

# ASYMPTOTIC BEHAVIOR OF SOLUTIONS OF A NONLINEAR BOUNDARY-VALUE PROBLEM IN BOUNDARY-LAYER THEORY

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**Abstract**

**Full Text**

**HYDROMECHANICS**

**G. V. GIL' and A. D. MYSHKIS**

**ASYMPTOTIC BEHAVIOR OF SOLUTIONS  
OF A NONLINEAR BOUNDARY-VALUE  
PROBLEM IN BOUNDARY-LAYER THEORY**

*(Presented by Academician L. I. Sedov on 9 XI 1956)*

In boundary-layer theory the following boundary-value problem is well known:

$$\begin{aligned} y'''(t) + 2y(t)y''(t) + 2\beta(k^2 - [y'(t)]^2) &= 0 \quad (0 \leq t < \infty), \\ y(0) = y'(0) &= 0, \quad y'(\infty) = k \quad (\beta \geq 0, k > 0). \end{aligned} \quad (1)$$

This problem has been considered by a number of authors (see <sup>(1)</sup>, references on pp. 80, 83, 113, 118, 129; see also the later works <sup>(2-5)</sup>).

The purpose of this note is to investigate the asymptotic behavior, as  $t \rightarrow \infty$ , of the solution of the boundary-value problem (1).

It is of interest to obtain a method for investigating the asymptotic behavior of solutions of a broader class of boundary-value problems, including (1) as a special case.

Let  $y(t)$  ( $0 \leq t < \infty$ ) be a solution of the boundary-value problem (1).

**Lemma 1.**

$$y'(t) < k \quad (0 \leq t < \infty).$$

**Proof.** Suppose that the assertion is false and that  $\tau$  is the first root of the equation  $y'(t) = k$ . Since  $y'(t) < k$  ( $0 \leq t < \tau$ ), we have  $y''(\tau) \geq 0$ . If  $y''(\tau) = 0$ , then, by the uniqueness theorem for the solution of a differential equation with prescribed initial conditions,

$$y(t) \equiv y(\tau) + (t - \tau)k,$$

i.e.  $y'(0) = k$ , which is impossible. Hence  $y''(\tau) > 0$ , i.e.  $y'(t)$  increases at the point  $\tau$ . But  $y'(\infty) = k$ , and therefore  $y'(t)$  attains, at some point  $t_1 > \tau$ , its greatest value on the interval  $[\tau, \infty)$ . At the point  $t_1$  we have  $y'' = 0$ ,  $y''' \leq 0$ ; but  $y'(t_1) > k$ , whence

$$y'''(t_1) = -2\beta(k^2 - [y'(t_1)]^2) > 0$$

if  $\beta > 0$ . If, however,  $\beta = 0$ , then, as above,

$$y(t) \equiv y(t_1) + (t - t_1)y'(t_1),$$

i.e.  $y'(0) = y'(t_1) > 0$ .

The contradiction obtained in all cases proves Lemma 1.

We transform the boundary-value problem (1) to a form more convenient for further consideration. Let  $t_0 \geq 0$  be the last root of the equation  $y'(t) = 0$ . Then for  $t \geq t_0$  the function  $y(t)$  increases, and therefore  $t$ , and hence also  $y'(t)$ , is a function of  $y$ . Making the substitution of the unknown function

$$v = k^2 - [y'(t)]^2 = v(y) \quad (y_0 = y(t_0) \leq y < \infty),$$

we obtain the boundary-value problem

$$\begin{aligned} \sqrt{k^2 - v(y)} v''(y) + 2yv'(y) - 4\beta v(y) &= 0 \quad (y_0 < y < \infty), \\ v(y_0) = k^2, \quad v(\infty) = 0, \quad 0 < v(y) < k^2 & \quad (y_0 < y < \infty). \end{aligned} \quad (2)$$

**Lemma 2.**

$$v'(y) < 0 \quad (y_0 < y < \infty).$$

**Proof.** At the roots of the equation  $v'(y) = 0$

$$v''(y) = \frac{4\beta v}{\sqrt{k^2 - v}} > 0 \quad (\text{for } \beta > 0),$$

i.e., all stationary points of  $v$  must be minima of  $v(y)$ , which is impossible; in the case  $\beta = 0$ , with  $v'(y_1) = 0$  ( $y_1 > y_0$ ), we would have  $v(y) \equiv v(y_1)$ , which is also impossible. Lemma 2 is proved.

**Lemma 3.** *The function  $v''(y)$  has, for  $y > y_0$ , at most one zero.*

**Proof.** From relations (2) it is clear that for  $y > y_0$  the derivative  $v'''(y)$  exists and

$$\sqrt{k^2 - v} v''' - \frac{1}{2\sqrt{k^2 - v}} v'' v' + 2yv'' + 2(1 - 2\beta)v' = 0. \quad (3)$$

Hence, if  $\beta \neq \frac{1}{2}$ , then at the points where  $v''(y) = 0$ ,

$$(2\beta - 1)v''' = \frac{2(2\beta - 1)^2}{\sqrt{k^2 - v}} v' < 0$$

(by Lemma 2), i.e.  $v'''$  has one and the same sign. Our assertion follows from this.

Let now  $\beta = \frac{1}{2}$ . Then, if  $v''(y_1) = 0$  ( $y_1 > y_0$ ), by the uniqueness theorem for a solution applied to equation (3),

$$v(y) \equiv v(y_1) + (y - y_1)v'(y_1) \quad (y_0 < y < \infty).$$

Since  $v(\infty) = 0$ , we obtain  $v'(y_1) = 0$ , which contradicts Lemma 2. Lemma 3 is proved.

**Corollary.**

$$v'(\infty) = 0.$$

This relation, in view of the convergence of the integral  $\int_{y_0}^{\infty} v'(y) dy$ , follows at once from Lemma 3.

**Lemma 4.**

$$v''(\infty) = 0.$$

**Proof.** Since  $v'(y) < 0$  ( $y_0 < y < \infty$ ),  $v'(\infty) = 0$ , and for some  $y^* > y_0$  the function  $v''(y)$  ( $y \geq y^*$ ) has no zeros, it follows that  $v''(y) > 0$  ( $y \geq y^*$ ). It remains to prove that the set of all zeros of the function  $v'''(y)$  is bounded above. Suppose this is not so; then, at points where  $v'''(y) = 0$ , as a result of differentiation we obtain

$$\sqrt{k^2 - v} v^{IV} = v'' \left[ \frac{1}{4\sqrt{(k^2 - v)^3}} v'^2 + \frac{1}{2\sqrt{k^2 - v}} v'' + 4(\beta - 1) \right]. \quad (4)$$

If  $\beta \geq 1$ , then at these points, for  $y \geq y^*$ , we have  $v^{IV}(y) > 0$ , whence a contradiction easily follows.

If  $0 \leq \beta < 1$  and  $v''(y) \not\rightarrow 0$  (as  $y$  tends to  $\infty$ ), choose a number  $\varepsilon \in (0, \min\{8k(1-\beta), \overline{\lim}_{y \rightarrow \infty} |v''(y)|\})$ , and denote by  $F$  the closed set of points on the ray  $[y^*, \infty)$  at which  $v''(y) \geq \varepsilon$ , and by  $E$  its complement. Since  $\int_{y_0}^{\infty} v''(y) dy < \infty$ , the set  $E$  (as also  $F$ ) is unbounded. Take any of the intervals of which  $E$  consists. At its endpoints  $v''(y) = \varepsilon$ , and inside it  $v''(y) < \varepsilon$ . Hence, on this interval, at

some point  $\eta$ , a minimum of  $v''(y)$  is attained. Take a sequence of points  $\eta_i$  for which  $\lim \eta_i = \infty$ ; for sufficiently large  $i$ , at the points  $\eta_i$ ,

$$\frac{1}{2\sqrt{k^2 - v}} v'' + \frac{1}{4\sqrt{(k^2 - v)^3}} v'^2 < \frac{\varepsilon}{2k} + o(1) < 4(1 - \beta),$$

and from (4) we obtain  $v^{\text{IV}}(\eta_i) < 0$ . But this contradicts the fact that at the points  $\eta_i$  a minimum of  $v''(y)$  is attained. Lemma 4 is proved.

**Corollary.**

$$\lim_{y \rightarrow \infty} yv'(y) = 0.$$

This relation follows immediately from equation (2) and the boundary conditions.

**Lemma 5.** If  $\beta > 0$ , then

$$\lim_{y \rightarrow \infty} \frac{yv'(y)}{v(y)} = -\infty.$$

**Proof.** Let  $yv'v^{-1} = z$ . Then, with the aid of (2), we obtain

$$z = \frac{4\beta y}{\sqrt{k^2 - v}} + z \left( \frac{1}{y} - \frac{2y}{\sqrt{k^2 - v}} - \frac{z}{y} \right). \quad (5)$$

If for some  $y > y_0$  we have  $z'(y) = 0$ , then for this  $y$

$$z'' = 4\beta \frac{2(k^2 - v) + yv'}{2(k^2 - v)^{3/2}} + z \left( -\frac{1}{y^2} - \frac{2(k^2 - v) + yv'}{(k^2 - v)^{3/2}} + \frac{z}{y^2} \right).$$

In view of the corollary to Lemma 4 and the fact that  $z < 0$  ( $y \geq y_0$ ), for sufficiently large  $y$ , at the stationary points of  $z(y)$  we shall have  $z''(y) > 0$ . Hence it follows that  $z(y)$ , beginning with some  $y$ , is a monotone function.

Suppose that  $z(\infty) = c > -\infty$  ( $c \leq 0$ ). From relation (5) we obtain  $z'(\infty) = \infty$ , which is impossible. Thus  $z(\infty) = -\infty$ , as was required to prove.

**Theorem.** The solution of the boundary-value problem (1) and its first and second derivatives have, as  $t \rightarrow \infty$ , the following asymptotic representation:

$$y(t) = kt - C + t^{-2-2\beta+o(1)} e^{-kt^2+2Ct},$$

$$y'(t) = k - t^{-1-2\beta+o(1)} e^{-kt^2+2Ct},$$

$$y''(t) = t^{-2\beta+o(1)}e^{-kt^2+2Ct},$$

where  $C > 0$  is a certain constant (equal to  $\lim_{t \rightarrow \infty} (kt - y(t))$ ).

**Proof.** From equation (2) we obtain

$$\frac{v''(y)}{v'(y)} = -\frac{2y}{\sqrt{k^2 - v}} + \frac{4\beta v}{v'\sqrt{k^2 - v}} = -\frac{2y}{k} + 4yu(y),$$

where, by Lemma 5,  $u(y) \rightarrow 0$  as  $y \rightarrow \infty$ . Hence, integrating over the interval from some  $y_1 > y_0$ , we obtain

$$v'(y) = -C_1 e^{-\frac{y^2}{k} + 4 \int_{y_1}^y su(s) ds} \equiv -C_1 \chi(y) \quad (C_1 > 0; y > y_0) \quad (6)$$

and, consequently,

$$\lim_{y \rightarrow \infty} \frac{yv}{v'} = \lim_{y \rightarrow \infty} \left( -\frac{\int_y^\infty \chi(s) ds}{y^{-1}\chi(y)} \right) = -\frac{k}{2} \quad (7)$$

(the limit is easily found by l' Hospital' s rule),

$$v = \frac{C_2 + o(1)}{y} \chi(y) \quad \left( y_0 < y < \infty; C_2 = \frac{k}{2} C_1 \right). \quad (8)$$

Next, by l' Hospital' s rule we have

$$\lim_{y \rightarrow \infty} y^2 (k - \sqrt{k^2 - v}) = \lim_{y \rightarrow \infty} \frac{v'/2\sqrt{k^2 - v}}{-2y^{-3}} = \lim_{y \rightarrow \infty} \frac{-C_1 y^3 \chi(y)}{-4\sqrt{k^2 - v}} = 0,$$

whence, in view of (7),  $\lim_{y \rightarrow \infty} y^2 u(y) = -\frac{\beta}{2}$ . Therefore, by l' Hospital' s rule,

$$\lim_{y \rightarrow \infty} \frac{\int_{y_1}^y su(s) ds}{\ln y} = \lim_{y \rightarrow \infty} y^2 u(y) = -\frac{\beta}{2}. \quad (9)$$

Therefore, by virtue of (8),

$$\sqrt{k^2 - v} = k - y^{-1-2\beta+o(1)} e^{-\frac{y^2}{k}}. \quad (10)$$

Thus, first of all, since  $\sqrt{k^2 - v} = y'(t)$ , we have  $y(t) \sim kt$ , and on the basis of (10)

$$y'(t) = k - e^{-kt^2 + o(t^2)} \quad (t \geq t_0).$$

Integrating, we obtain

$$y = kt - C + \int_t^\infty e^{-k\tau^2 + o(\tau^2)} d\tau \quad \left( C = kt_0 - y_0 + \int_{t_0}^\infty e^{-k\tau^2 + o(\tau^2)} d\tau \right).$$

Put, for  $t \geq t_0$ ,

$$\int_t^\infty e^{-k\tau^2 + o(\tau^2)} d\tau = e^{-kt^2 + \gamma(t)t^2}.$$

Then it is easy to show that  $\gamma(\infty) = 0$ . Indeed, by l' Hospital' s rule,

$$\frac{\int_t^\infty e^{-\alpha\tau^2} d\tau}{e^{-\alpha t^2} t^{-1}} \xrightarrow{t \rightarrow \infty} \frac{1}{2\alpha} \quad (\alpha > 0);$$

but for any  $\varepsilon > 0$  ( $\varepsilon < k$ ), for sufficiently large  $t$ ,

$$e^{-(k+\varepsilon)t^2} < e^{-kt^2 + o(t^2)} < e^{-(k-\varepsilon)t^2},$$

and therefore

$$t^{-1} e^{-(k+\varepsilon)t^2} \left( \frac{1}{2(k+\varepsilon)} + o(1) \right) < \int_t^\infty e^{-k\tau^2 + o(\tau^2)} d\tau < t^{-1} e^{-(k-\varepsilon)t^2} \left( \frac{1}{2(k-\varepsilon)} + o(1) \right),$$

whence our assertion follows at once. Thus,

$$y = kt - C + e^{-kt^2 + o(t^2)}. \quad (11)$$

Hence  $C = \lim_{t \rightarrow \infty} (kt - y)$ ; from Lemma 1 it follows that the difference  $kt - y(t)$ , for  $0 \leq t < \infty$ , increases, and, since  $y(0) = 0$ ,  $C > 0$ .

Substituting (11) into (10), we obtain the expression for  $y'(t)$  given in the formulation of the theorem. Further, for  $t > t_0$ , according to (6),

$$y''(t) = \frac{-v'(y)}{2} = \frac{C_1}{2} e^{-\frac{y^2}{k} + 4 \int_{y_1}^y su(s) ds} = y^{-2\beta + o(1)} e^{-\frac{y^2}{k}}$$

(see (9)). Using (11), we obtain the required expression for  $y''(t)$ .

Finally, integrating the expression for  $y'(t)$  and taking into account that

$$\int_t^\infty \tau^{-1-2\beta+o(1)} e^{-k\tau^2+C\tau} d\tau = t^{-2-2\beta+o(1)} e^{-kt^2+2Ct}$$

(the proof is analogous to the proof that  $\gamma(\infty) = 0$ ), and then comparing the result obtained with (11), we obtain the required representation of  $y(t)$ .

Let us note in conclusion that the value of the constant  $C$  is unknown. It would be important to obtain an estimate, or a method for approximate computation, or an expansion in a series of some kind, etc., for this constant.

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*Note: Figure translations are in progress. See original paper for figures.*

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