
AI translation · View original & related papers at
chinaxiv.org/items/chinaxiv-202306.00575

Two-proton sequential decay from excited states of ^{18}Ne (postprint)

Authors: YU Ning, MAGLIONE Enrico, FERREIRA Lidia

Date: 2023-06-18T00:00:00+00:00

Abstract

Two-proton radioactivity from ^{18}Ne is discussed in terms of sequential decay. The branch ratios for one-proton emission from excited states are calculated, which including spectroscopic factors, obtained from a Shell-model calculation with realistic interactions. The branch ratios show that the two-proton emission from the 1^- state of ^{18}Ne at 7.94 MeV is most likely to go through the sequential decay. The same mechanism is discussed for other excited states at higher energy by different interactions.

Full Text

Preamble

Two-proton sequential decay from excited states of ^{18}Ne

YU Ning^{1,*}, MAGLIONE Enrico², FERREIRA Lidia^{3}

¹Institute of Particle Physics, Huazhong Normal University, Wuhan 430079, China

²Dipartimento di Fisica “G.Galilei”, Padova, Italy and Istituto Nazionale di Fisica Nucleare, Padova, Italy

³Centro de Física das Interações Fundamentais, and Departamento de Física, Instituto Superior Técnico, Lisbon, Portugal

Abstract

Two-proton radioactivity from ^{18}Ne is discussed in terms of sequential decay. The branch ratios for one-proton emission from excited states are calculated, including spectroscopic factors obtained from a shell-model calculation with realistic interactions. The branch ratios show that two-proton emission from the 1^- state of ^{18}Ne at 7.94 MeV is most likely to proceed through sequential decay. The same mechanism is discussed for other excited states at higher energies using different interactions.

Key words: Two-proton radioactivity, Nuclear shell model, Branch ratios

Introduction

Proton radioactivity, first experimentally observed as ground-state decay at GSI in 1981, has provided crucial information on the structure of nuclei beyond the proton drip line. The more complex decay mode of two-proton radioactivity, proposed 50 years ago in a seminal article [1], opened a new window on nucleon-nucleon correlations and nuclear structure. In 2002, the simultaneous emission of two protons was observed for the first time in the decay of ^{45}Fe by Pfützner, Giovanazzo, and collaborators in experiments at GSI and GANIL [2,3].

Research in this field flourished following this breakthrough, and to date ^{54}Zn [4], ^{48}Ni [5], ^{19}Mg [6], ^{16}Ne [7], ^{17}Ne [8], ^{18}Ne [9], ^{10}C [10], ^{14}O [11], and ^{29}S [12] have been found to exhibit two-proton emission. Several theoretical approaches have been applied to this problem, including the di-proton model [13,14], R-matrix approach [15], continuum shell model [16], adiabatic hyperspherical approach [17], and the quantum three-body cluster approach [18], where tunneling through the barrier is treated dynamically. There are two distinct decay modes for simultaneous two-proton emission: (1) three-body direct breakup involving uncorrelated emission of the two protons, commonly referred to as democratic decay; and (2) ^2He cluster emission, where a pair of protons correlated in a quasi-bound 1S configuration breaks up upon emission (di-proton emission). The two protons exhibit strong angular and energy correlations, with the ^2He appearing as a resonance at $20\text{ MeV}/c$ in the two-proton relative momentum distribution [19].

Microscopic calculations for the one- and two-proton decays of the $6.15\text{ MeV } 1^-$ state of ^{18}Ne were presented in Ref. [20]. It was found that for two-proton emission, sequential decay through a ghost of the $1/2^+$ state is within a factor of three of the democratic decay assumption. The calculated width for di-proton emission is only about a factor of two smaller than that for sequential decay, indicating that the observed decay may be a combination of both processes. In the excitation-energy spectrum of ^{18}Ne from Ref. [9], it is notable that some states appear in the two-proton emission channel ($^{18}\text{Ne} \rightarrow ^{16}\text{O} + 2\text{p}$) but not in the one-proton emission channel ($^{18}\text{Ne} \rightarrow ^{17}\text{F} + \text{p}$). This suggests that ^{17}F is not produced in its ground state in these decays, making sequential decay the most likely mechanism for two-proton emission from these states.

In this paper, we present microscopic shell-model calculations for sequential two-proton decay in ^{18}Ne from excited states using different Hamiltonians.

2 Calculation and Discussion

The spectroscopic factor is the most important quantity needed to obtain the decay width. To calculate it, we perform a shell-model calculation to obtain the wave functions for ^{18}Ne . The model space includes the 0s , 0p , $1\text{s}0\text{d}$, and $1\text{p}0\text{f}$

orbits, with ^{16}O treated as a closed s^4p^{12} shell. The low-lying positive-parity states of ^{17}F and ^{18}Ne are described as $s^4p^{12}(\text{sd})^1$ and $s^4p^{12}(\text{sd})^2$ configurations, while the low-lying negative-parity states are treated as $1\hbar\omega$ excitations of the forms $s^4p^{11}(\text{sd})^2$, $s^4p^{12}(\text{pf})^1$ and $s^4p^{11}(\text{sd})^3$, $s^4p^{12}(\text{sd})^1(\text{pf})^1$. Thus, the emitted protons in ^{18}Ne and ^{17}F originate from the (sd)(pf) shells. Two Hamiltonians designed for these model spaces are selected for calculating the wave functions: the WBP and WBT interactions [21]. We employ a simple shell-model code developed by our group, in which spurious states are removed by the standard method [22] of adding a center-of-mass Hamiltonian to the interaction.

The calculated excitation energies of these low-lying states are shown in [Figure 1: see original paper]. Some states are in reasonable agreement with the experimental energies found in ^{18}Ne . The low-lying negative-parity states are dominated by the $s^4p^{11}(\text{sd})^3$ configuration, but the smaller $s^4p^{12}(\text{sd})^1(\text{pf})^1$ component is responsible for one- and two-proton decay. The shell-model spectroscopic factors are obtained from the wave functions of ^{18}Ne and ^{17}F .

Decays from positive-parity states of ^{18}Ne to positive-parity states of ^{17}F , and from negative-parity states of ^{18}Ne to negative-parity states of ^{17}F , can proceed via $0d$ -shell wave emission or $1s$ -shell wave emission. Decays from positive-parity states of ^{18}Ne to negative-parity states of ^{17}F , and from negative-parity states of ^{18}Ne to positive-parity states of ^{17}F , can proceed via $0f$ -shell wave emission or $1p$ -shell wave emission. Because the $s^4p^{11}(\text{sd})^3$ component in ^{18}Ne is much larger than that of $s^4p^{12}(\text{sd})^1(\text{pf})^1$, the spectroscopic factors are larger in channels involving positive-parity states in ^{18}Ne .

According to scattering theory, the half-life for decay from an initial state i to a final state f by single-particle emission is given by:

$$\ln 2/T_{1/2} = \Gamma$$

where the decay width can be obtained from the relation [25,26]:

$$\Gamma_{if} = S_{if}^j \Gamma_{sp}^j$$

Here S_{if}^j is the spectroscopic factor corresponding to the probability that removing a particle j with angular momentum j from an initial state i will lead to a final state f , and Γ_{sp}^j is the single-particle width. α_j is the asymptotic normalization of the proton single-particle wave function in a state of spin j .

[Figure 1: see original paper] shows WBP and WBT predictions for the low-lying $T = 1$ energy spectrum of ^{18}Ne . Some levels are labeled by J^π and E_x . The experimental data [23,24] are presented in the right column. The J^π of levels that are not labeled are unknown.

The total width for decay is a sum of partial widths:

$$\Gamma = \sum \Gamma_{if}$$

The branching ratios are simply the ratio between a partial decay width and the total width:

$$B_{if} = \Gamma_{if}/\Gamma$$

For the fourth 1^- state at 7.94 MeV in ^{18}Ne , we find that the spectroscopic factor for decay to the $1/2^-$ third excited state ($Q_{1p} = 0.9141$ MeV) is considerably larger than that for decay to the $5/2^+$ ground state ($Q_{1p} = 4.0184$ MeV) of ^{17}F . The spectroscopic factors and widths for each channel are shown in . From this table, we can see that the branching ratio for decay to the $1/2^-$ state is larger than those for decay to the ground state and first excited state because it has a larger spectroscopic factor, even though it has a smaller single-particle width. We conclude that the 7.94 MeV 1^- state is the most likely candidate for two-proton sequential decay. This explains why it can be observed in the two-proton emission channel but not in the one-proton channel. Other states in this situation include the 3^- state around 9–10 MeV (9.809 MeV for WBP and 10.099 MeV for WBT) and the 5^- state near 13 MeV (13.412 MeV for WBP and 13.200 MeV for WBT).

Spectroscopic factors from the state $J^\pi = 1^-$, $E_x = 7.94$ MeV. The channel $5/2^+$ proton is from the $0f_{7/2}$ shell and decays to the $5/2^+$ ground state in ^{17}F . The last line shows the total widths for single-proton emission from ^{18}Ne . $\otimes 0f_{7/2}$ means that the emitted proton originates from the $0f_{7/2}$ orbit.

E_x (MeV)	Channel	$\otimes 0f_{7/2}$	$\otimes 0f_{5/2}$	$\otimes 1p_{3/2}$	$\otimes 1p_{1/2}$	$\otimes 0d_{3/2}$	$\otimes 1s_{1/2}$	Expt.
Spectroscopic factor								
Γ_{sp} (keV)								
$\otimes(0f+1p)$								
$\otimes(0d+1s)$								
Total Γ								≤ 50 keV
BrWBP								
BrWBT								

3 Conclusion

We have presented preliminary results for proton decay branching ratios and decay widths in ^{18}Ne using a shell-model calculation. The results obtained for the branching ratio from the 1^- state at 7.94 MeV in ^{18}Ne show that this state is the most likely candidate for sequential two-proton decay. The 3^- and 5^- states can also be candidates for the same process.

References

1. Goldansky V. Nucl Phys, 1960, 19: 482-495.
2. Pfützner M, Badura E, Bingham C, et al. Eur Phys J A, 2002, 14: 279-285.
3. Giovinazzo J, Blank B, Chartier M, et al. Phys Rev Lett, 2002, 89: 102501.
4. Blank B, Bey A, Canchel C, et al. Phys Rev Lett, 2005, 94: 232501.
5. Dossat C, Bey A, Blank B, et al. Phys Rev C, 2005, 72: 054315.
6. Mukha I, Summerer K, Acosta L, et al. Phys Rev Lett, 2007, 99: 182501.
7. Mukha I, Giovinazzo J, Summerer K, et al. Phys Rev C, 2008, 77: 061303(R).
8. Zerguerras T, Blank B, Blumenfeld Y, et al. Eur Phys J, 2004, A20: 389-396.
9. Raciti G, Cardella G, Napoli M D, et al. Phys Rev Lett, 2008, 100: 192503.
10. Mercurio K, Charity R J, Shane R, et al. Phys Rev C, 2008, 78: 031602(R).
11. Bain C R, Woods P J, Coszach R, et al. Phys Lett B, 1996, 373: 35-39.
12. (Reference missing in original)
13. (References 13-14 missing in original)
14. Bartlett A J, Tostevin J A, Thompson I J. Phys Rev C, 2008, 78: 054603.
15. Rotureau J, Okolowicz J, Ploszajczak M. Phys Rev Lett, 2005, 95: 042503.
16. Nielsen E, Fedorov D V, Jensen A S, et al. Phys Rep, 2001, 347: 374-460.
17. Grigorenko L, Mukha I, Thompson I, et al. Phys Rev Lett, 2002, 88: 042502.
18. Brown B A, Barker F C, Millener D J. Phys Rev C, 2002, 65: 051309(R).
19. (Reference for WBP/WBT interactions missing in original)
20. (Reference for spurious state removal missing in original)
21. Bardayan D W, Blackmon J C, Brune C R, et al. Phys Rev C, 2000, 62: 055804.
22. Hahn K I, Garcia A, Adelberger E G, et al. Phys Rev C, 1996, 54: 1999-2013.
23. Maglione E, Ferreira L S, Liotta R J. Phys Rev Lett, 1998, 81: 538-541.
24. Ferreira L S, Maglione E, Liotta R J. Phys Rev Lett, 1997, 78: 1640-1643.

Supported by the Fundação para a Ciência e a Tecnologia (Portugal), Project: PTDC/FIS/68340/2006 and CERN/FP/116385/2010

Corresponding author. E-mail address: yuning_{ok}@hotmail.com

Note: Figure translations are in progress. See original paper for figures.

Source: ChinaXiv – Machine translation. Verify with original.