

Transitional Area of Ce⁴⁺ to Ce³⁺ in Sm_xCayCe_{1-x-y}O₂ with Various Doping and Oxygen Vacancy Concentrations: A GGA + U Study (Postprint)

Authors: WU Tong-Wei, JIA Gui-Xiao, WANG Xiao-Xia, Lei Li, AN Sheng-Li

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Abstract

In this work, we perform DFT+U periodic calculations to investigate the geometric and electronic structures, as well as oxygen vacancy formation energies, of Sm_xCayCe_{1-x-y}O₂ systems ($x = 0.0312, 0.0625, 0.125, \text{ and } 0.250$; $y = 0.0312, 0.0625, 0.125, \text{ and } 0.250$; $z = 0.0312, 0.0625, 0.125, 0.250, \text{ and } 0.50$) with varying oxygen vacancy and doping concentrations. The calculated results show that V1-Sm³⁺-V2 structures, where V1 and V2 (both nearest to Sm³⁺) are positioned along a face diagonal, exhibit the lowest energy configurations. Electronic structure analysis of the Sm_xCayCe_{1-x-y}O₂ systems reveals that excess electrons generated from oxygen vacancies are localized on the f-level traps of neighboring Ce ions, and that Ca²⁺ and Sm³⁺ co-doping effectively suppresses the reduction of Ce⁴⁺. To prevent the formation of Ce³⁺, both x and y must be greater than 0.0625 when $z = 0.125$, or z must be less than 0.125 when $x = y = 0.0625$. The Ce³⁺/Ce⁴⁺ transformation ratio k exhibits a clear monotonic increase with increasing oxygen vacancy concentration. The introduction of Sm³⁺ reduces k . Furthermore, doped Sm³⁺ can suppress the reduction of Ce⁴⁺ when the V1-Sm³⁺-V2 structure with a face-diagonal positional relationship exists in a less reducing atmosphere. It should be noted that the Sm_{0.25}Ce_{0.75}O_{1.5} system should be considered as Sm-doped Ce₂O₃.

Full Text

Transitional Area of Ce⁴⁺ to Ce³⁺ in Sm_xCayCe_{1-x-y}O₂ with Various Doping and Oxygen Vacancy Concentrations: A GGA + U Study

WU Tong-Wei; JIA Gui-Xiao; WANG Xiao-Xia; LI Lei; AN Sheng-Li

- (1) School of Materials and Metallurgy, Inner Mongolia University of Science and Technology, Baotou, Inner Mongolia 014010, China;
- (2) The Inner Mongolia Autonomous Region Key Lab of Novel Functional Ceramics and Devices, Baotou, Inner Mongolia 014010, China

ABSTRACT

In this work, we perform DFT+U periodic calculations to investigate the geometric and electronic structures as well as oxygen vacancy formation energies of $\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ systems ($x = 0.0312, 0.0625, 0.125$ and 0.250 ; $y = 0.0312, 0.0625, 0.125$ and 0.250 ; $z = 0.0312, 0.0625, 0.125, 0.250$ and 0.50) with varying oxygen vacancy and doping concentrations. Our calculations reveal that the V1-Sm³⁺-V2 structures, where V1 and V2 (both nearest to Sm³⁺) adopt a face-diagonal positional relationship, represent the lowest-energy configurations. Analysis of the electronic structures shows that excess electrons originate from oxygen vacancies and become localized on the f-level traps of neighboring Ce atoms, while Ca²⁺ and Sm³⁺ co-doping effectively suppresses the reduction of Ce⁴⁺ to Ce³⁺. To prevent the formation of Ce³⁺, both x and y must exceed 0.0625 when $z = 0.125$, or z must be less than 0.125 when $x = y = 0.0625$. The Ce³⁺/Ce⁴⁺ transformation ratio k exhibits a clear monotonic increase with rising oxygen vacancy concentration, while the introduction of Sm³⁺ reduces k . Furthermore, Sm³⁺ doping can inhibit Ce⁴⁺ reduction when the V1-Sm³⁺-V2 structure with a face-diagonal relationship exists in less reduced atmospheres. It should be noted that the $\text{Sm}_{0.25}\text{Ce}_{0.75}\text{O}_{1.5}$ system is best described as Sm-doped Ce₂O₃.

Keywords: cerium oxide; oxygen vacancies; doping; electronic structures; GGA+U

1 INTRODUCTION

CeO₂-based materials are important components in solid oxide fuel cells (SOFC) due to their high ionic conductivity. Pure CeO₂ is a mixed ionic and electronic conductor. When a neutral oxygen vacancy is created in CeO₂, the two excess electrons left behind become localized on the f-level traps of two neighboring Ce atoms, formally reducing Ce⁴⁺ to Ce³⁺. To improve ionic conductivity and reduce electronic conductivity, CeO₂ is commonly doped with alkaline earth or rare earth metals such as Ca²⁺, Gd³⁺, Sm³⁺, and La³⁺, or through their co-doping.

The ionic and electronic conductivities of CeO₂ have been widely studied. Mogensen's investigation of Ca- or Gd-single-doped CeO₂ systems found that the conduction carrier depends on the atmosphere, with conductivity measured in air being primarily ionic and the electronic contribution negligible. In reduced atmospheres, both ionic conductivity (σ_i) and electronic conductivity (σ_e) are related to oxygen partial pressure. The highest σ_i for CeO₂ doped with 20 mol% Ca reaches 0.209 S/cm in H₂ atmosphere at 850 °C, which is larger than that in

air. $\text{Ce}_{1-x}\text{M}_x\text{O}_2$ ($\text{M} = \text{Gd}^{3+}$ or Sm^{3+}) at $x = 0.15, 0.2$ and $(\text{Ce}_{0.8}\text{Sm}_{0.2})_{1-0.05}\text{Ca}_{0.05}\text{O}_2$ compounds exhibit the largest ionic conductivity, with the latter achieving 0.126 S/cm at $800 \text{ }^\circ\text{C}$ in air. Additionally, the ionic radii and valence of dopants, oxygen vacancy and dopant concentrations, and dopant-vacancy and vacancy-vacancy defect associations—which decrease the number of mobile vacancies available—all affect the ionic conductivity of CeO_2 .

Theoretically, the distributions of oxygen vacancies and dopants, their formation energies, geometric and electronic structures, and oxygen ion migration energies of CeO_2 systems have been investigated at the atomic level. Among rare earth metals, Sm^{3+} and among alkaline earth metals, Ca^{2+} yield the smallest oxygen ion migration energies in CeO_2 systems, and their doping can effectively suppress electronic conductivity. Both theoretical and experimental studies have shown that Ca^{2+} and Sm^{3+} co-doping can better improve ionic conductivity. Previous work has examined the effects of different oxygen vacancy and doping concentrations on the distributions of oxygen vacancies, dopants, and Ce^{3+} in CeO_2 systems. Murgida's study demonstrated that oxygen vacancy concentration affects the distribution of Ce^{3+} , with excess electrons preferring to localize on cation sites that maximize the mean Ce^{3+} coordination number, and two vacancies tending to be second-nearest neighbors. For doped CeO_2 , studies have found that dopants prefer to occupy nearest neighbor (NN) and next-nearest neighbor (NNN) positions relative to oxygen vacancies. Independent of oxygen vacancy and doping concentrations, Ce^{3+} is distributed around the nearest neighbor oxygen vacancy.

However, a systematic theoretical study on the $\text{Ce}^{3+}/\text{Ce}^{4+}$ electronic properties of CeO_2 with different oxygen vacancy and doping concentrations is lacking. Since Ca and Sm doping can effectively restrain the reduction of Ce^{4+} to Ce^{3+} and doped CeO_2 exhibits high ionic conductivity, we investigate the effects of various oxygen vacancy and Ca-, Sm-doping concentrations on the electronic structures ($\text{Ce}^{3+}/\text{Ce}^{4+}$) of CeO_2 systems. We select different numbers of oxygen vacancies or Ca and Sm dopants in the same supercell to obtain a series of vacancy and doping concentrations, ultimately mapping the transitional area of Ce^{4+} to Ce^{3+} and the $\text{Ce}^{3+}/\text{Ce}^{4+}$ change ratio k under different conditions.

2 MODELS AND COMPUTATIONAL DETAILS

2.1 Models

CeO_2 has a fluorite-type structure ($\text{Fm}\bar{3}\text{m}$ space group) with one formula unit per primitive unit cell. In this work, we consider various numbers of oxygen vacancies or Ca^{2+} and Sm^{3+} dopants to produce different oxygen vacancy and doping concentrations in a $2 \times 2 \times 2$ supercell, namely $\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ systems with various x, y and z ($x = 0.0312, 0.0625, 0.125$ and 0.250 ; $y = 0.0312, 0.0625, 0.125$ and 0.250 ; $z = 0.0312, 0.0625, 0.125, 0.250$ and 0.50), see Table 1. When the doping concentration exceeds 0.30, $\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ systems transition from doped to alloyed compositions; however, for descriptive

convenience, we refer to all compositions as doped systems.

Doped atoms are uniformly distributed in the $2 \times 2 \times 2$ supercell to obtain systems with various doping concentrations. Hooper's study on Sm-doped CeO₂ systems found that dopant-vacancy interactions follow a hybrid NN/NNN mixture distribution as Sm³⁺ concentrations increase, with NN distributions being most favored. Our previous studies found that the first oxygen vacancy (V1) occupies the nearest neighbor (NN) position to the dopant, and we employ the same structural model here. For CeO₂ systems where more than one oxygen vacancy is nearest to a single Sm³⁺, these vacancies are selected according to the NN distribution rule, meaning the second oxygen vacancy nearest to a Sm³⁺ is introduced and designated as V2. Thus, V1 and V2 can adopt three distinct positional relationships within the cube of eight coordinated O²⁻ for one Sm³⁺: the side configuration in Fig. 1a [Figure 1: see original paper], the face-diagonal configuration in Fig. 1b, and the body-diagonal configuration in Fig. 1c.

Calculation results show that structures with the face-diagonal relationship between V1 and V2 are the most stable, consistent with pure CeO₂ systems containing double oxygen vacancies. For Ca-doped CeO₂ systems, we consider models where Ca²⁺ has either NN or NNN distribution relationships with one oxygen vacancy.

2.2 Computational Details

All calculations were performed using the Vienna ab initio simulation package (VASP). Ce 5s5p6s4f5d, O 2s2p, Ca 3s3p4s, and Sm 5s5p6s4f5d electrons were treated as valence electrons. Structures were relaxed until forces on each ion fell below 0.02 eV/Å and total energy converged within 1×10^{-6} eV. A plane-wave cutoff energy of 400 eV, a $3 \times 3 \times 3$ Monkhorst-Pack k-point mesh, and a Gaussian smearing parameter of 0.20 eV were employed.

Standard DFT formulations often fail to describe strongly correlated electrons due to inadequate treatment of electron correlation. This limitation can be corrected using a DFT+U method, where introducing a Hubbard parameter U modifies the self-interaction error and improves description of correlation effects. This methodology has been widely applied to reduced CeO₂ systems. Theoretical work has shown that the U value for Ce should exceed 5.0 eV. In this study, we use U = 6.0 eV for Ce and U = 8.0 eV with J = 0.65 eV for Sm, consistent with previous work. Exchange-correlation effects were described using the Perdew-Burke-Ernzerhof (PBE) functional within the generalized gradient approximation (GGA). The calculated lattice constant from GGA+U is 5.48 Å, in good agreement with the experimental value of 5.41 Å.

The formation energies of oxygen vacancies (E_{Vo}) for CeO_{2-x} systems, $E_{\text{Vo-Sm}}$ for Sm_xCe_{1-x}O_{2-y} systems, and $E_{\text{Vo-CaSm}}$ for Sm_xCa_yCe_{1-x-y}O_{2-z} systems are defined as:

$$E_V^o = E[X] + \frac{p}{2}E[O_2] - E[Y]$$

where $E[X]$ and $E[Y]$ are the total energies of pure or doped CeO₂ systems with and without oxygen vacancies, respectively. $E[O_2]$ is the energy of an O₂ molecule placed in a 10 Å × 10 Å × 10 Å supercell, and p is the number of oxygen vacancies.

3 RESULTS AND DISCUSSION

3.1 Geometric Structures

Optimized model structures with the lowest energies were obtained. The V1-Sm³⁺-V2 structures, where V1 and V2 (both nearest to Sm³⁺) adopt a face-diagonal positional relationship, represent the lowest-energy configurations, consistent with pure CeO₂ systems containing double oxygen vacancies. The introduction of Sm³⁺, Ca²⁺, and oxygen vacancies into CeO₂ produces significant geometric distortions, in agreement with previous work. Representative geometric structures of CeO_{1.9688}, Sm_{0.0312}Ce_{0.9688}O_{1.9688}, Sm_{0.0312}Ce_{0.9688}O_{1.9375}, Sm_{0.0312}Ca_{0.0312}Ce_{0.9376}O_{1.9688}, and Sm_{0.0312}Ca_{0.0312}Ce_{0.9376}O_{1.9375} systems are displayed in Fig. 2 [Figure 2: see original paper]. Other Sm_{*x*}Ce_{*y*}Ce_{1-*x-y*}O₂₋ systems with different oxygen vacancy and doping concentrations exhibit similar geometric distortions.

The oxygen vacancy creates a region of effective positive potential, causing neighboring O²⁻ ions to move toward the vacancy while neighboring Ce⁴⁺ ions move away. In the CeO_{1.9688} system (Fig. 2a), three of the four O²⁻ ions near the oxygen vacancy move toward it while the fourth moves away. As discussed in the electronic structure analysis below, this O²⁻ is bridged by two Ce³⁺ ions that have greater negative potential than Ce⁴⁺. In the Sm_{0.0312}Ce_{0.9688}O_{1.9688} system (Fig. 2b), the behavior is similar to CeO_{1.9688}, except that one of the two Ce³⁺ ions is replaced by Sm³⁺. In the Sm_{0.0312}Ce_{0.9688}O_{1.9375} system (Fig. 2c), the movement of four O²⁻ ions toward the two vacancies and two O²⁻ ions toward Sm³⁺ results from the combined attraction of V1 and V2 to O²⁻, meaning an O²⁻ bridged by Sm³⁺ and Ce³⁺ is not repelled, unlike in systems with a single oxygen vacancy. In the Sm_{0.0312}Ca_{0.0312}Ce_{0.9376}O_{1.9688} system (Fig. 2d), the four O²⁻ ions near oxygen vacancies are driven toward the vacancy while neighboring Ce⁴⁺ ions move away. The Sm_{0.0312}Ca_{0.0312}Ce_{0.9376}O_{1.9375} system (Fig. 2e) is similar to the Sm_{0.0312}Ce_{0.9688}O_{1.9375} system (Fig. 2c).

3.2 Oxygen Vacancy Formation Energies

Oxygen vacancy formation energies (E_{Vo}) for CeO₂₋ systems, E_{Vo-Sm} for Sm_{*x*}Ce_{1-*x*}O₂₋ systems, and $E_{Vo-CaSm}$ for Sm_{*x*}Ce_{*y*}Ce_{1-*x-y*}O₂₋ systems are listed in Table 1. E_{Vo} increases monotonically with increasing x , as shown in Table 1 and Fig. 3 [Figure 3: see original paper]. For Sm_{*x*}Ce_{1-*x*}O₂₋ and

$\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ systems, $E_{\text{Vo-Sm}}$ with a given x and $E_{\text{Vo-CaSm}}$ with given x and y are larger at higher λ , while both decrease with increasing x (or x and y) at fixed λ . Notably, the introduction of Ca and Sm promotes spontaneous oxygen vacancy formation, similar to Fergus' s findings for Sm-doped CeO_2 systems.

3.3 Electronic Structures

When oxygen vacancies form in CeO_2 , the Ce 4f state splits into two states: an empty Ce 4f_empty state and an occupied defect Ce^{3+} 4f_full state in the range between O 2p and Ce 4f_empty, consistent with our previous work and other studies, see Fig. 4 [Figure 4: see original paper]. Total electronic densities of states (DOS), partial electronic densities of states (PDOS), and localization electronic densities of states (LDOS) from the defect state of Ce^{3+} were calculated for various $\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ systems with different x , y , and λ , as shown in Figs. 4 and 5 [Figure 5: see original paper].

3.3.1 Excess Electron Distribution The PDOS of the Ce 4f state for $\text{CeO}_{1.9688}$, $\text{CeO}_{1.9375}$, $\text{CeO}_{1.875}$, and $\text{CeO}_{1.75}$ systems (Fig. 4a) shows a new peak appearing in the -1.2 to 0 eV range for $\text{CeO}_{1.9688}$, $\text{CeO}_{1.9375}$, and $\text{CeO}_{1.875}$ systems, and in the -0.80 to 0.60 eV range for the $\text{CeO}_{1.75}$ system. These peaks are fully occupied by Ce^{3+} electrons. Compared to the $\text{CeO}_{1.75}$ system, the Fermi level of the other systems shifts upward by approximately 0.50 eV due to decreased oxygen vacancy concentration. Excess electrons arise from oxygen vacancies and localize on the f-level traps of neighboring Ce atoms, as visually confirmed by the corresponding LDOS of the defect Ce^{3+} state in Figs. 4b-e. These findings are consistent with theoretical studies of $\text{CeO}_{1.9843}$, $\text{CeO}_{1.9687}$, $\text{CeO}_{1.9375}$, and $\text{CeO}_{1.875}$ and $\text{CeO}_2(111)$, (110), and (100) surfaces.

The PDOS of the Ce 4f state for $\text{Sm}_x\text{Ce}_{1-x}\text{O}_2$ ($0.0312 \leq x \leq 0.25$, $0.0312 \leq \lambda \leq 0.5$) systems (Fig. 5a) shows similar behavior to CeO_2 systems, with a new peak appearing in the O 2p-Ce 4f_empty range. The corresponding LDOS of the defect Ce^{3+} state (Figs. 5c-j) reveals that excess electrons from oxygen vacancies localize on the f-level traps of neighboring Ce atoms, as in CeO_2 systems. For the $\text{Sm}_{0.0312}\text{Ce}_{0.9688}\text{O}_{1.9375}$ system (Fig. 5j), two oxygen vacancies in the $2 \times 2 \times 2$ supercell should induce four charge-compensating cations; however, calculations find only three, possibly due to the V1-Sm³⁺-V2 structure with a face-diagonal relationship in less reduced atmospheres, where doped Sm³⁺ can suppress Ce⁴⁺ reduction.

To better control excess electrons, we explored $\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ ($0.0312 \leq x \leq 0.25$, $0.0312 \leq y \leq 0.25$, $0.0312 \leq \lambda \leq 0.5$) systems with various x , y , and λ . The DOS in Fig. 5b shows that except for the $\text{Sm}_{0.0625}\text{Ca}_{0.0625}\text{Ce}_{0.875}\text{O}_{1.875}$ system, $\text{Sm}_x\text{Ca}_y\text{Ce}_{1-x-y}\text{O}_2$ systems contain no Ce^{3+} . Thus, to avoid Ce^{3+} formation, both x and y must exceed 0.0625 when $\lambda = 0.125$, or λ must be less than 0.125 when $x = y = 0.0625$. For these systems, the corresponding LDOS

of the defect Ce^{3+} state (Figs. 5c-k) shows that excess electrons from oxygen vacancies localize on the f-level traps of neighboring Ce atoms, as in CeO_2 -systems.

3.3.2 Transitional Area of Ce^{4+} to Ce^{3+} and Ce^{3+}/Ce^{4+} Change Ratio k To visually understand the transitional area of Ce^{4+} to Ce^{3+} in $Sm_xCa_yCe_{1-x-y}O_2$ - systems with various oxygen vacancy and doping concentrations, we plotted the transitional areas for $Sm_xCe_{1-x}O_2$ - (Fig. 6a [Figure 6: see original paper]) and $Sm_xCa_yCe_{1-x-y}O_2$ - systems (Fig. 6b). Fig. 6a shows that no Ce^{3+} exists in $Sm_xCe_{1-x}O_2$ - systems with $x = 0.167$ and 0.0833 , while Ce^{3+} exists in systems with $x = 0.0312$ and 0.0312 , $x = 0.0625$ and 0.0625 , $x = 0.125$ and 0.125 , $x = 0.167$ and 0.166 , and $x = 0.250$ and 0.250 .

For $Sm_xCa_yCe_{1-x-y}O_2$ - systems with $y = 0$, substitution of Ce^{4+} by Ca^{2+} creates two excess electrons while substitution by Sm^{3+} creates one excess electron. To comprehensively reflect the doping concentration effect on the Ce^{4+} to Ce^{3+} transition ratio, we transform the effect of one Ca^{2+} to that of two Sm^{3+} (ignoring ionic radius differences) and plot the corresponding transitional area in Fig. 6b. This shows that no Ce^{3+} exists in $Sm_xCa_yCe_{1-x-y}O_2$ - systems with $x = 0.0936$ and 0.0625 , $x = 0.334$ and 0.166 , $x = 0.375$ and 0.250 , and $x = 0.750$ and 0.50 , while Ce^{3+} exists in systems with $x = 0.187$ and 0.125 .

The Ce^{3+}/Ce^{4+} change ratio k in $Sm_xCe_{1-x}O_2$ - systems with various oxygen vacancy and doping concentrations is shown in Fig. 6c. The ratio k exhibits a clear monotonic increase with vacancy concentration, while Sm^{3+} introduction reduces k , as seen in $Sm_{0.0312}Ce_{0.9688}O_{1.9375}$, $Sm_{0.0625}Ce_{0.9375}O_{1.875}$, $Sm_{0.125}Ce_{0.875}O_{1.75}$, $Sm_{0.166}Ce_{0.834}O_{1.668}$, and $Sm_{0.25}Ce_{0.75}O_{1.5}$ systems. Thus, Sm^{3+} doping suppresses Ce^{4+} reduction to Ce^{3+} . Notably, the $Sm_{0.25}Ce_{0.75}O_{1.5}$ system can be considered as Sm-doped Ce_2O_3 because it shows no unoccupied defect states between the occupied Ce^{3+} states and unoccupied $Ce\ 4f_{empty}$ (Fig. 5a), consistent with other theoretical work.

4 CONCLUSION

This work investigates the influence of Sm^{3+} single-doping and Ca^{2+} and Sm^{3+} co-doping, along with oxygen vacancies and their concentrations, on the geometric and electronic structures of $Sm_xCa_yCe_{1-x-y}O_2$ - systems. Geometric structure analysis shows that V1- Sm^{3+} -V2 structures with a face-diagonal positional relationship between V1 and V2 (both nearest to Sm^{3+}) are readily obtained, consistent with pure CeO_2 systems. Electronic structure studies reveal that oxygen vacancies contribute excess electrons that localize on the f-level traps of neighboring Ce atoms in $Sm_xCa_yCe_{1-x-y}O_2$ - systems. The Ce^{3+}/Ce^{4+} change ratio k depends on both dopant and oxygen vacancy concentrations. Additionally, the presence of V1- Sm^{3+} -V2 structures with face-diagonal relationships in less reduced atmospheres may enable doped Sm^{3+} to suppress Ce^{4+} reduction. It should be noted that the $Sm_{0.25}Ce_{0.75}O_{1.5}$ system can be considered as Sm-doped Ce_2O_3 .

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