representing  $G(\lambda)$  in the form

$$G(\lambda) = \int_0^\infty g(x) P(x, \lambda) dx,$$

let us consider the corresponding  $\mathcal{L}_7^2$  function

$$\Omega(\lambda)G(\lambda) = \omega(\lambda)\omega^*(\lambda) \int_0^\infty g(x)P(x, \lambda) dx. \quad (9)$$

Taking into account the relations (8), as well as the structure of the functions  $\omega(\lambda)$ ,  $\omega^*(\lambda)$ , we find, after integration by parts, that

$$\int_0^\infty g(x)\omega^*(\lambda)e^{\frac{1}{2}ix(\lambda+\sqrt{\lambda^2-1})} dx = \int_0^\infty f(x)e^{\frac{1}{2}ix(\lambda+\sqrt{\lambda^2-1})} dx,$$

$$\int_0^\infty g(x)\omega(\lambda)e^{\frac{1}{2}ix(\lambda-\sqrt{\lambda^2-1})} dx = \int_0^\infty f(x)e^{\frac{1}{2}ix(\lambda-\sqrt{\lambda^2-1})} dx,$$

where on the right-hand sides there is one and the same function  $f(x) \in \mathcal{L}^2(0, \infty)$ .

The equality (9) can now be written in the form

$$\Omega(\lambda)G(\lambda) = \int_0^\infty f(x)Q(x, \lambda) dx,$$

which, together with (7), proves the representation (6).

Physico-Technical Institute of Low Temperatures  
Academy of Sciences of the Ukrainian SSR

Received  
6 VI 1961

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*Note: Figure translations are in progress. See original paper for figures.*

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