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# MATHEMATICS

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**Abstract**

**Full Text**

**MATHEMATICS**

**Ya. I. ZHITOMIRSKII**

**ON THE CAUCHY PROBLEM FOR A PARABOLIC EQUATION OF SECOND ORDER WITH VARIABLE COEFFICIENTS**

*(Presented by Academician I. G. Petrovskii, 15 V 1957)*

Consider the equation

$$\frac{\partial u}{\partial t} + Lu = f(x, t), \tag{1}$$

where

$$L = - \sum_{i,j=1}^n \frac{\partial}{\partial x_i} a_{ij}(x) \frac{\partial}{\partial x_j} + \sum_{i=1}^n b_i(x) \frac{\partial}{\partial x_i} + c(x), \quad x = (x_1, \dots, x_n).$$

Introduce the notation:

$$\Delta(x) = \det \|a_{ij}(x)\|, \quad \Delta_i(x) = \begin{vmatrix} a_{11} \dots a_{1,i-1} & b_1 & a_{1,i+1} \dots a_{1n} \\ \dots & \dots & \dots \\ a_{n1} \dots a_{n,i-1} & b_n & a_{n,i+1} \dots a_{nn} \end{vmatrix}.$$

Suppose that the coefficients of the operator  $L$  satisfy the following conditions:

$$a_{ij}(x) = a_{ji}(x) \tag{2}$$

for all  $x$ ,

$$\sum_{i,j=1}^n a_{ij}(x) \xi_i \xi_j \geq \alpha^2 \sum_{i=1}^n \xi_i^2 \quad (\alpha > 0) \tag{3}$$

and the coefficient  $c(x)$  is bounded from below:

$$c(x) > -c \quad (c > 0). \quad (4)$$

Let there exist a differentiable function  $g(x)$  such that

$$\frac{\partial g}{\partial x_i} = \frac{\Delta_i(x)}{\Delta(x)}. \quad (5)$$

To equation (1) we append the initial condition

$$u(x, 0) = 0. \quad (1')$$

In the present paper we study questions of existence and uniqueness of the solution of the Cauchy problem (1)–(1') in the whole space  $(x_1, \dots, x_n)$  in the class of rapidly increasing functions.

Following M. I. Vishik (<sup>1</sup>), we define a generalized solution of the Cauchy problem (1)–(1'). For this we shall need to introduce into consideration certain Hilbert spaces of functions.

Let us first consider the space  $\Omega_L$  of finite twice differentiable functions with the scalar product

$$(u, v) = \int_{-\infty}^{\infty} \dots \int u(x, t)v(x, t)e^{-g(x)} dx. \quad (6)$$

We shall show that the operator  $L$  is symmetric on  $\Omega_L$ . Let  $\Omega$  be a domain in the space  $(x_1, \dots, x_n)$ ; let  $\Gamma$  be its boundary; and let  $\mathbf{n}$  be the outward normal to  $\Gamma$ . Putting, in the well-known Ostrogradsky formula,

$$\int_{\Omega} \dots \int \operatorname{div} \mathbf{R} dx = \int_{\Gamma} \dots \int (\mathbf{R}, \mathbf{n}) d\sigma$$

(( $\mathbf{R}, \mathbf{n}$ ) is the ordinary scalar product of two vectors)

$$\mathbf{R} = \left\{ v(x)e^{-g(x)} \sum_{j=1}^n a_{ij}(x) \frac{\partial u}{\partial x_j} \right\},$$

and taking for  $\Omega$  a sphere of radius  $r$ , then letting  $r$  tend to infinity, we easily obtain the formula

$$- \int_{-\infty}^{\infty} \dots \int \sum_{i,j=1}^n \frac{\partial}{\partial x_i} a_{ij}(x) \frac{\partial u}{\partial x_j} v(x)e^{-g(x)} dx = \quad (7)$$

$$= \int_{-\infty}^{\infty} \dots \int \sum_{i,j=1}^n a_{ij}(x) \frac{\partial u}{\partial x_j} \frac{\partial v}{\partial x_i} e^{-g(x)} dx - \int_{-\infty}^{\infty} \dots \int \sum_{i,j=1}^n a_{ij}(x) \frac{\partial u}{\partial x_j} v(x) e^{-g(x)} \frac{\Delta_i(x)}{\Delta(x)} dx.$$

From (7) it follows that

$$\begin{aligned} (Lu, v) &= \int_{-\infty}^{\infty} \dots \int \sum_{i,j=1}^n a_{ij}(x) \frac{\partial u}{\partial x_j} \frac{\partial v}{\partial x_i} e^{-g(x)} dx + \\ &+ \sum_{j=1}^n \int_{-\infty}^{\infty} \dots \int \left( - \sum_{i=1}^n a_{ij}(x) \frac{\Delta_i(x)}{\Delta(x)} + b_j(x) \right) \frac{\partial u}{\partial x_j} v(x) e^{-g(x)} dx + \\ &+ \int_{-\infty}^{\infty} \dots \int c(x) u(x) v(x) e^{-g(x)} dx. \end{aligned}$$

But, by virtue of the definition of the quantities  $\Delta_i(x)$  and  $\Delta(x)$ , it is obvious that

$$\sum_{i=1}^n a_{ij}(x) \frac{\Delta_i(x)}{\Delta(x)} = b_j(x) \quad (j = 1, \dots, n).$$

Therefore

$$(Lu, v) = \int_{-\infty}^{\infty} \dots \int \sum_{i,j=1}^n a_{ij}(x) \frac{\partial u}{\partial x_j} \frac{\partial v}{\partial x_i} e^{-g(x)} dx + \int_{-\infty}^{\infty} \dots \int c(x) u(x) v(x) e^{-g(x)} dx. \quad (8)$$

Interchanging in (8) the roles of the functions  $u(x)$  and  $v(x)$ , and taking (2) into account, we obtain the symmetry of the operator  $L$  on  $\Omega_L$ :

$$(Lu, v) = (u, Lv). \quad (9)$$

From formula (8) and conditions (3), (4) it follows that  $(Lu, u) > -c(u, u)$ . Making in (1) the substitution of the unknown function  $u(x, t) = e^{k_1 t} v(x, t)$  and choosing  $k_1$  so that  $k = k_1 - c > 0$ , we obtain that the operator  $L_1 = L + k$  satisfies the condition

$$(L_1 u, u) > k(u, u). \quad (10)$$

Therefore one may assume that  $L$  also satisfies condition (10).

We now introduce in the space  $\Omega_L$  a new metric by the formula

$$\{u, u\}' = (Lu, u). \quad (11)$$

Complete the space  $\Omega_L$  with respect to the metrics (6) and (11). The resulting complete Hilbert spaces will be denoted respectively by  $H$  and  $H'$ .

Since the bilinear form  $(Lu, v)$  is bounded in the metric (11), it can be realized as the bilinear form of some bounded symmetric operator  $\tilde{L}$  in  $H'$

$$(Lu, v) = \{\tilde{L}u, v\}'. \quad (12)$$

We now consider the operator  $S = \frac{\partial}{\partial t} + L$ , defined on the space of functions  $\Omega_S$  satisfying the following conditions:

- 1)  $u(x, 0) = 0$ ;
- 2)  $u(x, t) \in \Omega_L$  for all  $t \in [0, T]$ ;
- 3)  $\frac{\partial u}{\partial t}$  exists in  $H$  for  $t \in [0, T]$ , and moreover

$$\int_0^T \left( \frac{\partial u}{\partial t}, \frac{\partial u}{\partial t} \right) dt < \infty;$$

- 4)

$$\int_0^T (Lu, Lu) dt < \infty.$$

In the space of functions  $\Omega_S$  we introduce two metrics by the formulas

$$[u, v] = \int_0^T (u, v) dt; \quad (13)$$

$$\{u, v\}_1 = \int_0^T \{u, v\}' dt. \quad (14)$$

Complete the space  $\Omega_S$  with respect to the metrics (13) and (14). The resulting complete Hilbert spaces will be denoted respectively by  $\mathfrak{A}$  and  $\mathfrak{A}'$ . Since

$$\left| \left[ \frac{\partial u}{\partial t}, v \right] \right| \leq \left[ \frac{\partial u}{\partial t}, \frac{\partial u}{\partial t} \right] \cdot [v, v] \leq C_1 \int_0^T (Lv, v) dt \leq C_2 \{v, v\}_1,$$

the bilinear form  $\left[\frac{\partial u}{\partial t}, v\right]$  can be realized in  $\mathfrak{A}'$  in the form of the form of some operator  $\tilde{T}$

$$\left[\frac{\partial u}{\partial t}, v\right] = \{\tilde{T}u, v\}_1.$$

At the same time

$$[Lu, v] = \int_0^T (Lu, v) dt = \int_0^T \{\tilde{L}u, v\}' dt = \{\tilde{L}u, v\}_1.$$

Thus,

$$[Su, v] = \{\tilde{T}u, v\}_1 + \{\tilde{L}u, v\}_1. \quad (15)$$

Now assuming that  $v(x, t)$  belongs to the space  $\Omega^*$ , consisting of functions  $v(x, t)$  belonging to  $\Omega_S$  and satisfying the condition  $v(x, T) = 0$  (obviously,  $\Omega^*$  is dense in  $\Omega_S$  in the metric (14)), one may in (15) transfer the operator  $\tilde{T}$  to the function  $v(x, t)$

$$[Su, v] = \{u, \tilde{T}^*v\}_1 + \{u, \tilde{L}^*v\}_1. \quad (16)$$

From (16) it follows that, if  $u(x, t)$  is an ordinary solution of equation (1), then

$$[f, v] = \{u, (\tilde{T}^* + \tilde{L}^*)v\}_1. \quad (17)$$

By a generalized solution of problem (1)–(1') we shall mean a function  $u(x, t)$  belonging to the space  $\mathfrak{A}'$ , for which equality (17) holds for all  $v(x, t)$  from the space  $\Omega^*$  (1).

Since the operator  $L$  satisfies conditions (9) and (10) (symmetry and positive definiteness), it follows, by the results of (1) (see also (2)), that the following theorem holds.

**Theorem.** *For any function  $f(x, t)$  satisfying the condition*

$$\int_0^T (f, f) dt < \infty,$$

*there exists a generalized solution of problem (1)–(1'), and it is unique.*

Let us consider, as an example, an equation of the form (1), in which

$$Lu = -\Delta u + \sum_{i=1}^n b_i(x_i) \frac{\partial u}{\partial x_i} + c(x)u,$$

where  $\Delta$  is the  $n$ -dimensional Laplace operator;  $b_i$  depends only on  $x_i$ ;  $x = (x_1, \dots, x_n)$ . It is easy to see that conditions (2)–(5) are satisfied here.

In this case

$$g(x) = \sum_{i=1}^n \int_0^{x_i} b_i(y) dy.$$

Then for any function  $f(x, t)$  satisfying the condition

$$\int_0^T \int_{-\infty}^{\infty} \dots \int f^2(x, t) e^{-g(x)} dx dt < \infty,$$

there exists a generalized solution of the Cauchy problem (1)–(1'), and it is unique.

In particular, for  $n = 1$ ,  $b(x) = x^{2m-1}$ , we obtain uniqueness of the solution of the Cauchy problem (1)–(1') in the class of functions  $u(x, t)$  satisfying the conditions

$$|u(x, t)| \leq A_1 e^{A_2 |x|^{2m-\varepsilon}}, \quad \left| \frac{\partial u(x, t)}{\partial x} \right| \leq A_3 e^{A_4 |x|^{2m-\varepsilon}},$$

for any  $\varepsilon > 0$ ;  $0 \leq t \leq T$ ;  $A_1, A_2, A_3, A_4$  are constants.

These examples show that the uniqueness class for the solution of the Cauchy problem (1)–(1') can be extended without bound (in the sense of the admissible growth as  $|x| \rightarrow \infty$  of the functions entering it), depending on the behavior as  $|x| \rightarrow \infty$  of the coefficients of equation (1).

The mixed problem for equation (1) for the half-space  $x_i > 0$  is posed as follows: one seeks a solution of equation (1) satisfying condition (1') and the boundary condition

$$u(x_1, \dots, x_n, t) \Big|_{x_i=0} = 0. \quad (1'')$$

Problem (1)–(1')–(1'') is reduced in the usual way <sup>(1)</sup> to problem (1)–(1'). For this it is necessary to impose on the functions of the space  $\Omega_L$  the additional restriction  $u(x, t) = 0$  for  $x_i \leq 0$ . Then, as above:

**Theorem.** *For any function  $f(x, t)$  satisfying the condition*

$$\int_0^T (f, f) dt < \infty,$$

*there exists a generalized solution of problem (1)–(1')–(1''), and it is unique.*

The examples analyzed above show that in this case as well the uniqueness class for the solution of problem (1)–(1')–(1'') contains functions whose admissible growth as  $|x| \rightarrow \infty$  can increase without bound as the growth as  $|x| \rightarrow \infty$  of the coefficients of equation (1) increases.

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*Note: Figure translations are in progress. See original paper for figures.*

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