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**Abstract**

**Full Text**

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### MATHEMATICS

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#### A SUFFICIENT CONDITION FOR THE SOLVABILITY OF THE MULTIDIMENSIONAL MOMENT PROBLEM

*(Presented by Academician S. L. Sobolev on III 30 1960)*

Let  $\Phi_1$  and  $\Phi_2$  be linear spaces with involution, and let  $A$  and  $B$  be linear operators in  $\Phi_1$  and  $\Phi_2$ , respectively, real with respect to the involution, i.e.  $A\varphi^* = (A\varphi)^*$  for every  $\varphi \in \Phi_1$  and  $B\psi^* = (B\psi)^*$  for every  $\psi \in \Phi_2$ . Let  $\Phi = \Phi_1 \otimes \Phi_2$  be the tensor product of  $\Phi_1$  and  $\Phi_2$ , i.e. the set of elements  $x$  of the form  $\sum_{k=1}^n \varphi_k \otimes \psi_k$ , where  $\varphi_k \in \Phi_1$ ,  $\psi_k \in \Phi_2$ . The operators  $A$  and  $B$  and the involution are defined in  $\Phi$  in the natural way. Suppose that a scalar product  $(x, y)$ , real with respect to the involution, is given in  $\Phi$ , i.e.  $(x^*, y^*) = \overline{(x, y)}$ , and that  $A$  and  $B$  are symmetric operators in this scalar product. As a consequence of the reality of  $A$ , and also of  $B$ , they have equal deficiency indices. Denote by  $H$  the completion of  $\Phi$  with respect to the scalar product  $(x, y)$ . For fixed  $\psi_0 \in \Phi_2$ , obviously,  $(\varphi_1 \otimes \psi_0, \varphi_2 \otimes \psi_0)$  is a scalar product in  $\Phi_1$ .

**Theorem 1.** *Let  $\bar{A}_{\psi_0}$ , the closure of  $A$  in the Hilbert space  $H_{\psi_0}$  obtained by completing  $\Phi_1$  with respect to the scalar product  $(\varphi_1 \otimes \psi_0, \varphi_2 \otimes \psi_0)$ , be a self-adjoint operator for every fixed  $\psi_0 \in \Phi_2$ .*

*Then  $\bar{A}$ , the closure of  $A$  in  $H$ , is a self-adjoint operator, and in  $H$  there exists a self-adjoint extension  $\tilde{B}$  of the operator  $\bar{B}$  such that  $\bar{A}$  and  $\tilde{B}$  commute in  $H$ , i.e. their spectral families commute.*

The case when the closures of  $A$  and  $B$  in  $H$  are self-adjoint operators was considered by A. G. Kostyuchenko and B. S. Mityagin in <sup>(1)</sup>. The case when only  $A$  is self-adjoint was considered by R. S. Ismagilov under the assumption that  $\Phi$  is a Carleman space <sup>(2)</sup>. Theorem 1 generalizes R. S. Ismagilov's result, which makes it possible to obtain certain new results for the multidimensional moment problem and the representation of positive-definite functionals.

**Proof of Theorem 1.** The operators  $A$  and  $B$ , obviously, commute on  $\Phi$ . Therefore, for any non-real  $\lambda$  the equality

$$R_\lambda Bx = BR_\lambda x \quad \text{for } x \in \Phi_\lambda = (A - \lambda E)\Phi_1 \otimes \Phi_2, \quad (1)$$

holds, where  $R_\lambda$  is the resolvent of the operator  $\bar{A}$ . Denote by  $B_\lambda$  the operator  $B$  considered on the linear manifold  $\Phi_\lambda$ . Equality (1) holds for any  $x \in D_{\bar{B}_\lambda}$ , where  $\bar{B}_\lambda$  is the closure of  $B_\lambda$  in  $H$ , since  $R_\lambda$  is a bounded operator. We shall show that  $\bar{B}_\lambda \supset B$ , i.e.  $D_{\bar{B}_\lambda} \supset \Phi$  and for  $x \in \Phi$  one has  $\bar{B}_\lambda x = Bx$ .

Indeed, let  $\varphi_0$  be any element of  $\Phi_1$  and  $\psi_0$  any element of  $\Phi_2$ . From the condition of the theorem it follows that the set  $(A - \lambda E)\Phi_1$  is dense in  $H_{\psi_1}$ , where  $\psi_1 = (B - iE)\psi_0 \in \Phi_2$ , so that there exists a sequence  $\varphi_n^{(1)} = (A - \lambda E)\varphi_n$ , where  $\varphi_n \in \Phi_1$ , converging to  $\varphi_0$  in the norm  $H_{\psi_1}$ . But, since  $B$  is symmetric,  $\|(B - iE)x\|^2 = \|Bx\|^2 + \|x\|^2$ , so that convergence in  $H_{\psi_1}$  means that simultaneously

$$(A - \lambda E)\varphi_n \otimes \varphi_0 \rightarrow \varphi_0 \otimes \psi_0, \quad (A - \lambda E)\varphi_n \otimes B\psi_0 \rightarrow \varphi_0 \otimes B\psi_0$$

in the norm  $H$ . Consequently,  $\varphi_0 \otimes \psi_0 \in D_{\bar{B}_\lambda}$  and, hence,  $\Phi \subset D_{\bar{B}_\lambda}$ . On the other hand, since  $B_\lambda \subset B$ , we have  $\bar{B}_\lambda \subset \bar{B}$ . Thus,  $\bar{B}_\lambda = \bar{B}$ . Hence, for any  $x \in \Phi$  we have  $\bar{B}R_\lambda x = R_\lambda \bar{B}x$ . It follows that for any  $y \in D_{B^*}$ ,

$$(\bar{B}R_\lambda x, y) = (R_\lambda x, B^*y) = \int_{-\infty}^{\infty} \frac{1}{\lambda - \xi} d(E_\xi x, B^*y).$$

Next,

$$(R_\lambda \bar{B}x, y) = \int_{-\infty}^{\infty} \frac{1}{\lambda - \xi} d(E_\xi \bar{B}x, y).$$

Comparing the last two equalities, from the uniqueness of the Stieltjes transform we have, for any  $\xi$ , that  $(E_\xi \bar{B}x, y) = (E_\xi x, B^*y)$ . The equality

$$(\bar{B}x, E_\xi y) = (x, E_\xi B^*y)$$

holds for all  $x \in \Phi$ . It follows that if  $y \in D_{B^*}$ , then also  $E_\xi y \in D_{B^*}$ , and

$$B^*E_\xi y = E_\xi B^*y, \quad (2)$$

i.e., the operator  $B^*$  commutes with  $E_\xi$ .

Further,  $D_{B^*} = D_{\bar{B}} \dot{+} M_+ \dot{+} M_-$ , where  $B^*M_+ = iM_+$  and  $B^*M_- = -iM_-$ . From (2) it follows that  $E_\xi M_\pm \subset M_\pm$  and  $E_\xi D_{\bar{B}} \subset D_{\bar{B}}$ . In consequence of the reality of  $A$  and  $B$ , the operators  $B^*$  and  $E_\xi$  will also be real. Therefore  $M_- = M_+^*$ , i.e., if  $x \in M_+$ , then  $x^* \in M_-$ . Since every self-adjoint extension of the operator  $B$  is a part of  $B^*$ , in order that some extension  $\tilde{B}$  commute with  $E_\xi$ , it is necessary that from  $y \in D_{\tilde{B}}$  it should follow that  $E_\xi y \in D_{\tilde{B}}$ . Let

$$D_{\tilde{B}} = D_{\bar{B}} \dot{+} M_+ \dot{+} VM_+,$$

where  $V$  is an isometric operator mapping  $M_+$  onto  $M_-$ . Then

$$E_\xi D_{\tilde{B}} = E_\xi D_{\bar{B}} \dot{+} E_\xi M_+ \dot{+} E_\xi VM_+,$$

and, in order that  $E_\xi D_{\tilde{B}} \subset D_{\tilde{B}}$ , it is necessary that  $E_\xi V M_+ = V E_\xi M_+$ . Let  $x_0$  be an arbitrary element of  $M_+$ . Denote by  $M_+^{(1)}$  the closure of the linear span of the vectors  $E_\Delta x_0$ , where  $\Delta$  is an arbitrary interval.  $M_+^{(1)} \subset M_+$ . Let  $M_-^{(1)} = M_+^{(1)*} \subset M_-$ . Define the operator  $V^{(1)}$  from  $M_+^{(1)}$  into  $M_-^{(1)}$  as follows:

$$V^{(1)} E_\Delta x_0 = E_\Delta x_0^*,$$

and on the remaining vectors by linearity and continuity.  $V^{(1)}$  maps  $M_+^{(1)}$  isometrically onto  $M_-^{(1)}$ , since the operator  $E_\xi$  and the scalar product  $(x, y)$  are real with respect to the involution. It is obvious that

$$E_\xi V^{(1)} M_+^{(1)} = V^{(1)} E_\xi M_+^{(1)}$$

for any  $\xi$ . If the orthogonal complement to  $M_+^{(1)}$  in  $M_+$  is not empty, then choose there an arbitrary element  $y_0$  and do the same, etc., until  $M_+$  is exhausted.

Thus one can construct an isometric operator  $V$  such that  $E_\xi V M_+ = V E_\xi M_+$ . The corresponding extension  $\tilde{B}$  will commute with  $E_\xi$ :  $\tilde{B} E_\xi x = E_\xi \tilde{B} x$  for all  $x \in D_{\tilde{B}}$ . Such an extension  $\tilde{B}$  is not unique, if  $\bar{B}$  is not a self-adjoint operator. Let  $R_\mu$  be the resolvent of  $\tilde{B}$ . Since

$$(\tilde{B} - \mu E) E_\xi x = E_\xi (\tilde{B} - \mu E) x,$$

then, applying  $R_\mu$  to both sides of the equality, we obtain

$$E_\xi R_\mu y = R_\mu E_\xi y$$

for all  $y \in H$ . Using again the spectral representation of the resolvent and the uniqueness of the Stieltjes transform, we shall have  $E_\xi F_\eta = F_\eta E_\xi$ , where  $F_\eta$  is the spectral family of  $\tilde{B}$ . The theorem is proved.

Let now  $\Phi_1$  and  $\Phi_2$  be nuclear algebras, and let the scalar product be given by a continuous positive-definite functional  $T(x)$ , i.e. one such that  $T(x \cdot x^*) \geq 0$ . Then, assuming the conditions fulfilled

of Theorem 1, from the general scheme in (1) we obtain a representation for  $T(x)$ :

$$T(x) = \int X_\lambda(x) d\sigma(\lambda), \quad (3)$$

where  $X_\lambda$  are the common eigenfunctionals of the operators  $A'$  and  $B'$ , adjoint to  $A$  and  $B$  in  $\Phi'$ . The measure  $\sigma(\lambda)$ , generally speaking, is not unique.

Let us give some applications of Theorem 1.

1. Let  $c_{mn}$  be a double sequence of moments, i.e., for any sequence  $\{\xi_{mn}\}$

$$\sum_{m,n} c_{m_1+m_2, n_1+n_2} \xi_{m_1, n_1} \bar{\xi}_{m_2, n_2} \geq 0. \quad (4)$$

The question is posed of when the two-dimensional moment problem is solvable (see (3)), i.e., when there exists a nonnegative measure  $\sigma(\lambda, \mu)$  such that

$$c_{mn} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \lambda^m \mu^n d\sigma(\lambda, \mu).$$

**Theorem 2.** *The two-dimensional moment problem is solvable if, for any fixed  $n_0$ , the one-dimensional moment problem  $a_m = c_{m, 2n_0} + c_{m, 2(n_0+1)}$  is determinate; moreover, the measure  $\sigma(\lambda, \mu)$ , generally speaking, is not unique.*

Theorem 2 follows from Theorem 1 if we assume that  $\Phi_1$  and  $\Phi_2$  are spaces of sequences in which only a finite number of terms are nonzero, and  $A$  and  $B$  are shift operators, i.e.,  $A\{\xi_m\} = \{\xi_{m-1}\}$  and  $B\{\eta_n\} = \{\eta_{n-1}\}$ ; the involution is simply passage to the complex-conjugate sequence, and the scalar product is defined by means of (4).

2. Let  $f(x, y)$  be a continuous function positive-definite in the sense of Bochner in the rectangle  $P = [-a, a; -b, b]$ , i.e.,

$$\int_P f(x_1 - x_2, y_1 - y_2) \varphi(x_1, y_1) \overline{\varphi(x_2, y_2)} dx_1 dx_2 dy_1 dy_2 \geq 0 \quad (5)$$

for all finite infinitely differentiable functions from  $P$ . In this case  $\Phi_1 = K(-a, a)$ ,  $\Phi_2 = K(-b, b)$ . The operators  $A$  and  $B$  are respectively  $i d/dx$  and  $i d/dy$ . The involution is given by the equality  $[\varphi(x, y)]^* = \overline{\varphi(-x, -y)}$ , and the scalar product is determined by (5).

The inequality holds

$$\begin{aligned} & \int_P f(x_1 - x_2, y_1 - y_2) \varphi(x_1) \overline{\varphi(x_2)} \psi(y_1) \overline{\psi(y_2)} dP \leq \\ & \leq \left( \int_{-b}^b |\psi(y)| dy \right)^2 \int_{-a}^a \int_{-a}^a f(x_1 - x_2, 0) \varphi(x_1) \overline{\varphi(x_2)} dx. \end{aligned} \quad (6)$$

The function  $f(x, 0)$  is obviously a positive-definite function. From (6) it follows that convergence in the scalar product generated by  $f(x, 0)$  implies convergence in the scalar product (5) for any fixed  $\psi(y)$ , which is stronger than the requirement of Theorem 1. Therefore the following theorem holds:

**Theorem 3.** *If  $f(x, 0)$  admits a unique extension to the entire  $X$ -axis as a positive-definite function, then  $f(x, y)$  has the representation*

$$f(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{i(\lambda x + \mu y)} d\sigma(\lambda, \mu),$$

where  $\sigma(\lambda, \mu)$  is a nonnegative measure, generally speaking not unique, so that  $f(x, y)$  admits an extension to the whole plane as a positive-definite function.

Devinatz in (4) proved an analogous theorem under the assumption that  $\hat{f}(x, 0)$  on the interval  $(-a + \varepsilon, a - \varepsilon)$  admits a unique extension to the whole axis as a positive-definite function.

3. Theorem 1 is easily generalized to the case when  $\Phi = \Phi_1 \otimes \dots \otimes \Phi_k$ , the operators  $A_i$  ( $1 \leq i \leq k$ ) act in  $\Phi_i$ , and the closure  $\overline{A_i}$  ( $1 \leq i \leq k - 1$ ) is a self-adjoint operator in the Hilbert space obtained by completing  $\Phi_i$  with respect to the scalar product

$$(\varphi, \varphi)_i = (\psi_{10} \otimes \dots \otimes \psi_{i-1,0} \otimes \varphi \otimes \psi_{i+1,0} \otimes \dots \otimes \psi_{k0},$$

$$\psi_{10} \otimes \dots \otimes \psi_{i-1,0} \otimes \varphi \otimes \psi_{i+1,0} \otimes \dots \otimes \psi_{k0})$$

for arbitrary  $\psi_{j0} \in \Phi_j$  ( $j \neq i$ ). Then there exists a self-adjoint extension  $\tilde{A}_k$  of the operator  $A_k$  such that all  $\overline{A_i}$  ( $1 \leq i \leq k - 1$ ) commute with  $\tilde{A}_k$ . The fact that the  $\overline{A_i}$  commute with one another was proved in (1).

Theorem 2 takes the following form in the multidimensional case: if the one-dimensional moment problem

$$a_{n_i} = c_{2n_{10}, \dots, 2n_{i-1,0}, n_i, 2n_{i+1,0}, \dots, 2n_{k0}} + c_{2n_{10}, \dots, 2n_{i-1,0}, n_i, 2n_{i+1,0}, \dots, 2(n_{k0}+1)}$$

is uniquely solvable for  $1 \leq i \leq k - 1$  and for arbitrary  $n_{j0}$  ( $j \neq i$ ), then the multidimensional moment problem

$$c_{n_1, \dots, n_k} = \int \lambda_1^{n_1} \dots \lambda_k^{n_k} d\sigma(\lambda_1, \dots, \lambda_k)$$

is solvable and, generally speaking, not uniquely.

Similarly, if  $f(x_1, \dots, x_k)$  is a positive-definite function in a  $k$ -dimensional parallelepiped, then if  $f(x_1, 0, \dots, 0)$ ,  $f(0, x_2, \dots, 0)$ ,  $\dots$ ,  $f(0, \dots, x_{k-1}, 0)$  admit unique extensions to the whole axis as positive-definite functions, then  $f(x_1, \dots, x_k)$  also admits an extension to the whole  $k$ -dimensional space as a positive-definite function.

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*Note: Figure translations are in progress. See original paper for figures.*

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